

# Heavy element nucleosynthesis: the r-process and the *vr*-process

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FRIB-TA Topical Program: Theoretical  
Justifications and Motivations for Early High-  
Profile FRIB Experiments  
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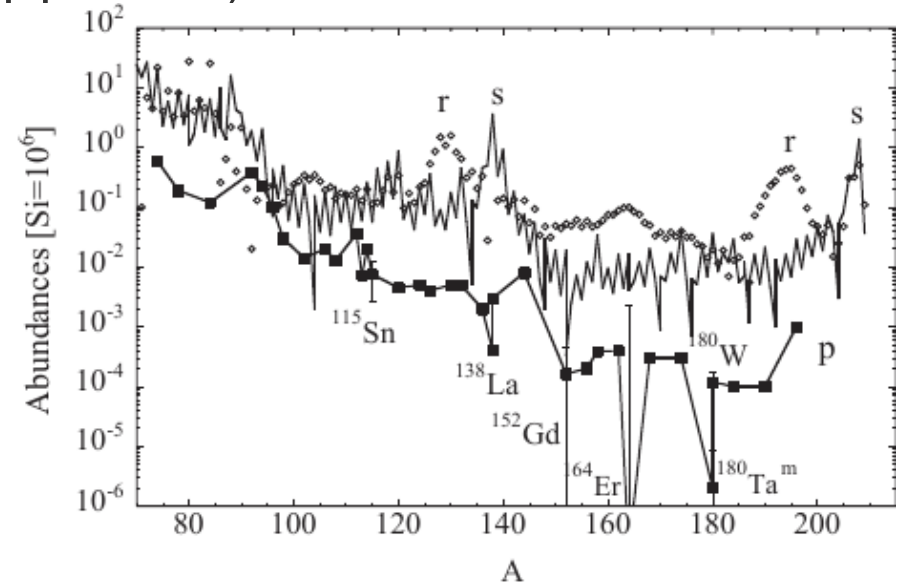
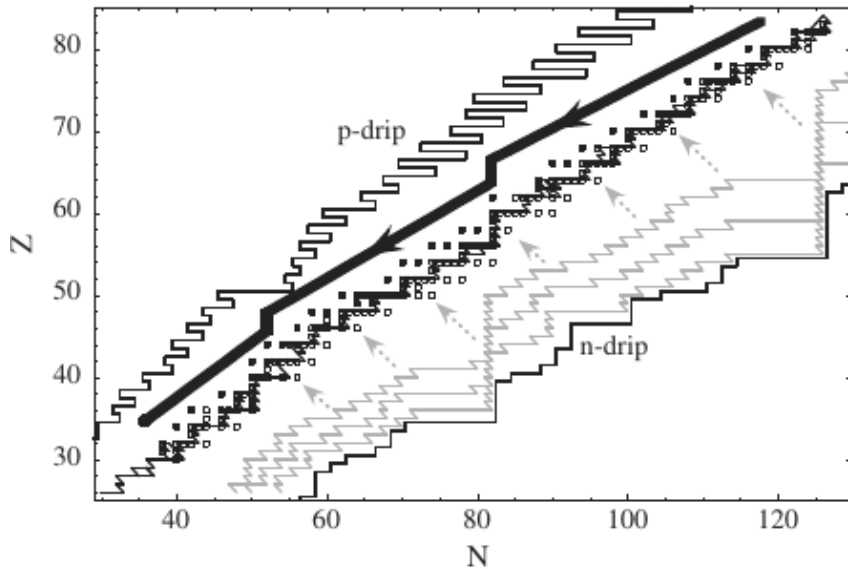


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Two main processes contribute to the nucleosynthesis beyond Iron: s-process, r-process and p-process ( $\gamma$ -process)



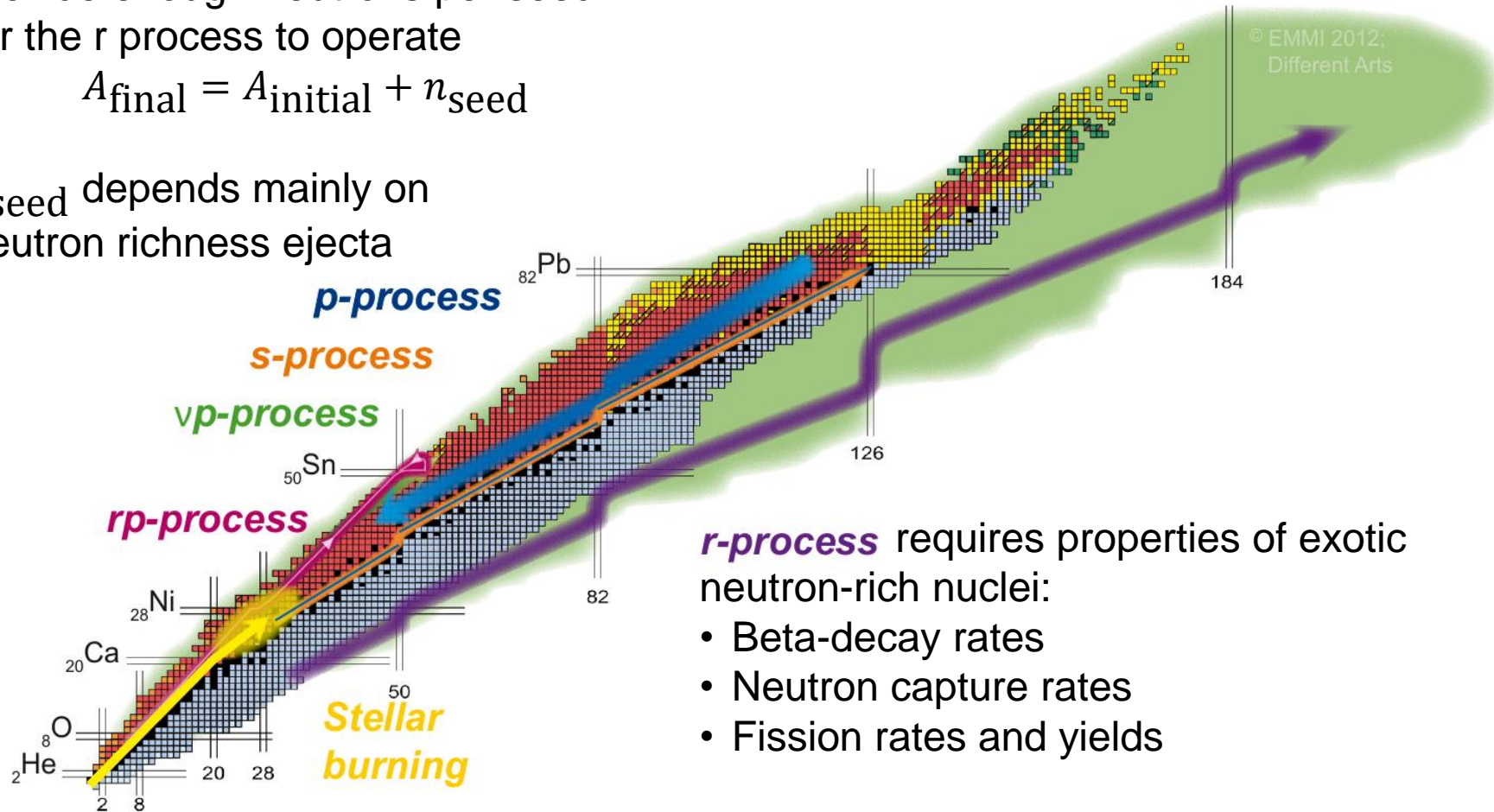
- s process: low neutron densities,  $n_n = 10^{10-12} \text{ cm}^{-3}$ ,  $\tau_n > \tau_\beta$   
(site: intermediate mass stars)
- r process: large neutron densities,  $n_n > 10^{20} \text{ cm}^{-3}$ ,  $\tau_n \ll \tau_\beta$   
(site: binary neutron star mergers?)
- Additional process(es) required to produce neutron-deficient p-nuclei
  - $\gamma$ -process: photodissociation material enriched by s-process
  - vp-process:  $(p,\gamma)$  and  $(n,p)$  reactions catalysed by  $\bar{\nu}_e + p \rightarrow n + e^+$

# R process needs

Astrophysical environment should provide enough neutrons per seed for the r process to operate

$$A_{\text{final}} = A_{\text{initial}} + n_{\text{seed}}$$

$n_{\text{seed}}$  depends mainly on neutron richness ejecta



**r-process** requires properties of exotic neutron-rich nuclei:

- Beta-decay rates
- Neutron capture rates
- Fission rates and yields

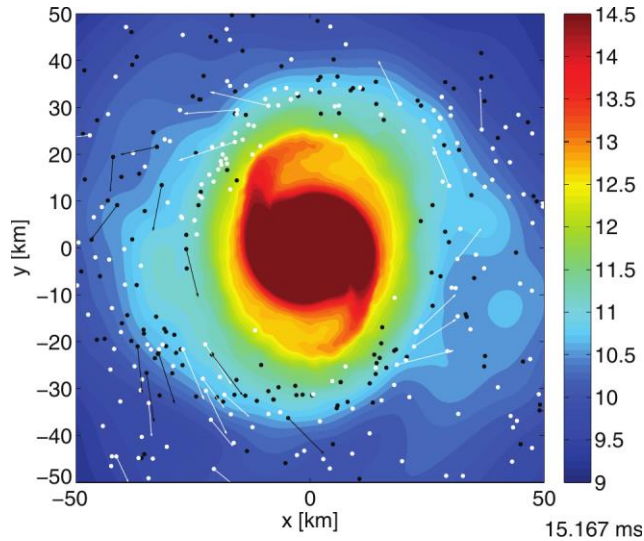
Benchmark against observations:

