Experiments with exotic nuclei II

Monday

~~Tuesday~~ *Monday*

Excited states

Experimental considerations: Reactions

Collectivity

Single-particle degrees of freedom

Transfer and knockout (only brief: See A. Wuosmaa lectures)

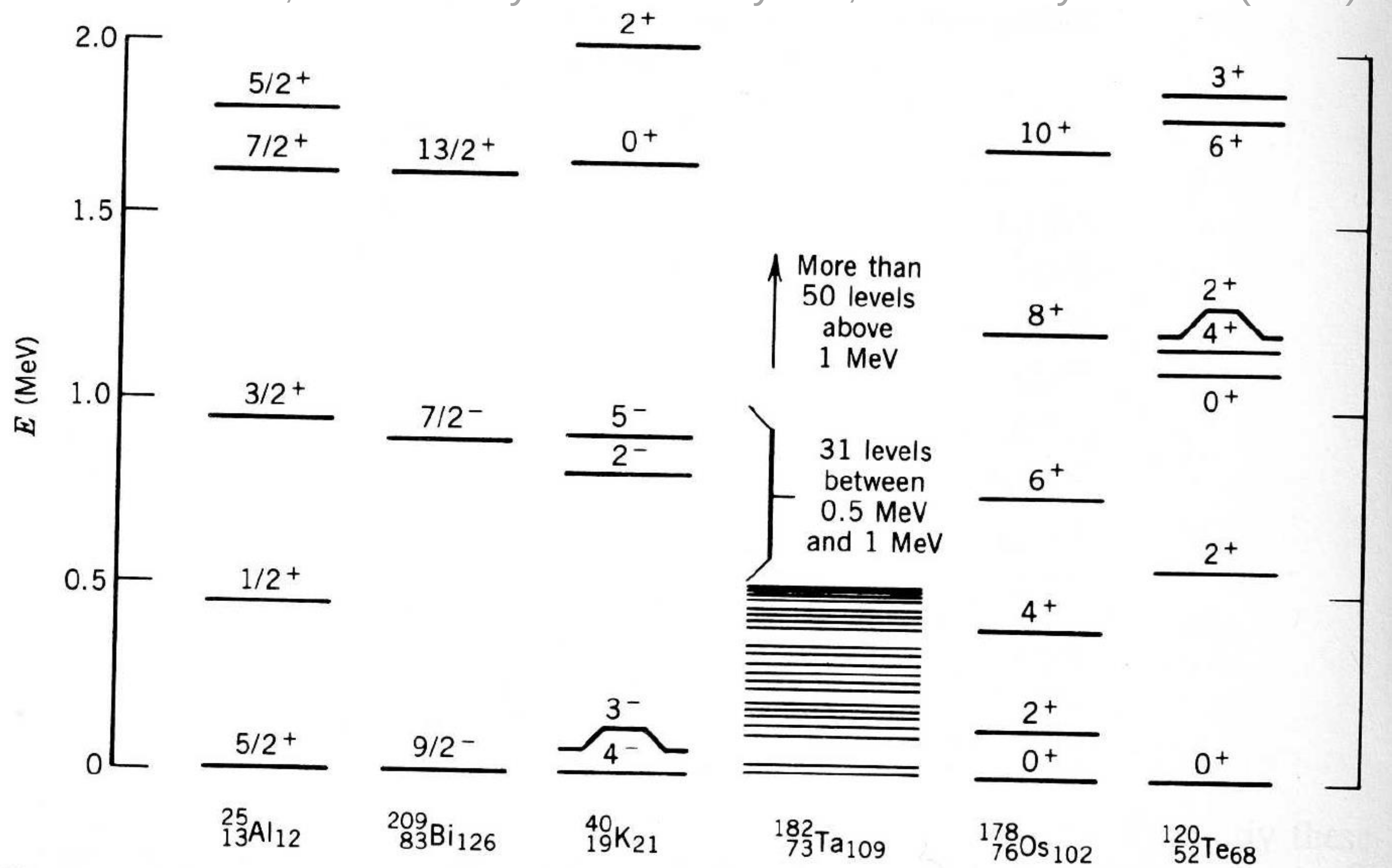
Excited-state lifetimes

Excited states

K. S. Krane, Introductory Nuclear Physics, John Wiley & Sons (1988)

Collective excitation: all nucleons outside a closed shell contribute coherently to the excitation (vibration, rotation)

Single-particle excitation: Excited states are formed by rearranging one or a few nucleons in their orbits



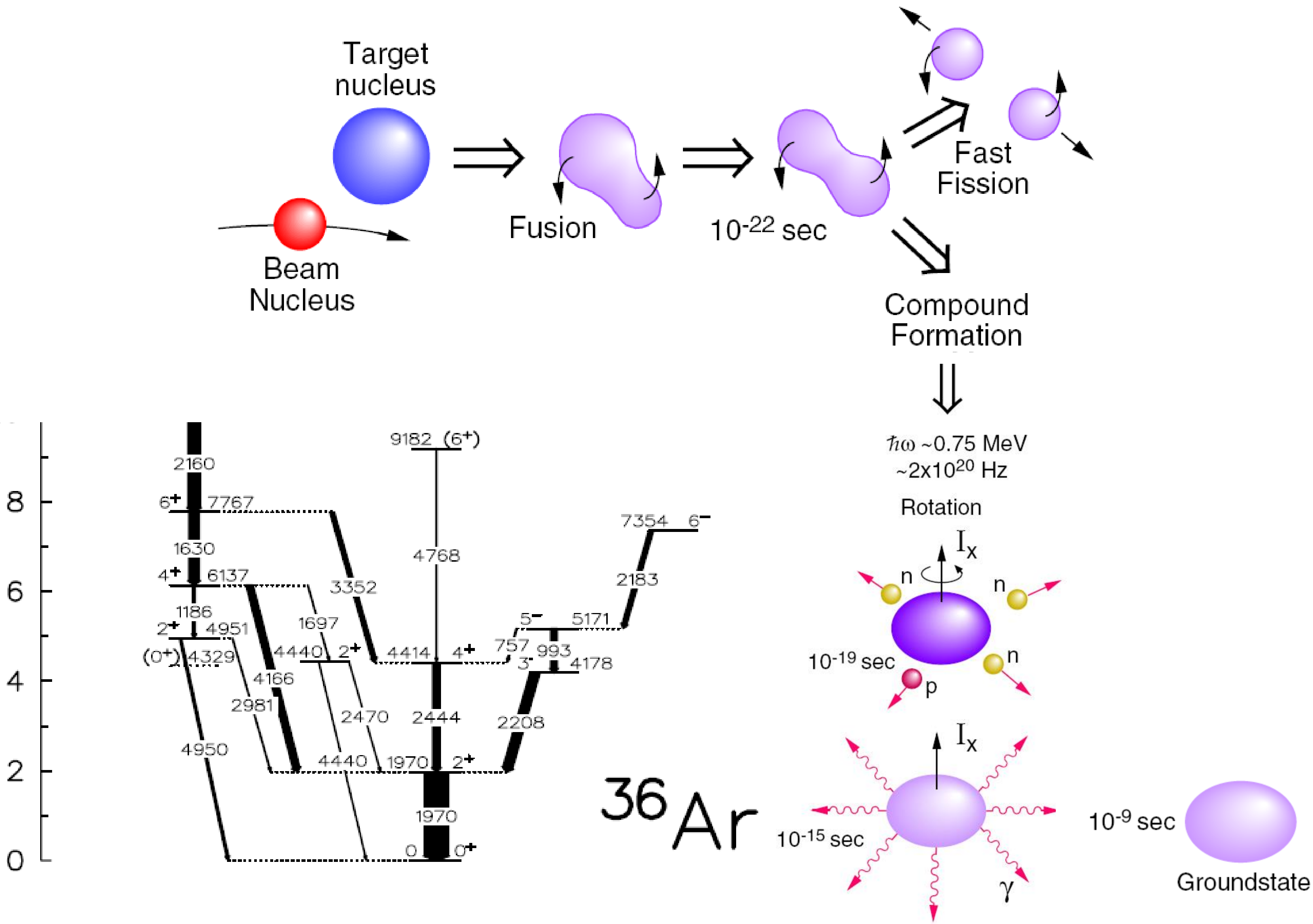
- In nuclei, the energy scales are close:

$E_{rot} \sim E_{vib} \sim E_{sp}$ (MeV)

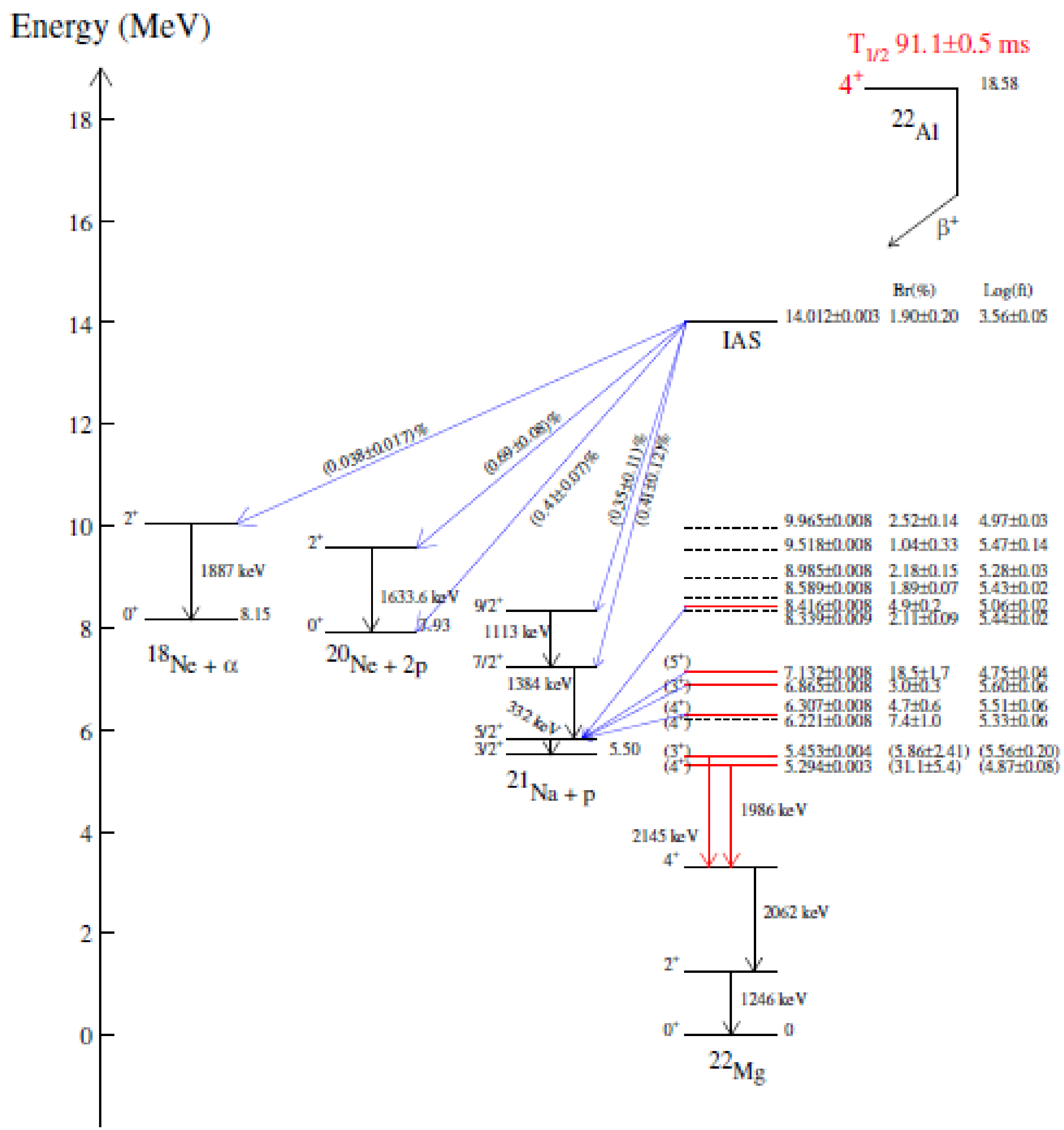
In your hands-on measurement, you will be able to study ^{24}Mg, a well-deformed rotor

Collective and single-particle excitation can be separated but interact strongly

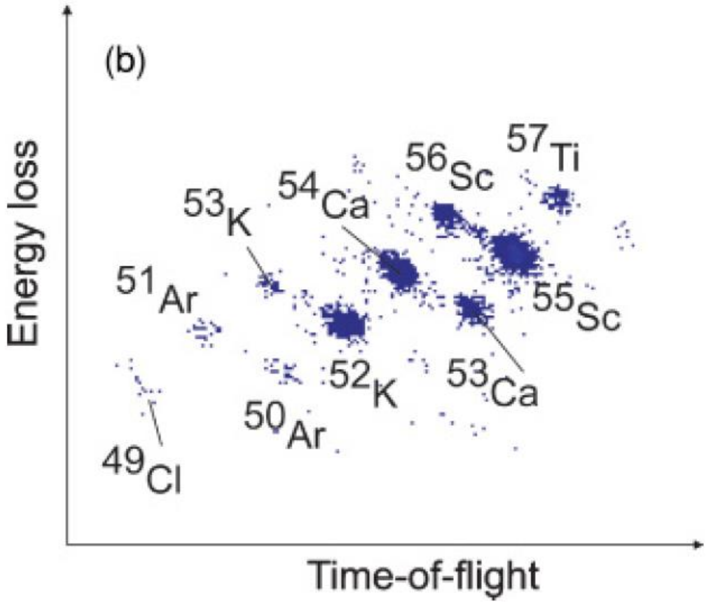
Population of excited states - Reactions