- Indirect: Solar and stellar abundances (contribution many events, chemical evol.)
- Direct: Kilonova electromagnetic emission (single event, sensitive Atomic and Nuclear Physics)

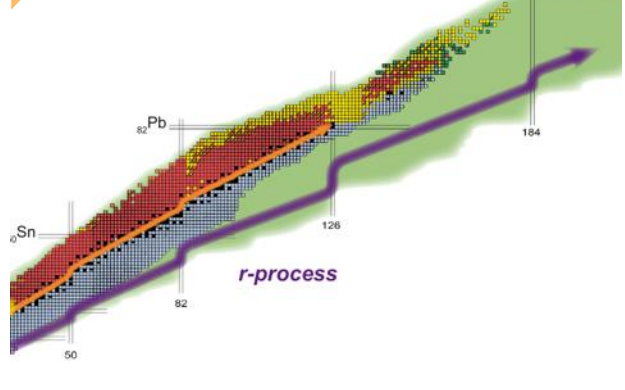
# Pipeline for r-process in mergers

## Simulations

Bauswein et al, ApJ 773, 78 (2013)

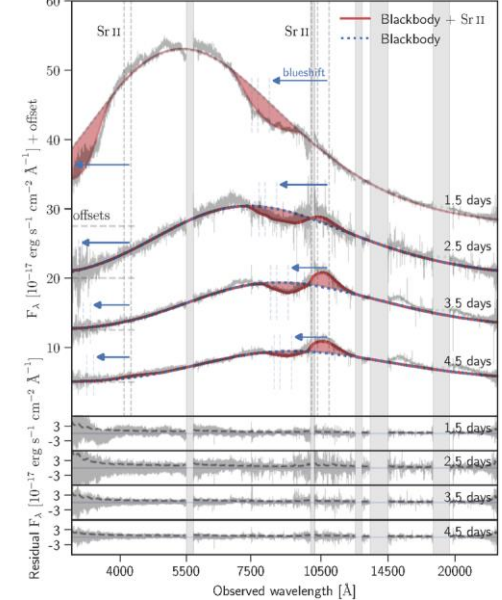


## Nucleosynthesis



## Light curve and spectra modelling

Watson et al, Nature 574, 497 (2019)



- Properties ejecta: proton-to-nucleon ratio ( $Y_e$ )
- Role of equation of state
- Role of neutrinos

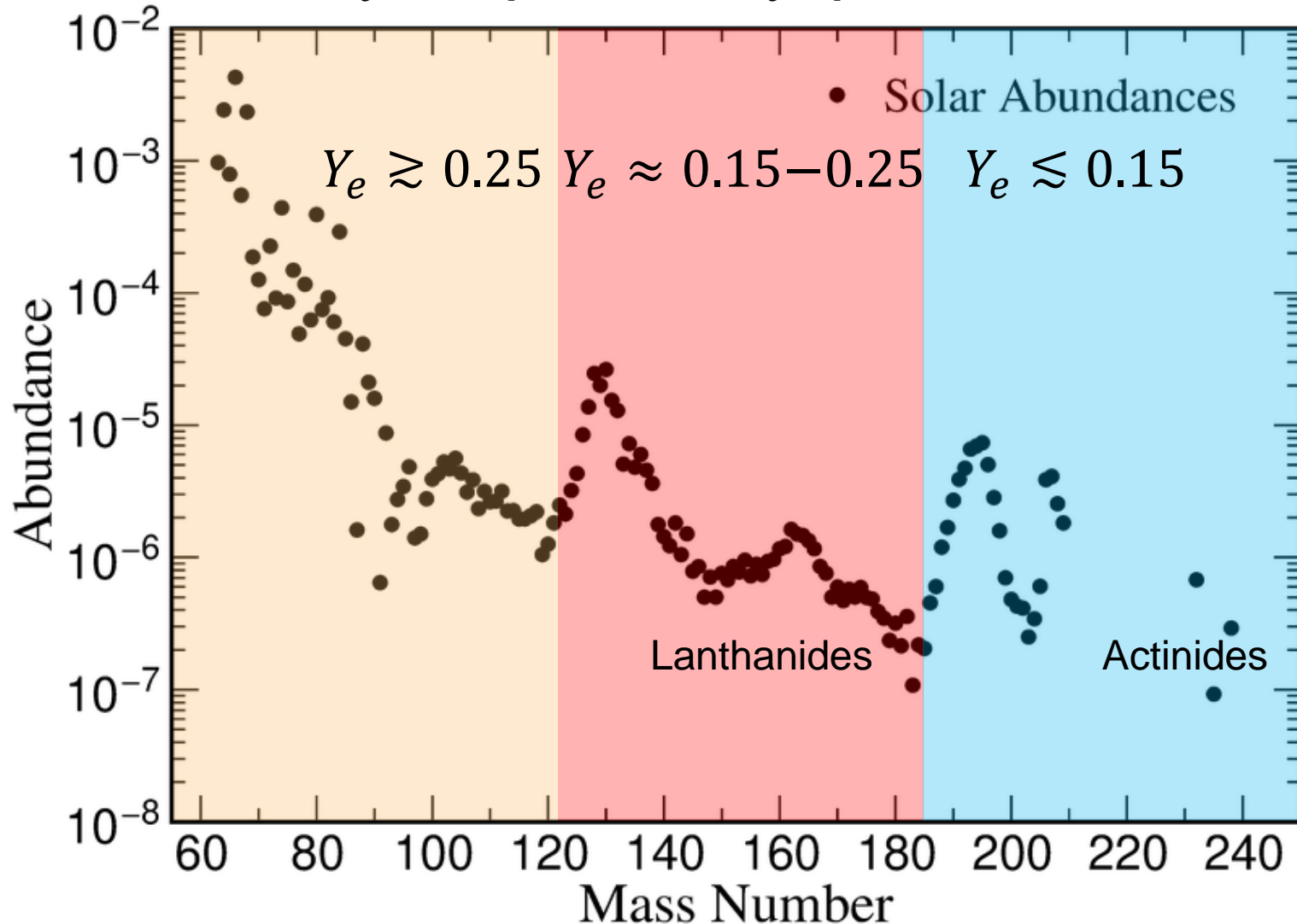
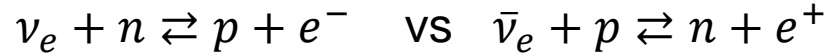
- Physics of neutron-rich and heavy nuclei

- Radioactive energy deposition
- Thermalization decay products (Barnes+ 2016, Kasen+ 2019)
- Spectra formation: atomic data depends on ejecta evolution (LTE vs NLTE)

- Which r-process elements are produced in mergers?
- Are mergers the (main) r-process site?

# Nucleosynthesis dependence on $Y_e$

Nucleosynthesis mainly sensitive to proton-to-nucleon ratio,  $Y_e = n_p / (n_n + n_p)$

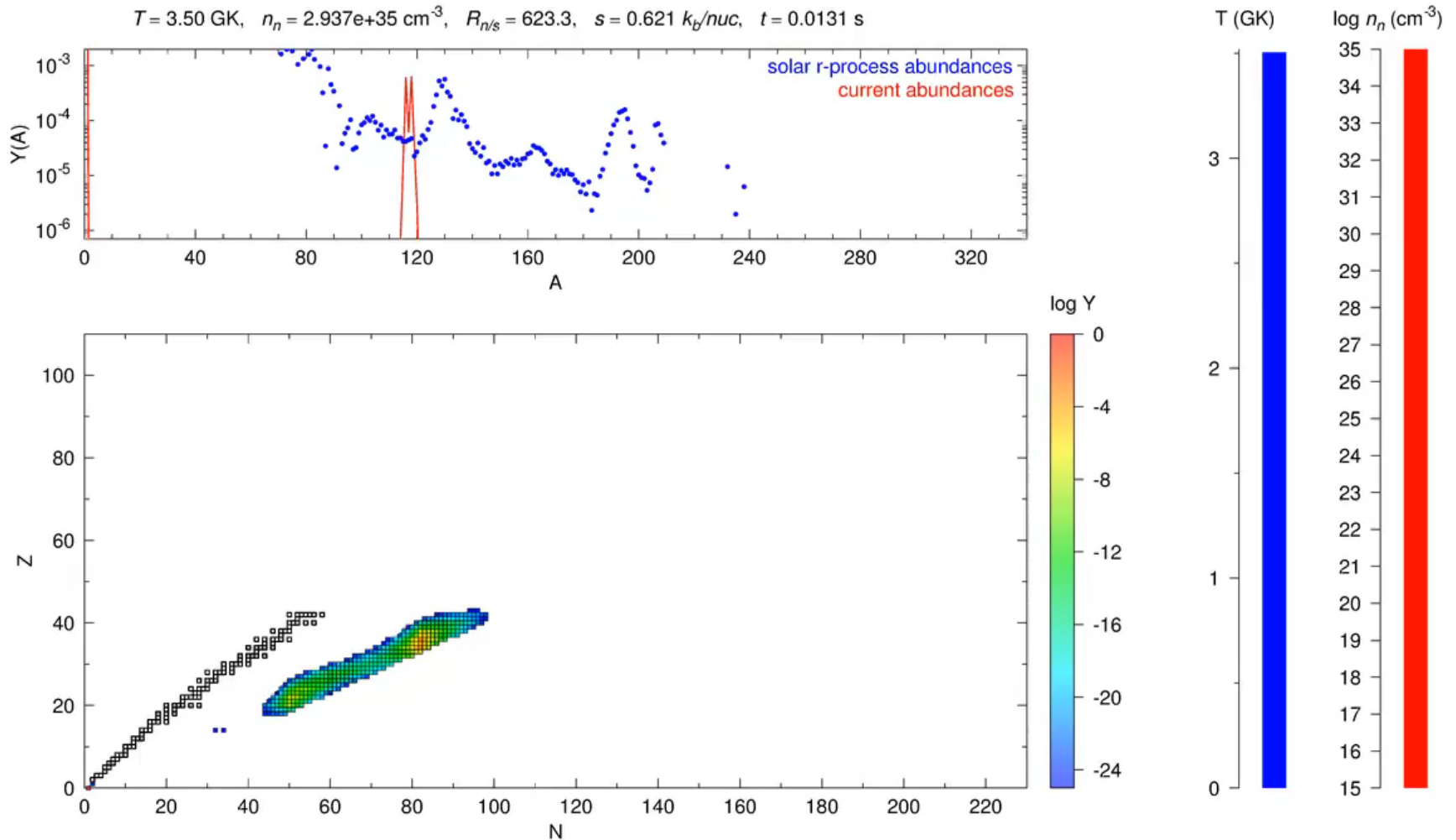


# Phases during the operation of the r-process

- **Weak freeze-out:** proton-to-nucleon ratio determined by (anti)neutrino absorption and their inverses
- **Seed production:** Charged particle reactions operating for  $T \gtrsim 2 \text{ GK}$  produce the seed nuclei and neutrons
- **Neutron-capture phase:** neutrons are captured on the available seed nuclei on a typical times of  $\sim 1 \text{ s}$ . Different equilibria are achieved:
  - $(n, \gamma) \rightleftharpoons (\gamma, n)$  equilibrium defines the r-process path that is mainly sensitive to the nuclear masses
  - Beta-flow equilibrium: abundance given element is proportional to the beta-decay half-lives. R-process peaks associated to nuclei with longest half-lives.
- **Freeze-out and decay to stability:** fully dynamical phase in which competition between neutron-captures, beta-decay (and fission) determines the final abundance pattern. Most sensitive phase to the nuclear input

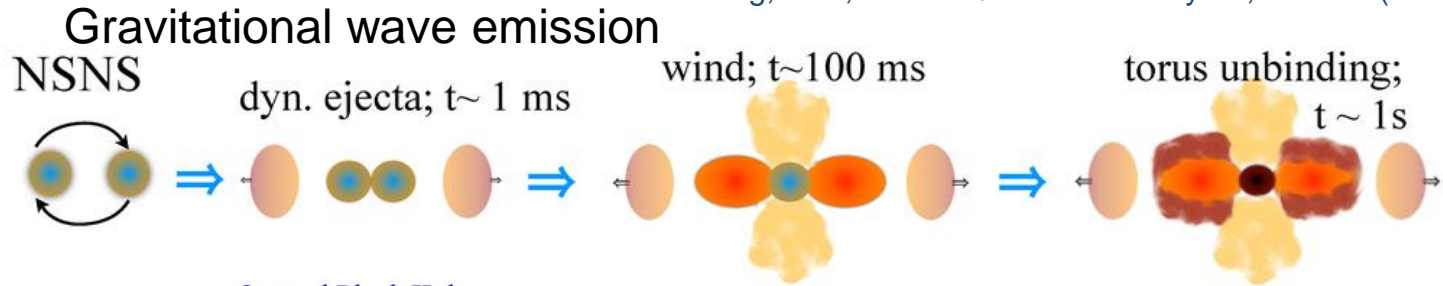
# R-process operation

Heavy elements produced by the r-process. Radioactive decay liberates energy



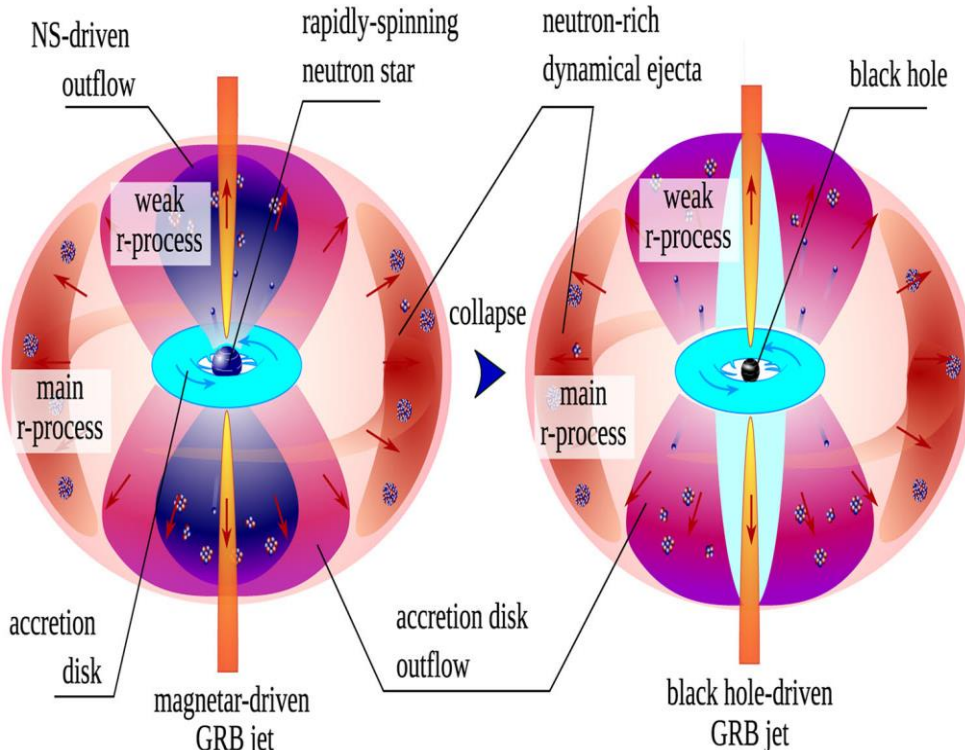
# Neutron star mergers: Different ejection mechanisms

S. Rosswog, et al, Class. Quantum Gravity 34, 104001 (2017).

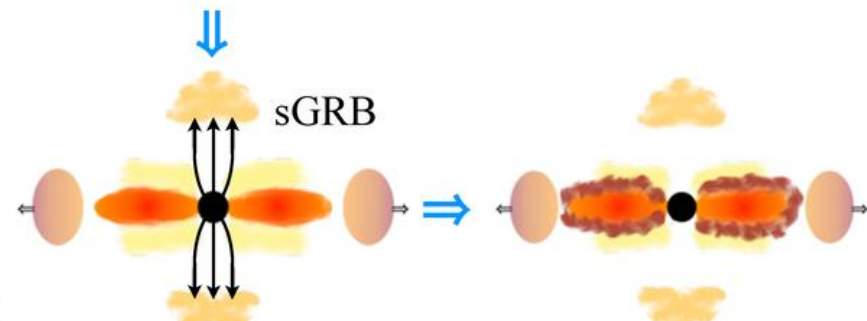


Central Neutron Star

Central Black Hole



BH formation



Two sources of ejecta:

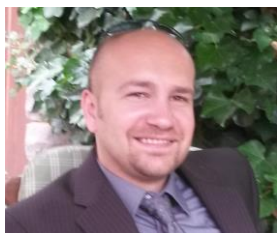
- Dynamical during the early phases of the merger ( $M \lesssim 0.01 M_{\odot}$ )
- Accretion disc on longer timescales ( $M \lesssim 0.05 M_{\odot}$ )
- Lifetime neutron-star determines impact neutrinos