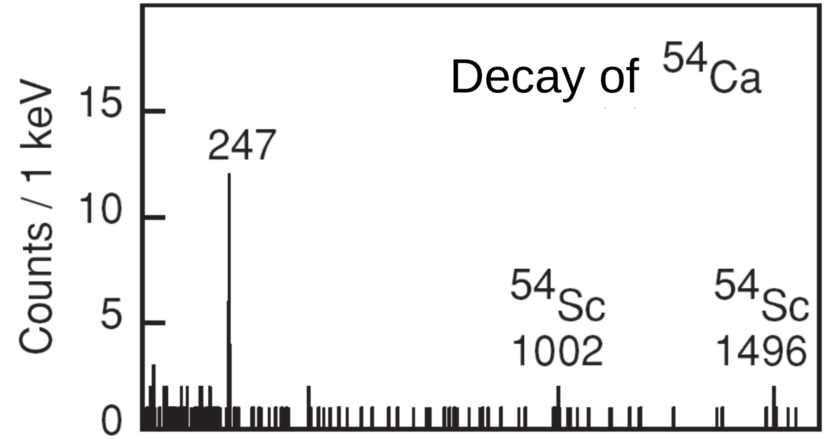
Population of excited states - Decays

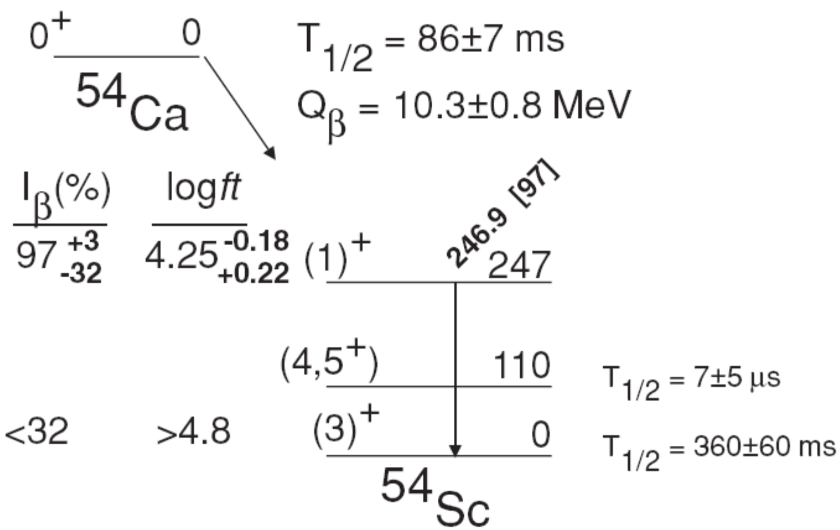


Excited states populated in β decay
Selectivity through selection rules



Total number of ^{54}Ca implants: 654 only



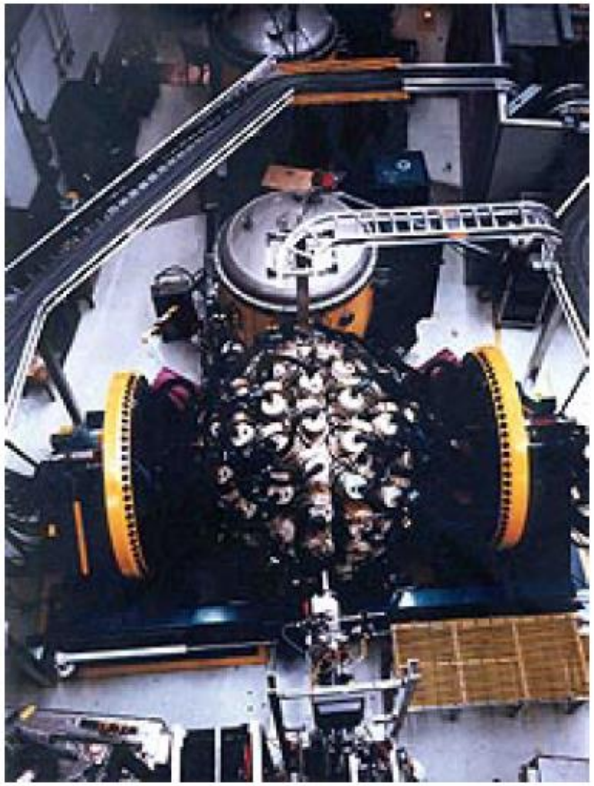


Selection rules in β decay, any textbook

Type	log(*ft*)	*L*	$\Delta\pi$	$\Delta\vec{J}$ $\vec{S}=\vec{0}$ Fermi	$\Delta\vec{J}$ $\vec{S}=\vec{1}$ Gam-Tel
super-allowed	2.9-3.7	0	+	0	0
allowed	4.4-6.0	0	+	0	0,1
first forbidden	6-10	1	-	0,1	0,1,2
second forbidden	10-13	2	+	1,2	1,2,3
third forbidden	>15	3	-	2,3	2,3,4

P. F. Mantica et al., PRC 77, 014313 (2008)

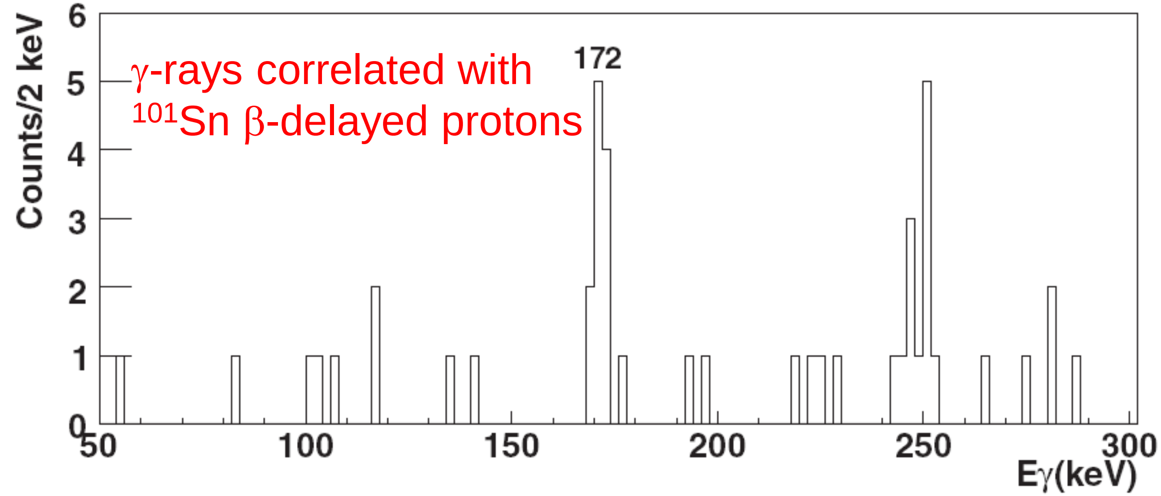
γ-ray spectroscopy tagged with β-delayed protons

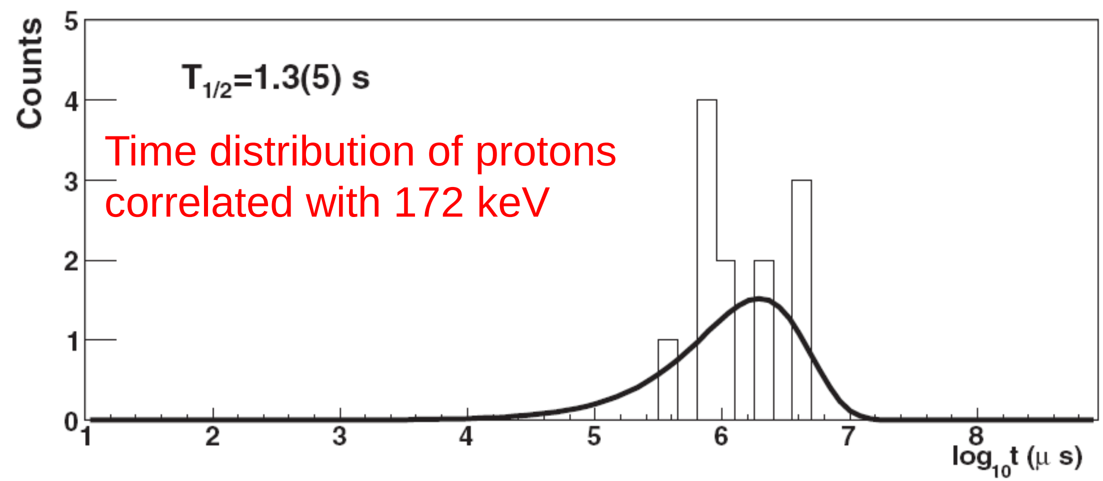


Single-neutron states above doubly magic ^{100}Sn:

$d_{5/2} - g_{7/2} \sim 172$ keV

$^{58}Ni + ^{46}Ti$ at 192 MeV (ATLAS/ANL)

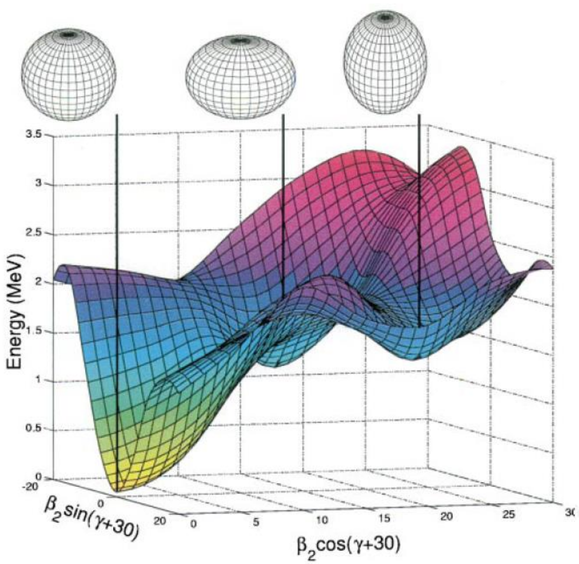


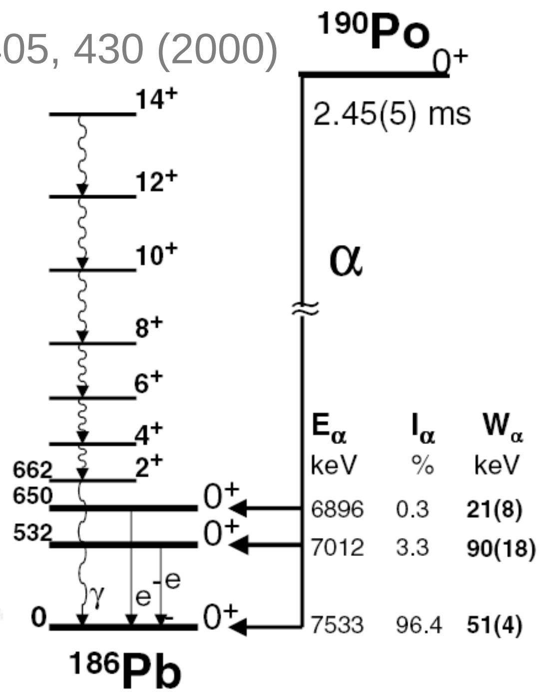


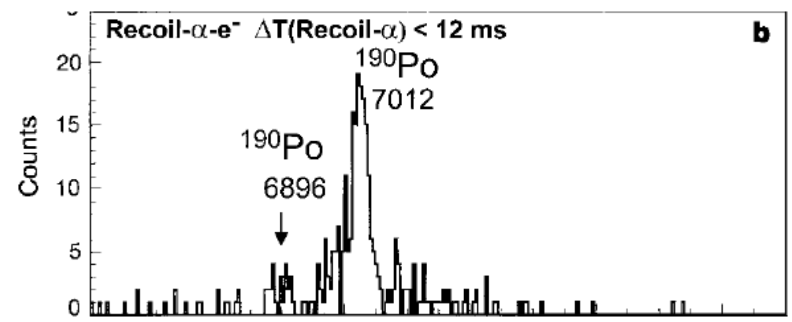
D. Seweryniak et al., PRL 101, 022504 (2007)

Excited states populated following α and proton emission

A. N. Andreyev et al., Nature 405, 430 (2000)

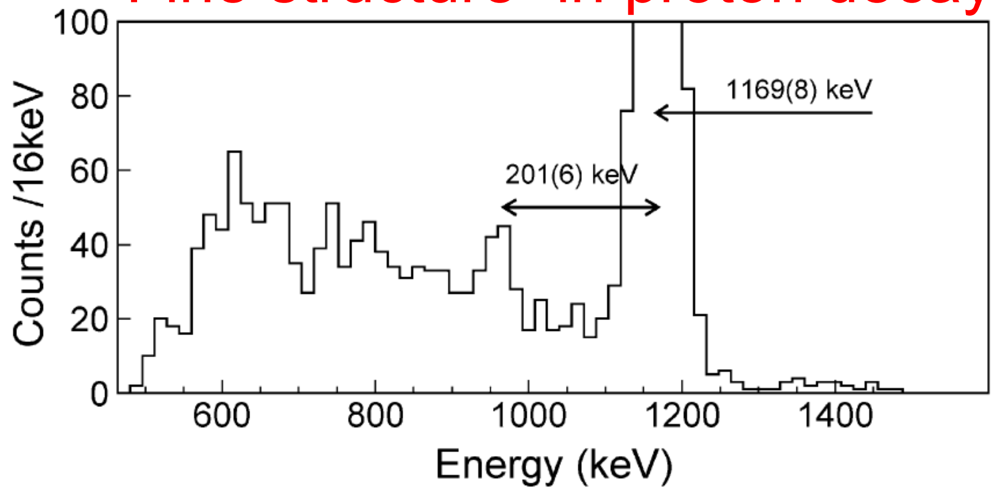






Ground state and first excited state (201 keV) of ^{140}Dy populated in proton decay of ^{141}Ho