S. Rosswog and O. Korobkin, Annalen Der Physik **2022**, 2200306 (2022).

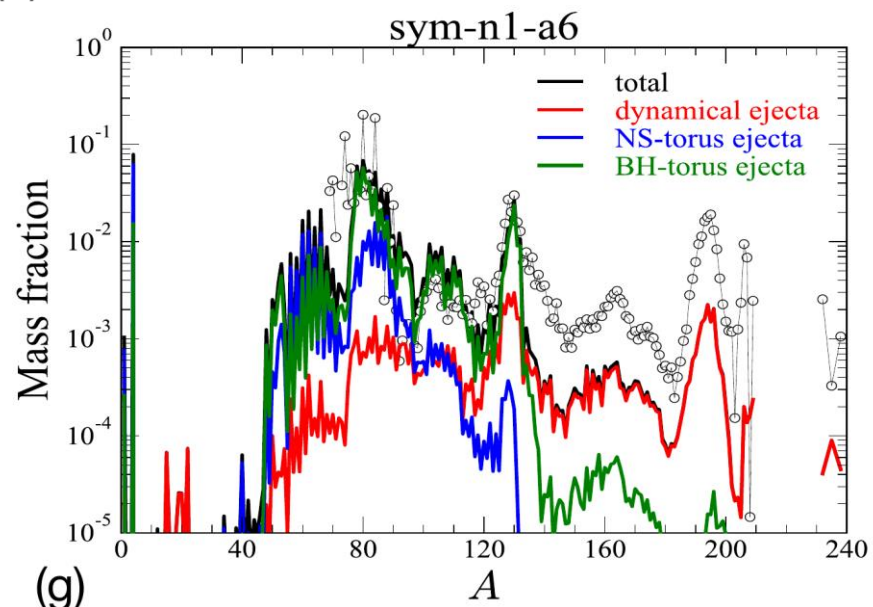
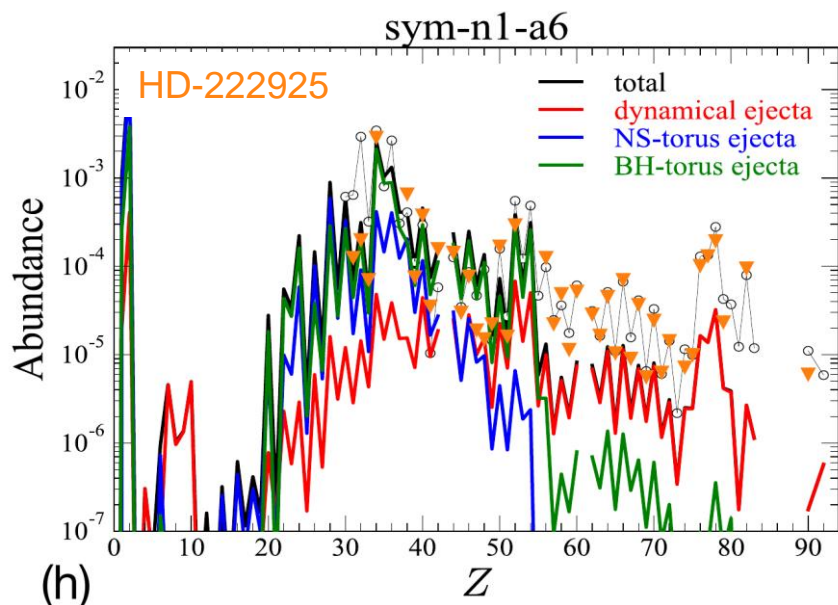
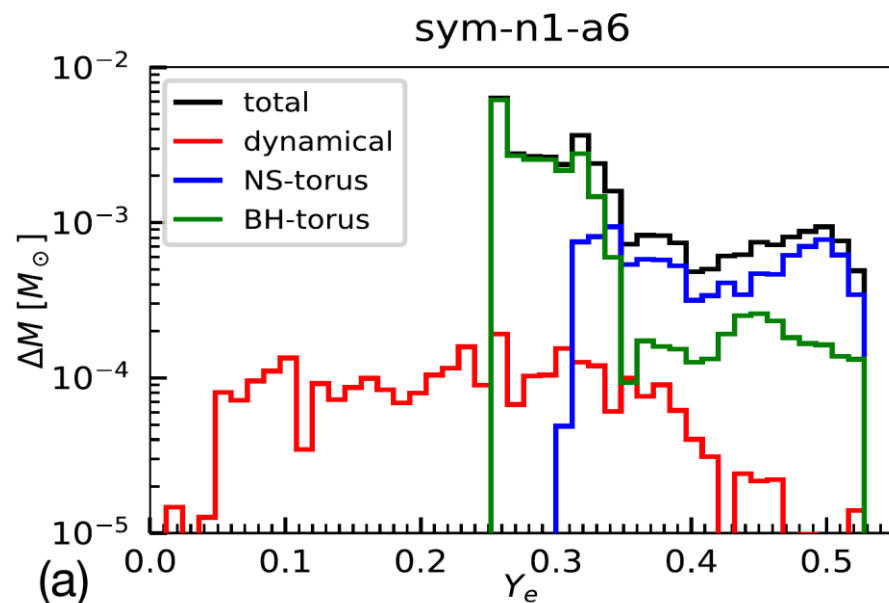


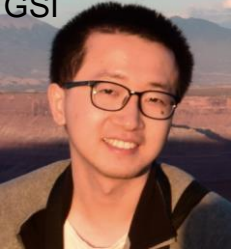
# Long term merger simulations

First long-term simulations with neutron star lifetimes 0.1-1 s and describe all components of the ejecta: dynamical, NS-remnant ejecta, and final viscous ejecta from BH torus.



Just et al, arXiv:2302.10928





arXiv:2305.11050v1 [astro-ph.HE] 18 May 2023

## Production of $p$ -nuclei from $r$ -process seeds: the $\nu r$ -process

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Schlossgartenstraße 2, D-64289 Darmstadt, Germany

<sup>3</sup>Max Planck Institute for Astrophysics, Karl-Schwarzschild-Straße 1, D-85748 Garching, Germany

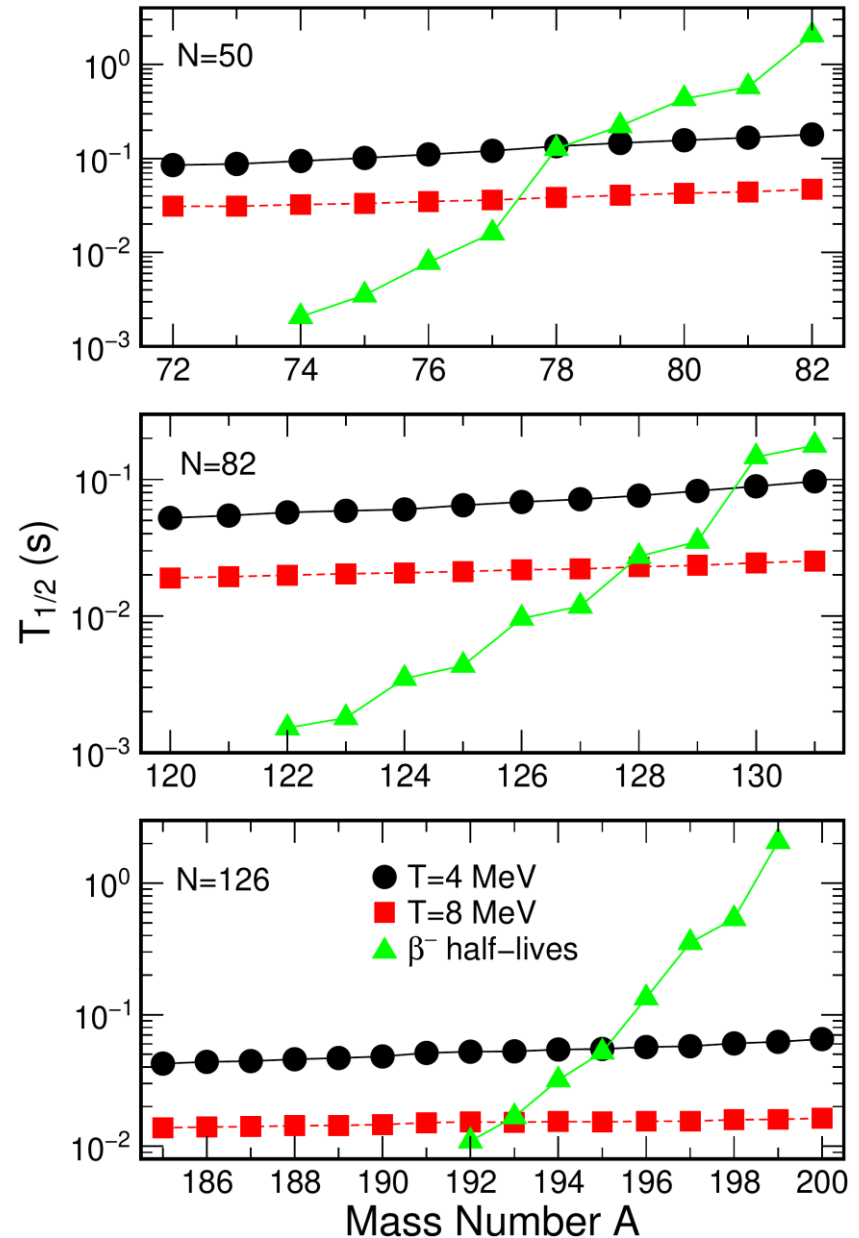
(Dated: May 19, 2023)

We present a *new* nucleosynthesis process that may take place on neutron-rich ejecta experiencing an intensive neutrino flux. The nucleosynthesis proceeds similarly to the standard  $r$ -process, a sequence of neutron-captures and beta-decays, however with charged-current neutrino absorption reactions on nuclei operating much faster than beta-decays. Once neutron capture reactions freeze-out the produced  $r$ -process neutron-rich nuclei undergo a fast conversion of neutrons into protons and are pushed even beyond the  $\beta$ -stability line producing the neutron-deficient  $p$ -nuclei. This scenario, which we denote as the  $\nu r$ -process, provides an alternative channel for the production of  $p$ -nuclei and the short-lived nucleus  $^{92}\text{Nb}$ . We discuss the necessary conditions posed on the astrophysical site for the  $\nu r$ -process to be realized in nature. While these conditions are not fulfilled by current neutrino-hydrodynamic models of  $r$ -process sites, future models, including more complex physics and a larger variety of outflow conditions, may achieve the necessary conditions in some regions of the ejecta.

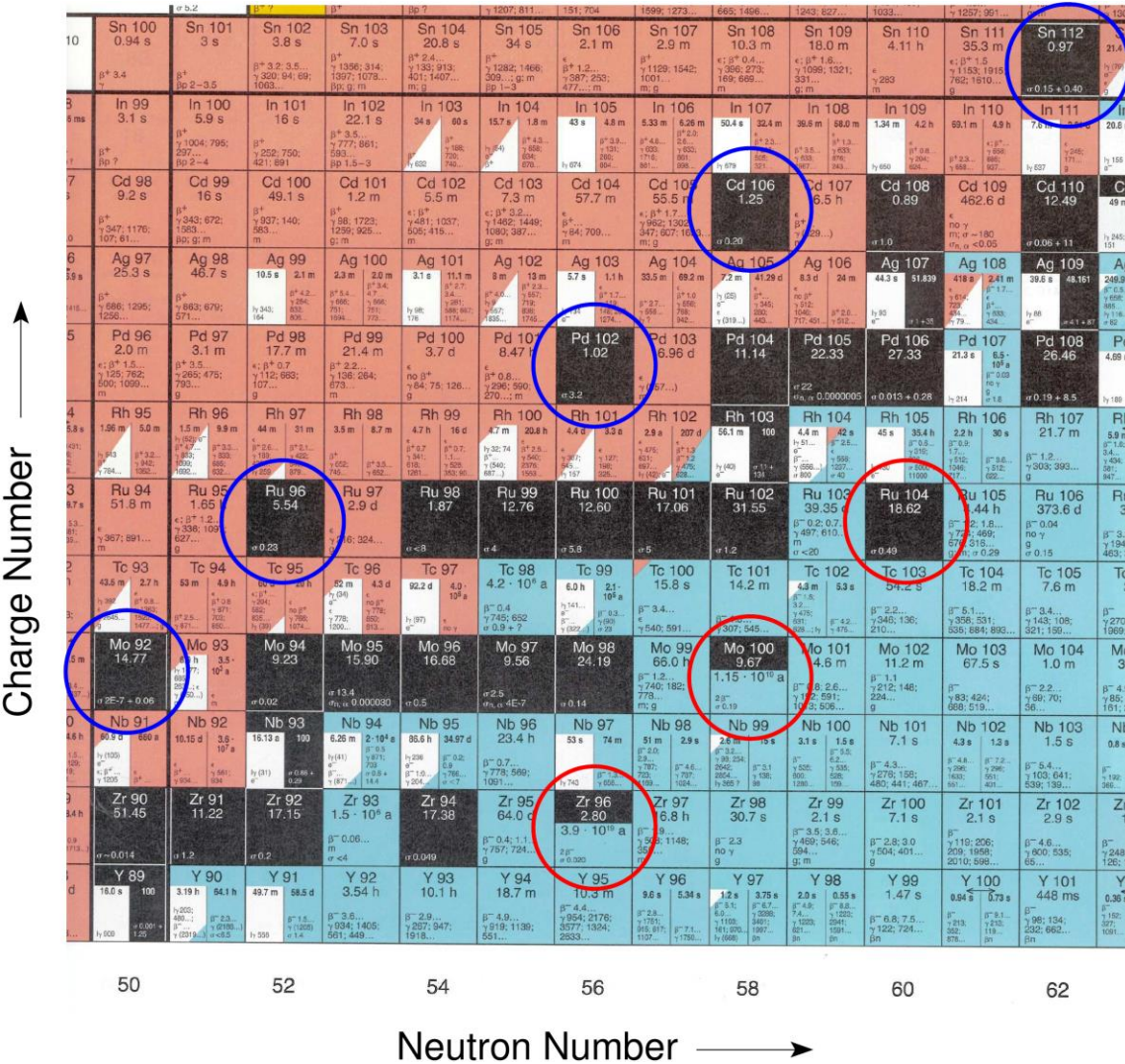
- A new nucleosynthesis process that may operate in binary neutron star mergers under strong neutrino fluxes when nuclei are present: charged-current neutrino-nucleus reactions faster than  $\beta^-$  decays.
- Novel mechanism for production of  $p$ -nuclei from neutron-rich nuclei.

# Role of neutrinos in r-process

- Large (anti)neutrino fluxes drive composition to  $Y_e \sim 0.5$  during alpha-particle formation (Meyer et al, 1995)
- large neutrino fluxes during the phase of neutron captures erode r-process peaks related to long beta-decays (Langanke, GMP, 2003)
- $\nu_e$  absorption cross sections  $\sim (N - Z)$
- What will be the resulting nucleosynthesis?



# Possible source of light p-nuclei and $^{92}\text{Nb}$



$\gamma$ -process fails to produce light p-nuclei  $^{92,94}\text{Mo}$  and  $^{96,98}\text{Ru}$  in solar proportions

Supernova neutrino winds:

- Ejecta with  $Y_e \sim 0.48$  produce  $^{92}\text{Mo}$
- $\nu p$ -process ( $Y_e \gtrsim 0.55$ ) produces  $^{94}\text{Mo}$ ,  $^{96,98}\text{Ru}$ .

Long-lived  $^{92}\text{Nb}$  present in early solar system. Cannot be produced by the  $\nu p$ -process