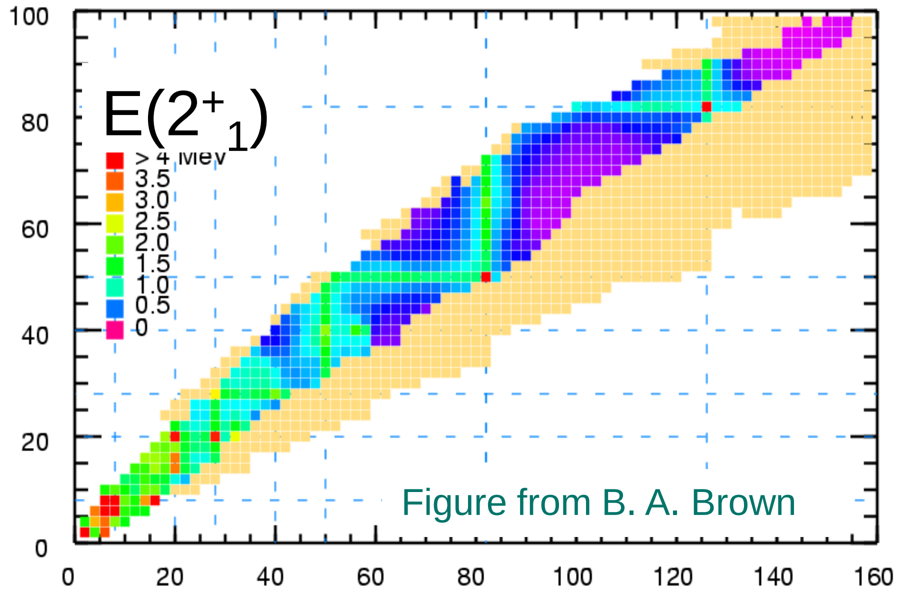
“Fine structure” in proton decay



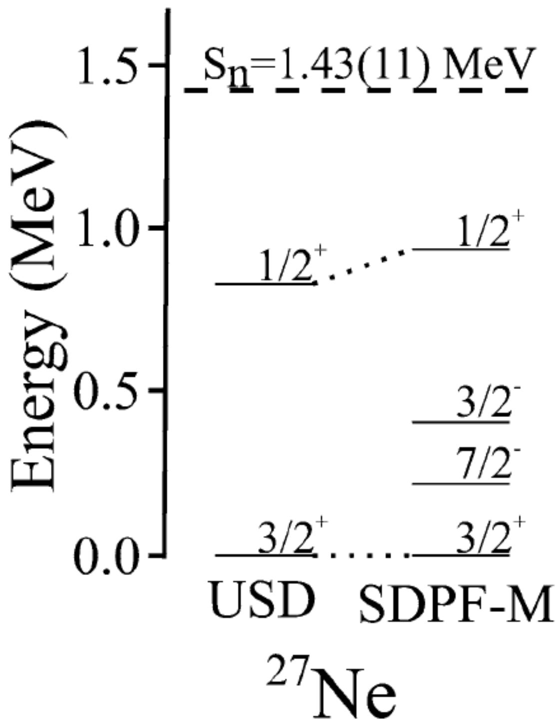
M. Karny et al., PLB 664, 52 (2008)

Structure information from excited states

As one indicator of shell closures



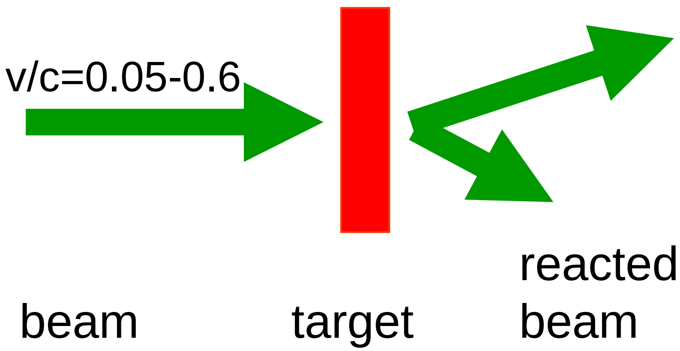
Guide model calculations



J. R. Terry et al., Phys. Lett. B 640, 86 (2006)

Experimental considerations: *Reactions*

Nuclear reactions – cross section



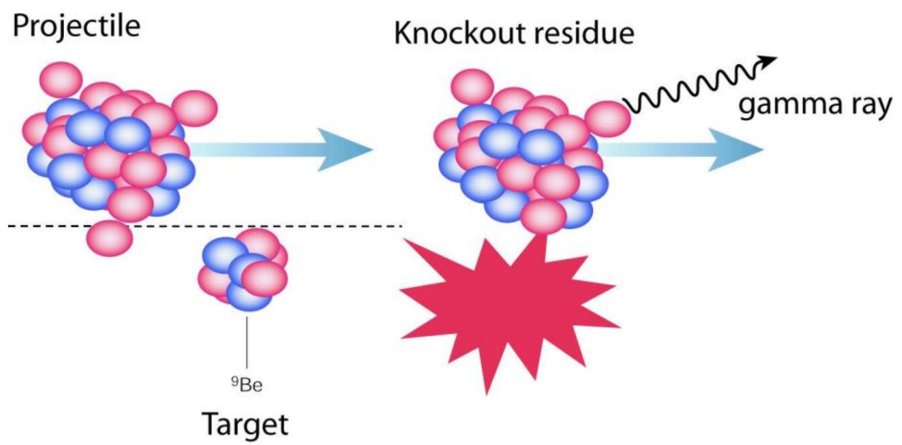
- The choice of the target depends on the reaction hat is desired

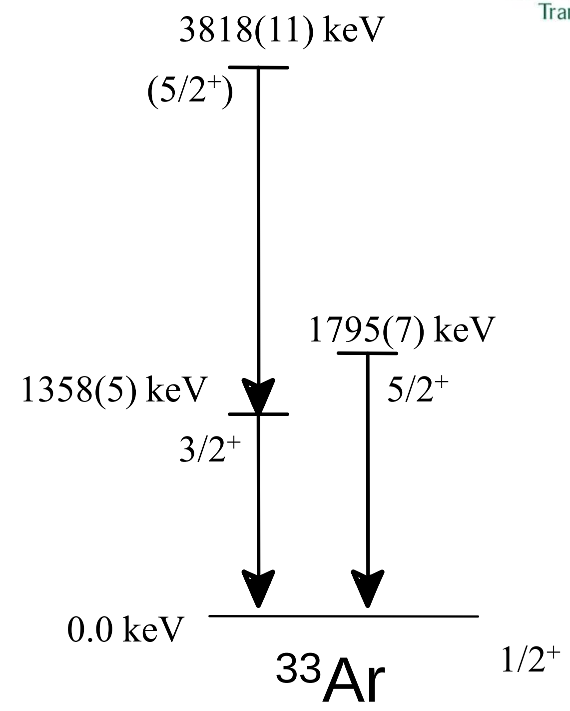
- $N_R = \sigma \times N_T \times N_B$
  - σ Cross section
  - N_T Atoms in target
  - N_B Beam rate
  - N_R Reaction rate

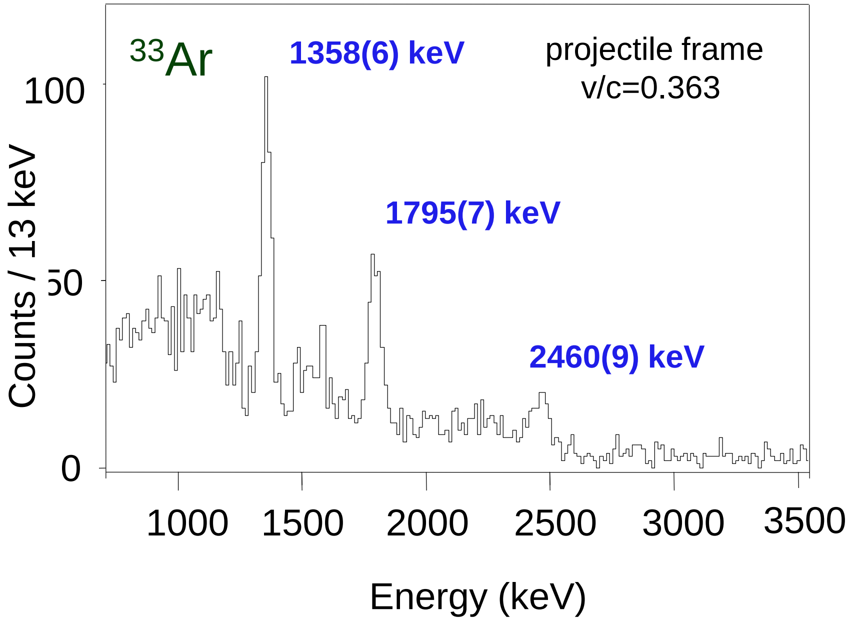
- Reactions
  - Inelastic scattering
  - Nucleon transfer
  - Fusion, fusion-evaporation
  - Breakup/fragmentation