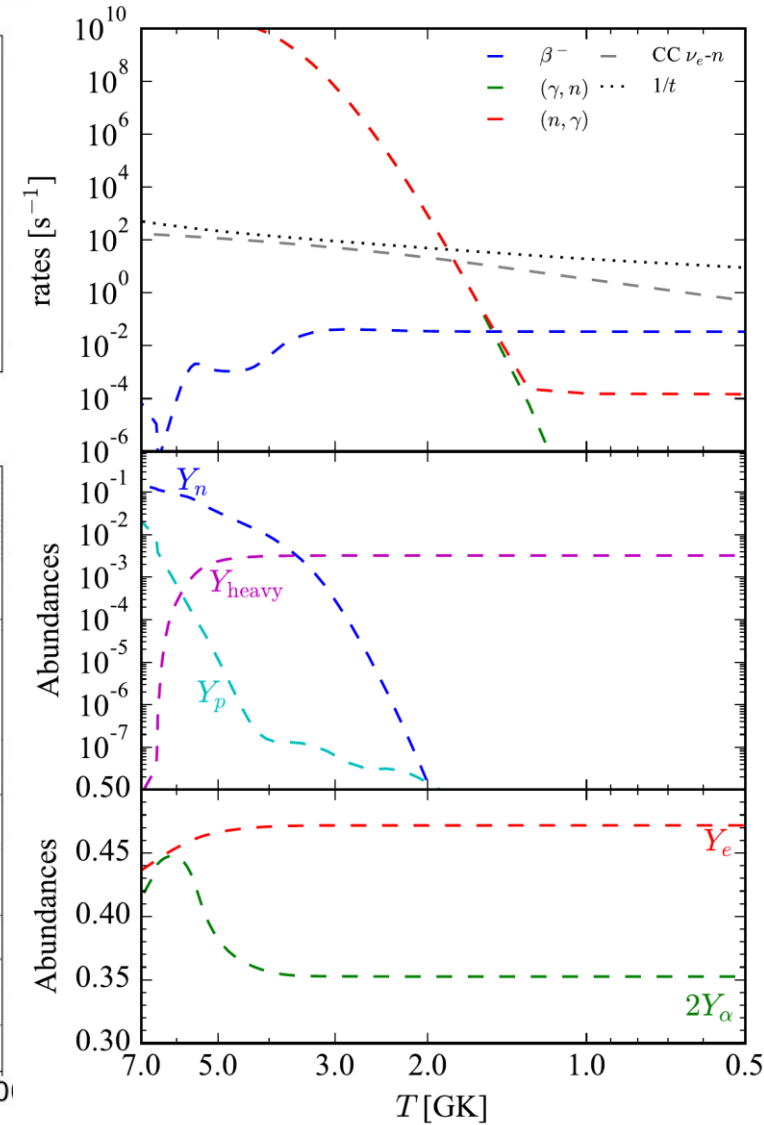
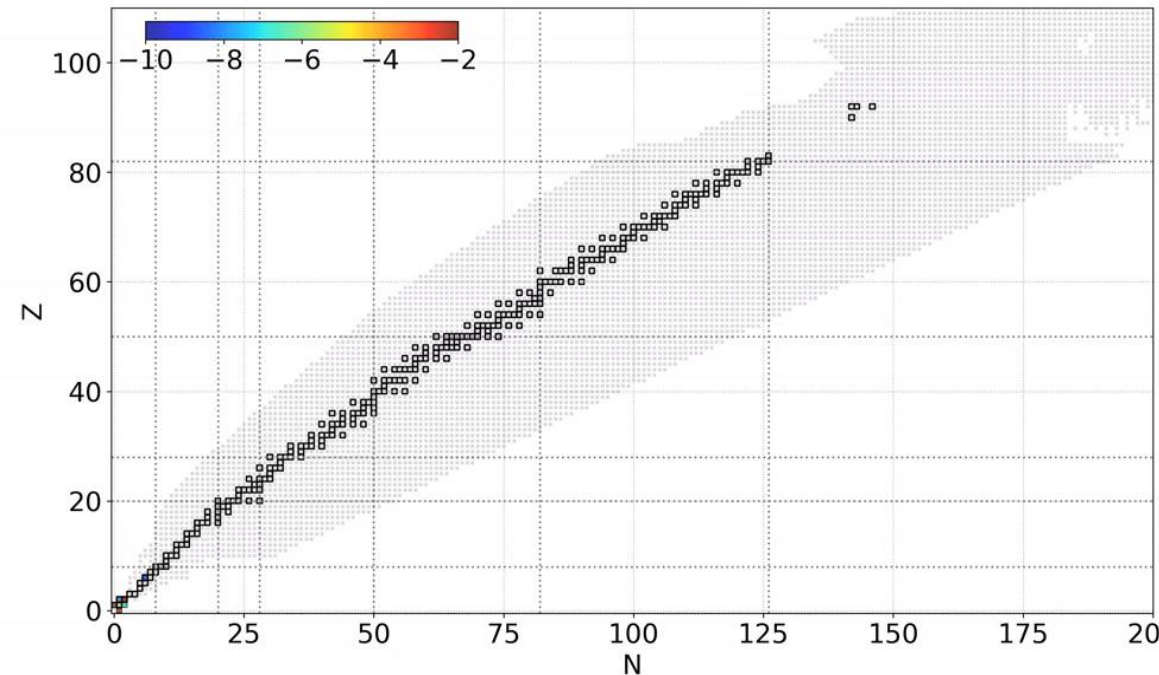
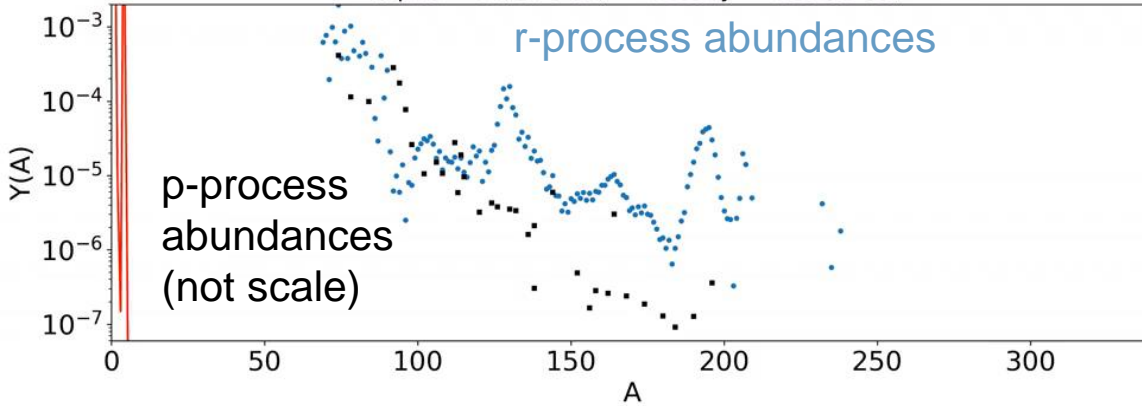
Can we produce all those nuclei in the same environment including heavier p-nuclei?

# Phases during the operation of the vr-process

- **Seed production:** Strong neutrino fluxes drive material to  $Y_e \sim 0.5$
- **Neutron-capture phase:** neutrons are used relatively fast by two competing mechanisms:
  - $n(\nu_e, e^-)p$  converts neutrons into protons that are captured in medium mass nuclei
  - $A(\nu_e, e^-X)$   $X = n, p, \alpha$  speeds up the decay of nuclei and the build up of heavy nuclei
- **Fast “decay” to stability and beyond:**  
 $A(\nu_e, e^-X)$  reactions drive material to beta-stability and beyond
  - Neutrons, protons and alphas produced by both charged-current and neutral current spallation reactions.
  - Protons and alphas captured mainly in light nuclei
  - Equilibrium between  $A(\nu_e, e^-X)$  and  $A(n, \gamma)$  determines final abundance

# Nucleosynthesis (no neutrino-nucleus)

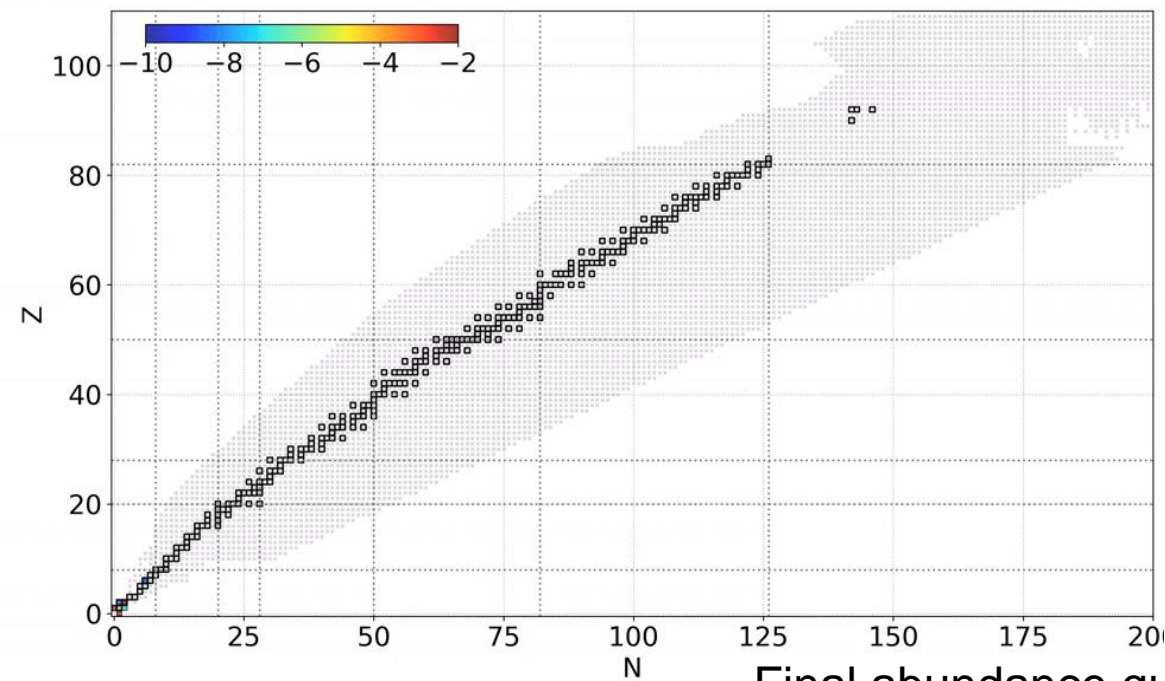
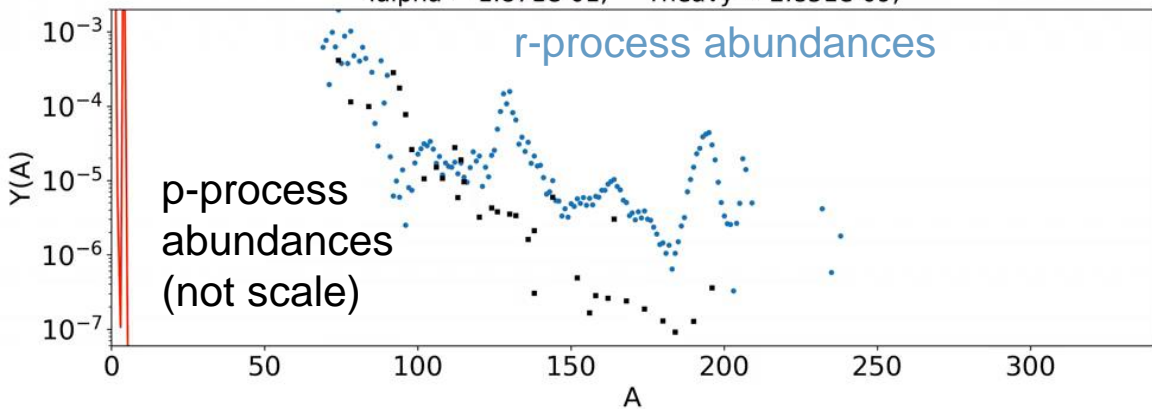
$i: 233; \quad t = 1.336e-03 \text{ s}; \quad T = 7.733e+00 \text{ GK}; \quad \rho = 1.954e+06 \text{ g cm}^{-3};$   
 $n_n = 2.657e+29 \text{ cm}^{-3}; \quad R_{n/s} = 1.170e+08; \quad S = 8.353e+01 \text{ kb/nuc};$   
 $Y_e = 4.228e-01; \quad Y_n = 2.258e-01; \quad Y_p = 7.141e-02;$   
 $Y_{\alpha} = 1.757e-01; \quad Y_{\text{heavy}} = 1.930e-09;$



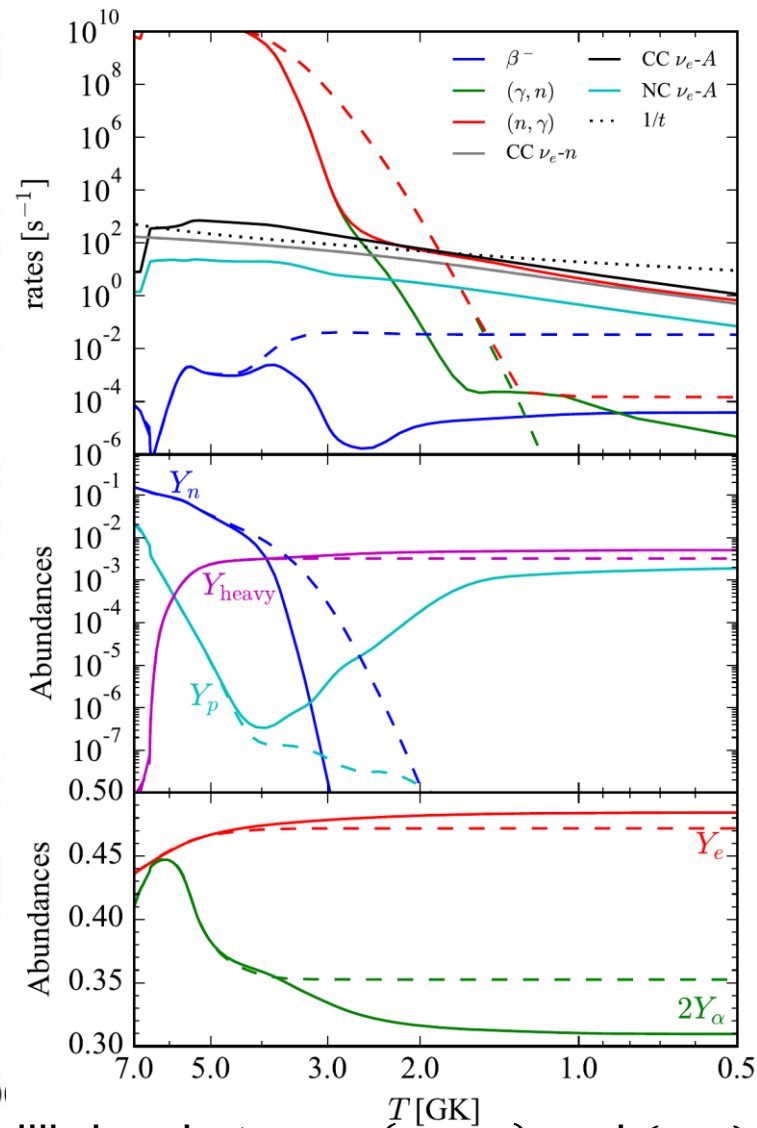
# Nucleosynthesis (with neutrino-nucleus)

$i: 232; \quad t = 1.494e-03 \text{ s}; \quad T = 7.532e+00 \text{ GK}; \quad \rho = 1.791e+06 \text{ g cm}^{-3};$   
 $n_n = 2.152e+29 \text{ cm}^{-3}; \quad R_{n/s} = 7.002e+07; \quad S = 8.355e+01 \text{ kb/nuc};$   
 $Y_e = 4.262e-01; \quad Y_n = 1.996e-01; \quad Y_p = 5.211e-02;$   
 $Y_\alpha = 1.871e-01; \quad Y_{\text{heavy}} = 2.851e-09;$

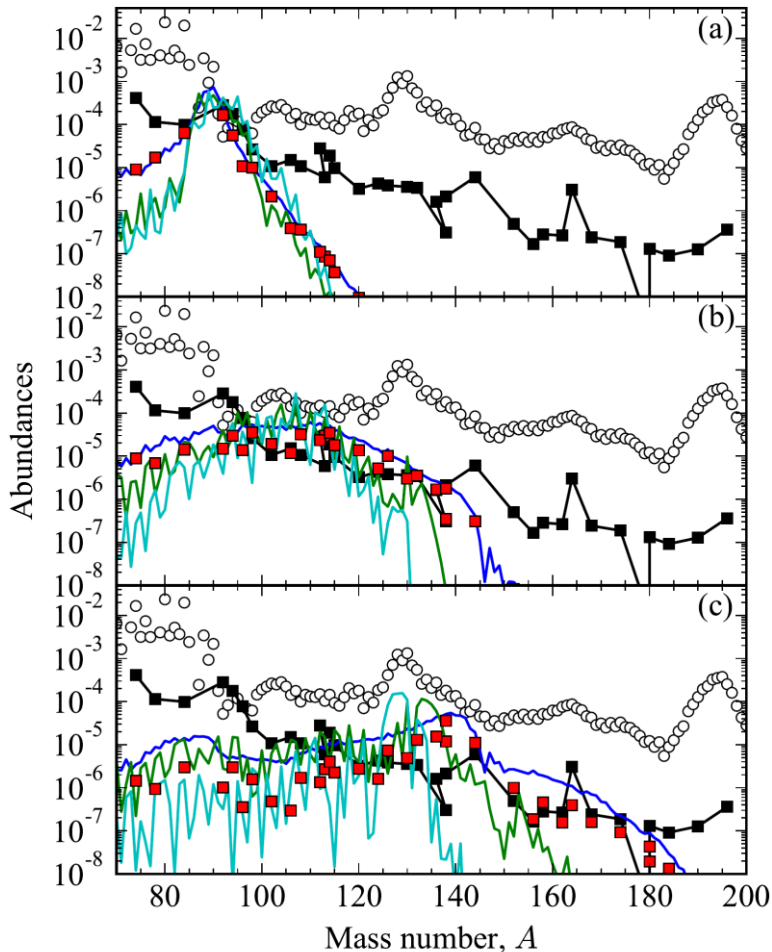
$\nu$ - $A$  cross sections from Sieverding, et al, ApJ 865, 143 (2018).



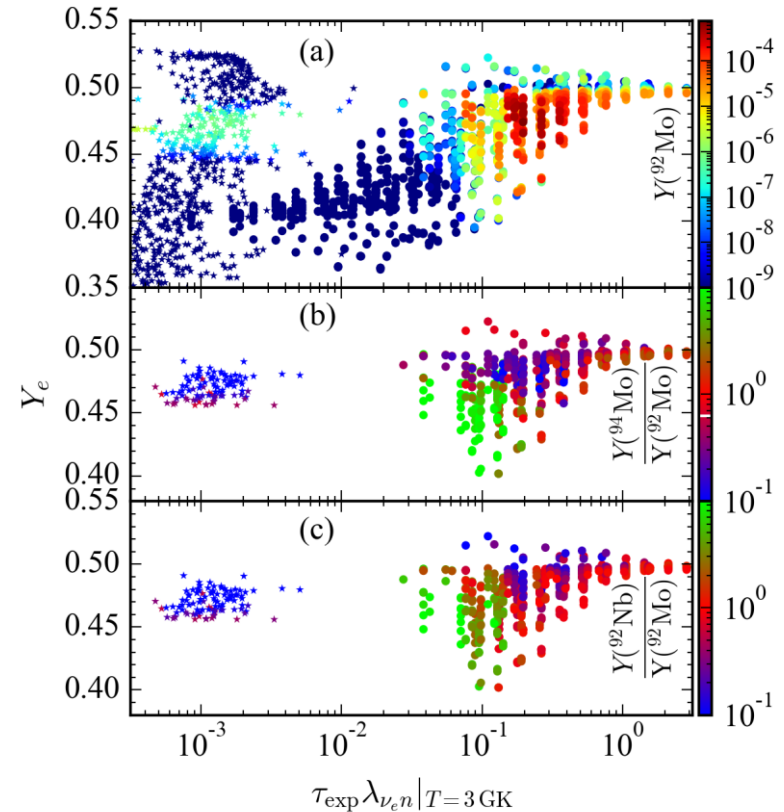
Final abundance equilibrium between  $(\nu_e, e^-)$  and  $(n, \gamma)$