- Experimental task
  - Identify and count incoming beam
  - Identify and count reacted beam
  - Tag the final state of the reaction residue
  - Measure scattering angles and momentum distributions

Excited states

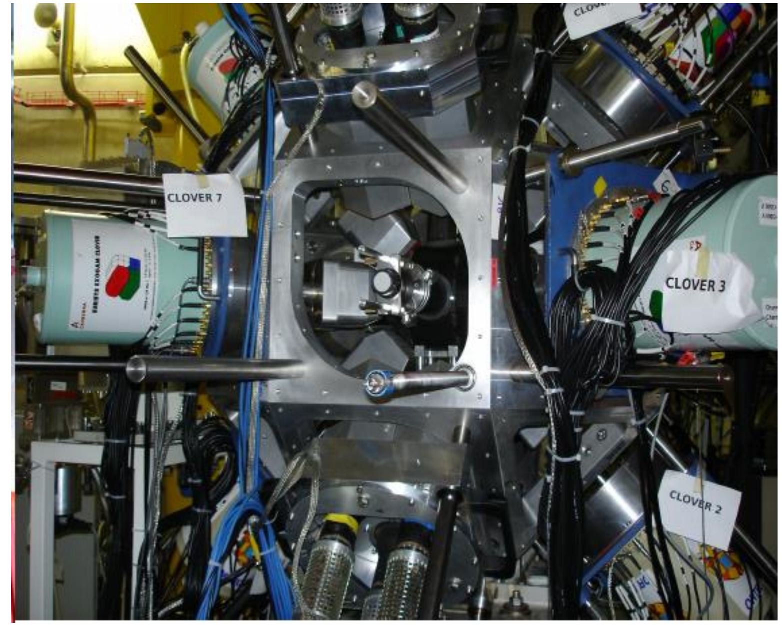


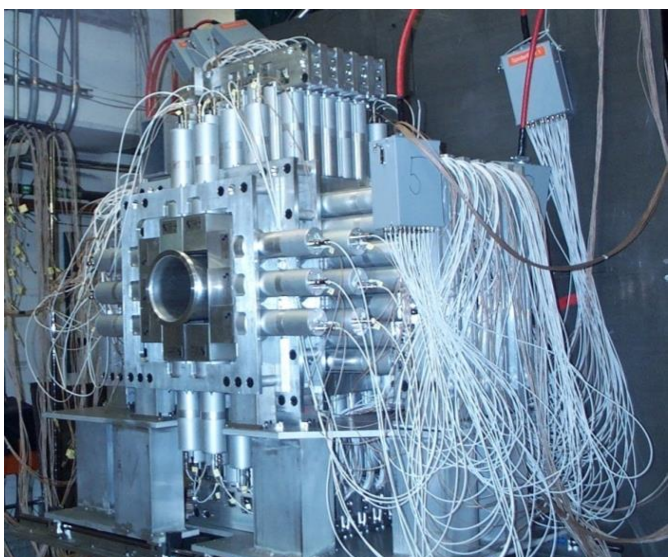




- Tag the population of excited states by measuring the decay γ rays. The γ-ray energy gives the energy difference between two states.

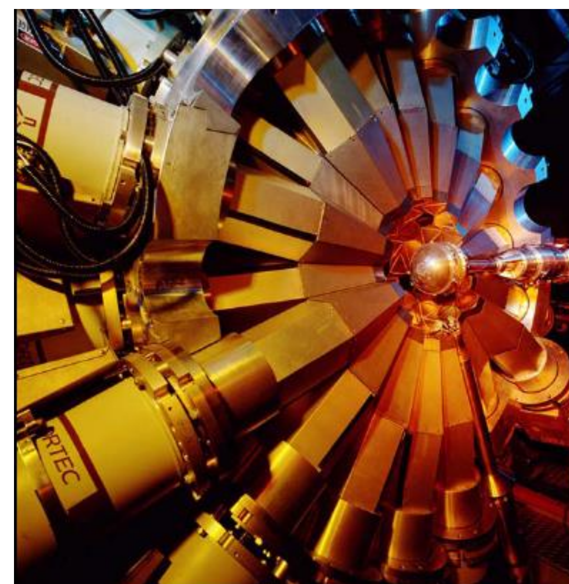
Gamma-rays to tag the final state





Germanium detectors:
Superior energy resolution, but modest efficiency

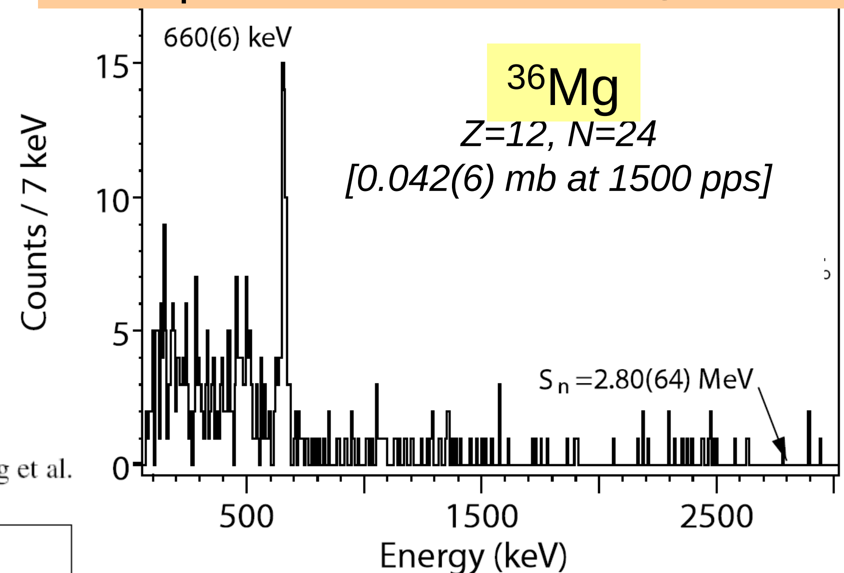
Scintillator-based:
High-efficiency, moderate resolution