Increasing neutrino fluence allows to produce heavier p-nuclei



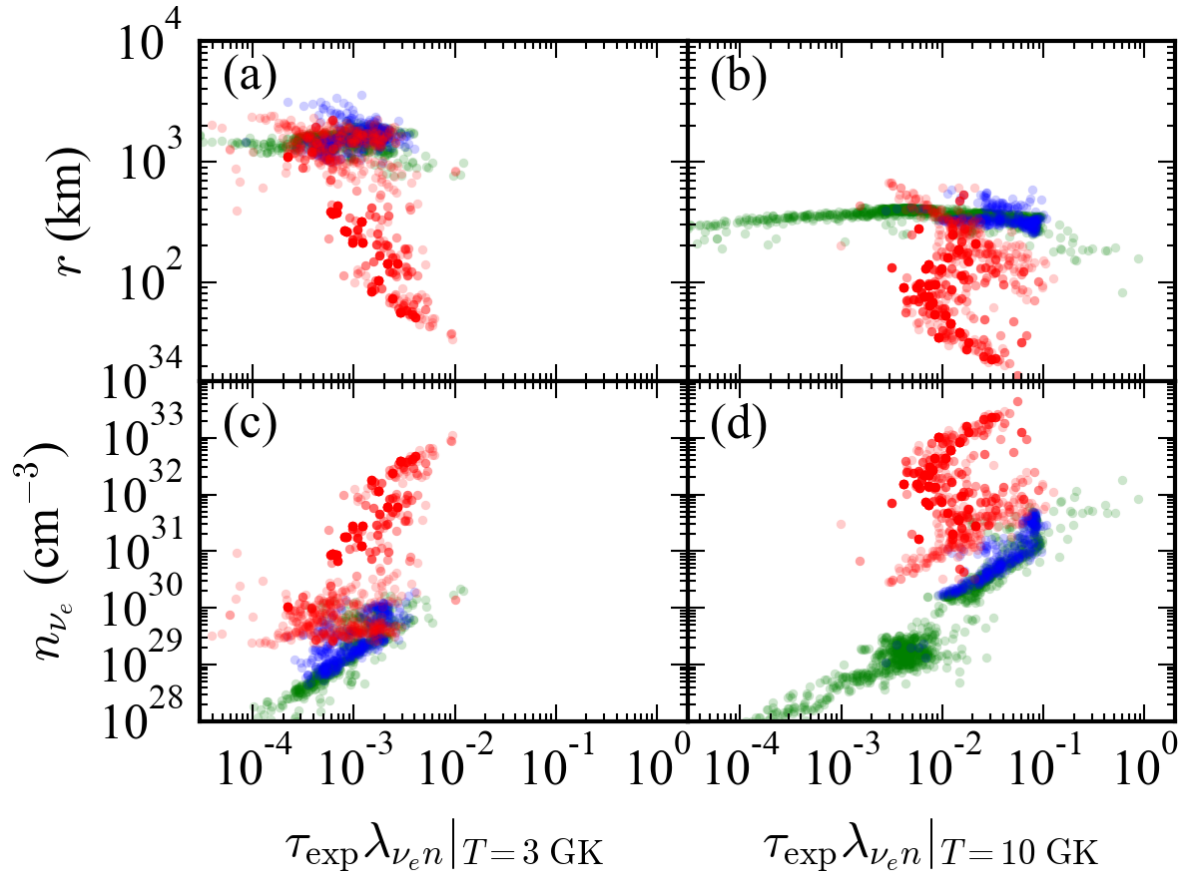
Dependence  $Y_e$  and neutrino fluence



Current neutrino-hydrodynamical models far from the necessary conditions

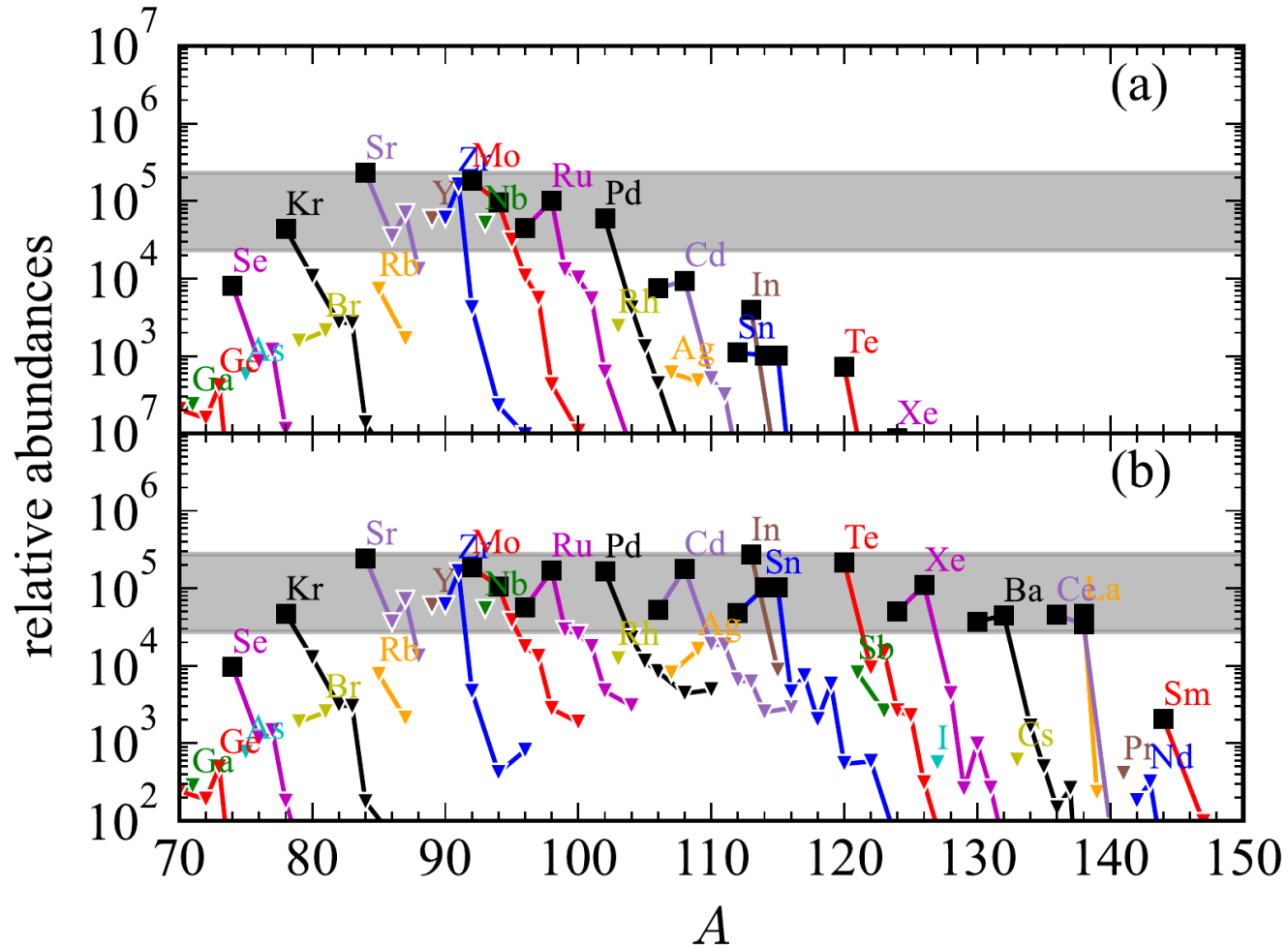


Dynamical, neutron-star torus, black-hole torus



- Material reaches the necessary fluence conditions but it is too hot form nuclei
- A non-thermal ejection mechanism is necessary (magnetic fields?)

# Coproduction of all p-nuclei



- All p-nuclei can be consistently produced
- Assuming the same astrophysical site produces both r-process and p-nuclei around 1% of the ejecta should reach  $\nu r$ -process conditions

- Multi-messenger observations (Gravitational and Electromagnetic waves) from binary neutron star mergers provide unique opportunities to study the production of heavy elements:
  - Neutron star mergers identified as one astrophysical site where the r-process operates
  - Kilonova observations provide direct evidence of the “in situ operation of the r-process”
- Strong synergies with FRIB and FAIR experiments
- $\nu r$ -process: new mechanism production p-nuclei:
  - Gamow-Teller and (spin-)dipole resonances near stability determine neutrino cross sections.
  - Important role of  $(n, \gamma)$  cross sections near stability.