Gamma-rays to tag the final state

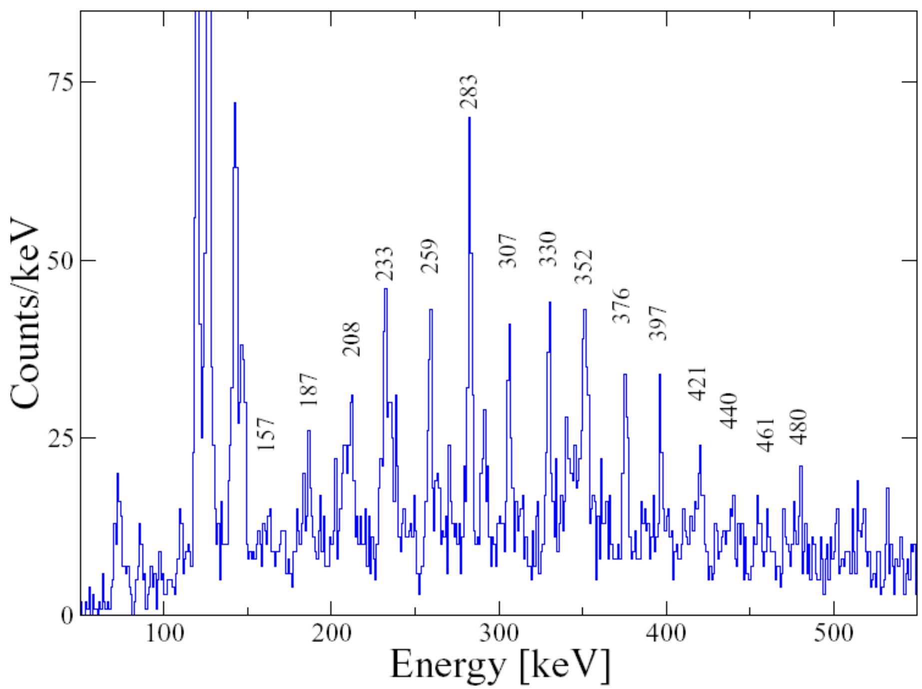
Two-proton knockout to ^{36}Mg. Only the first excited state was observed.

^{38}Si-2p at 83 MeV/u, SeGA @ NSCL



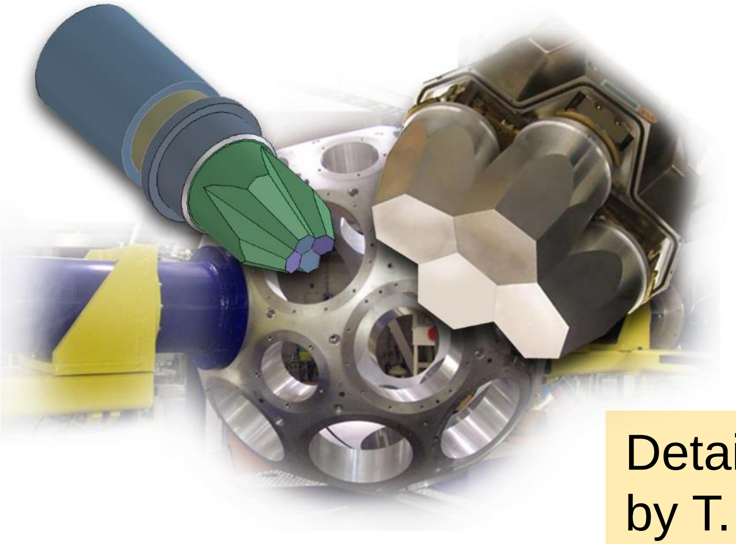
A. Gade et al., PRL 99, 072502 (2007)

^{48}Ca + ^{207}Pb ⇒ ^{253}No + 2n, JUROGAM+RITU+GREAT, R.-D. Herzberg et al.



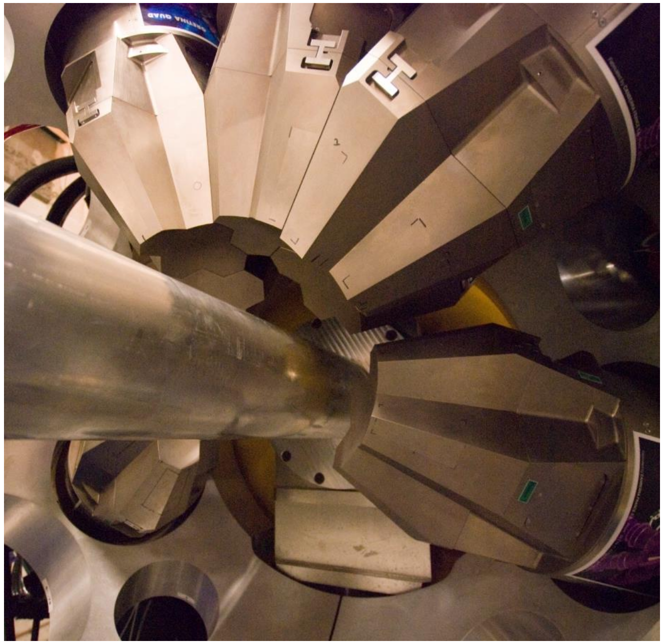
Low-energy fusion-evaporation reaction to produce ^{253}No. Many excited states are populated.

You will be doing an experiment with GRETINA



Lead institution: Lawrence Berkeley National Lab

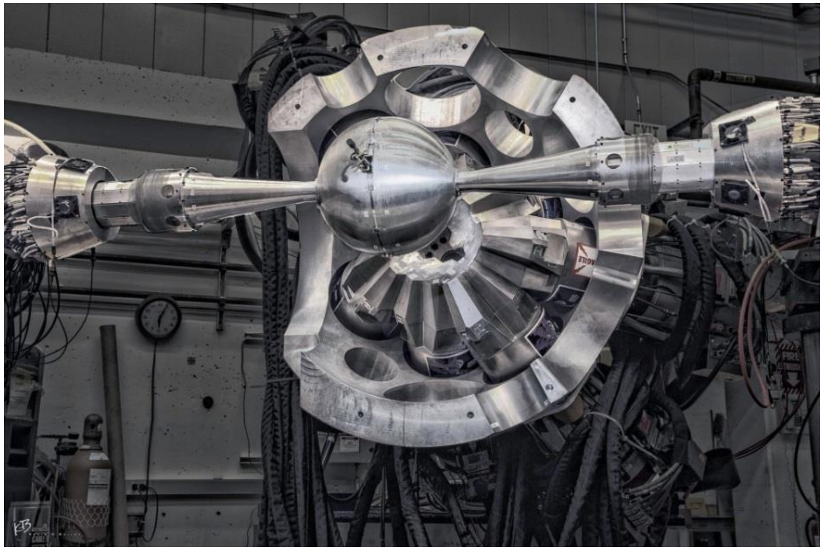
Array hosted by ANL and NSCL for experimental campaigns



Details in the lecture by T. Lauritsen

GRETINA is the next-generation Germanium-based gamma-ray tracking array:

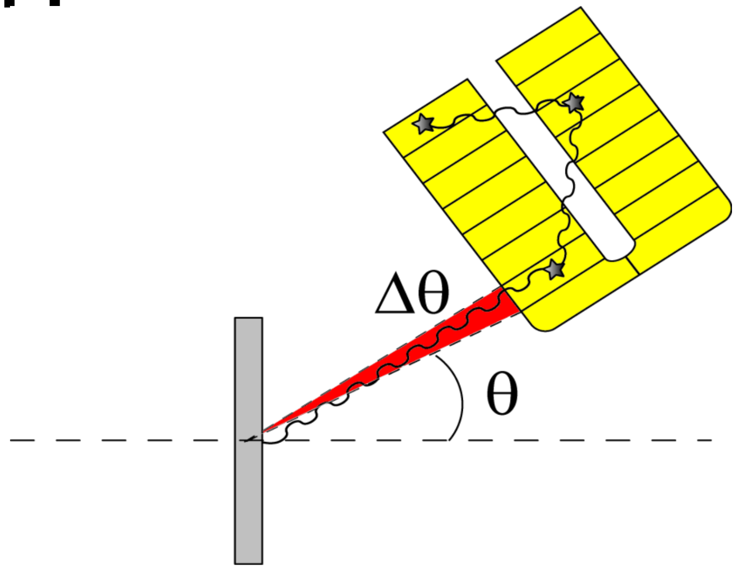
- By now: 9 modules with 4 crystals that are 36-fold segmented each (>1300 electronics channels)
- High segmentation and sub-segment position resolution through tracking of gamma-ray interactions inside the crystal



Why segmentation: Emission in flight → Doppler shift!

$$E = E_0 \frac{\sqrt{1-\beta_0^{\,2}}}{1-\boldsymbol{\beta}_0 \cdot \mathbf{e}}$$

$$\boldsymbol{\beta_0} \cdot \mathbf{e} = |\boldsymbol{\beta_0}| \cos \theta_0$$



E_0 γ-ray energy in the source frame

E γ-ray energy in the lab frame

$\boldsymbol{\beta_0}$ velocity of the source

θ_0 γ-ray angle of emission

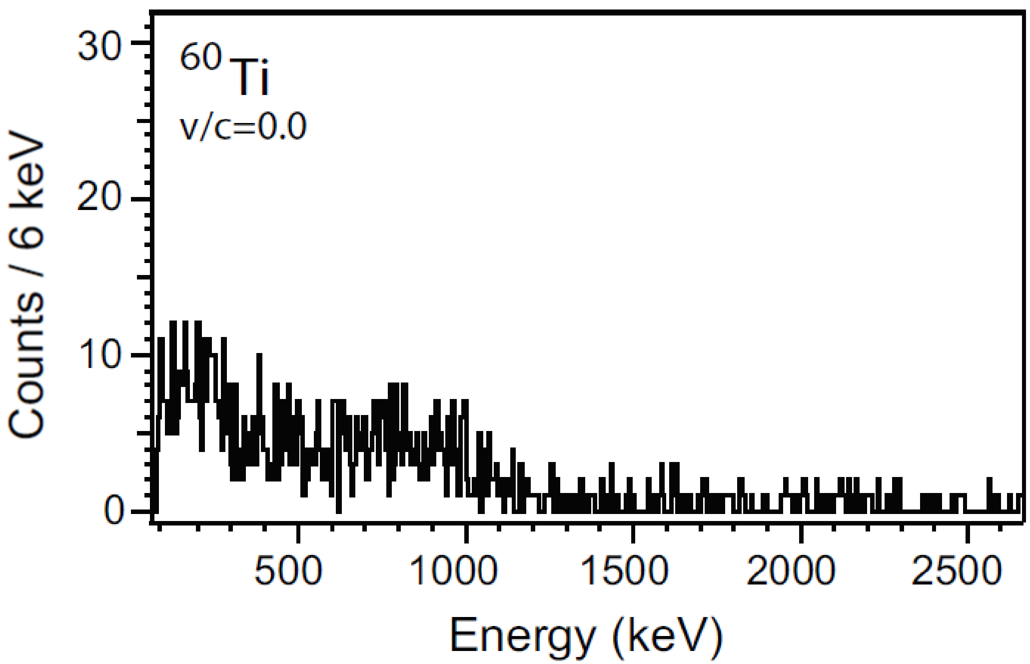
Example: SeGA geometry (NSCL)

Details in the lecture by T. Lauritsen

Gamma-rays to tag the final state

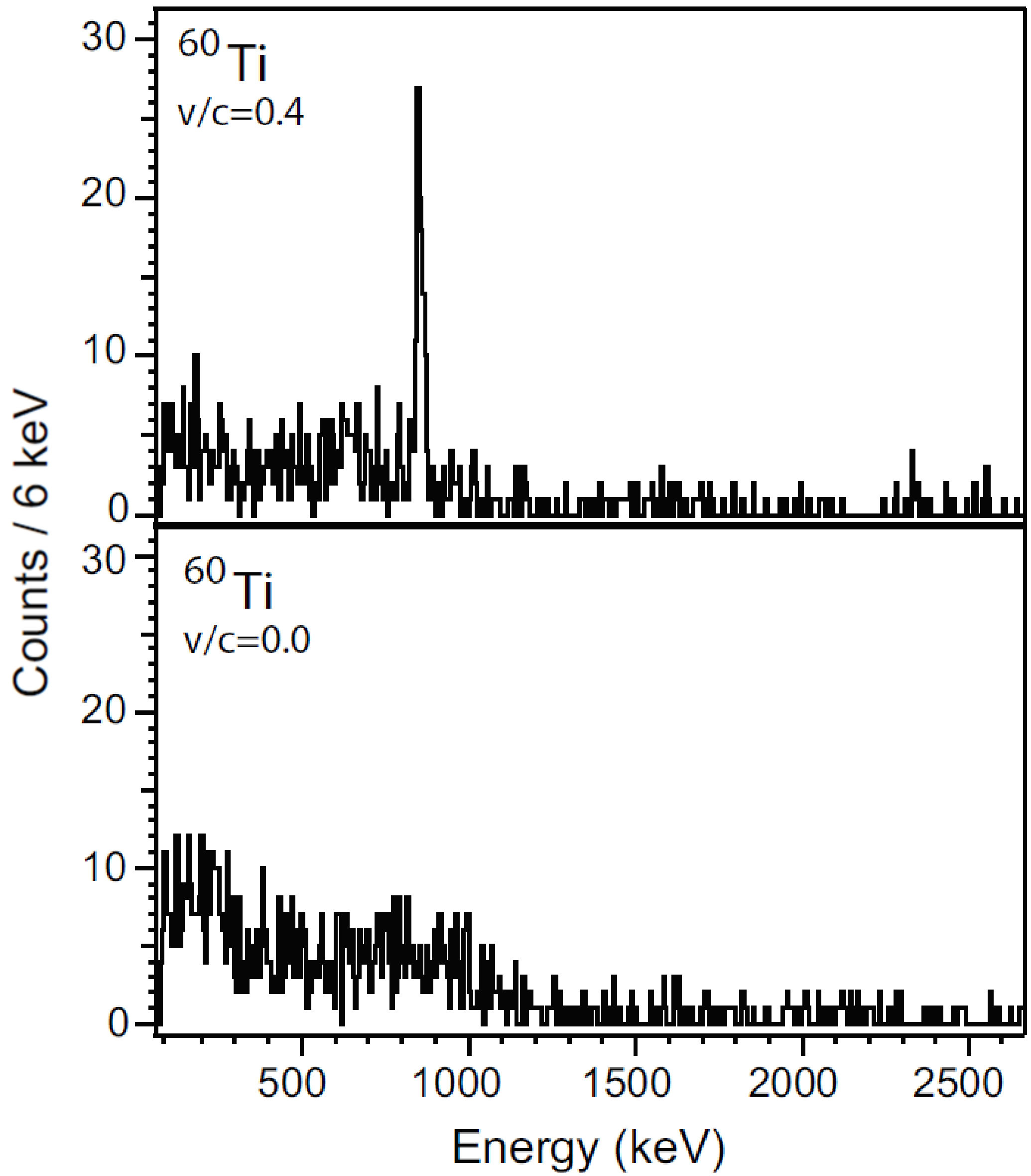
A. Gade, Eur. Phys. J. A 51, 118 (2015) - review

Making a histogram of energies as measured in the laboratory reference frame



Gamma-rays to tag the final state

A. Gade, Eur. Phys. J. A 51, 118 (2015) - review

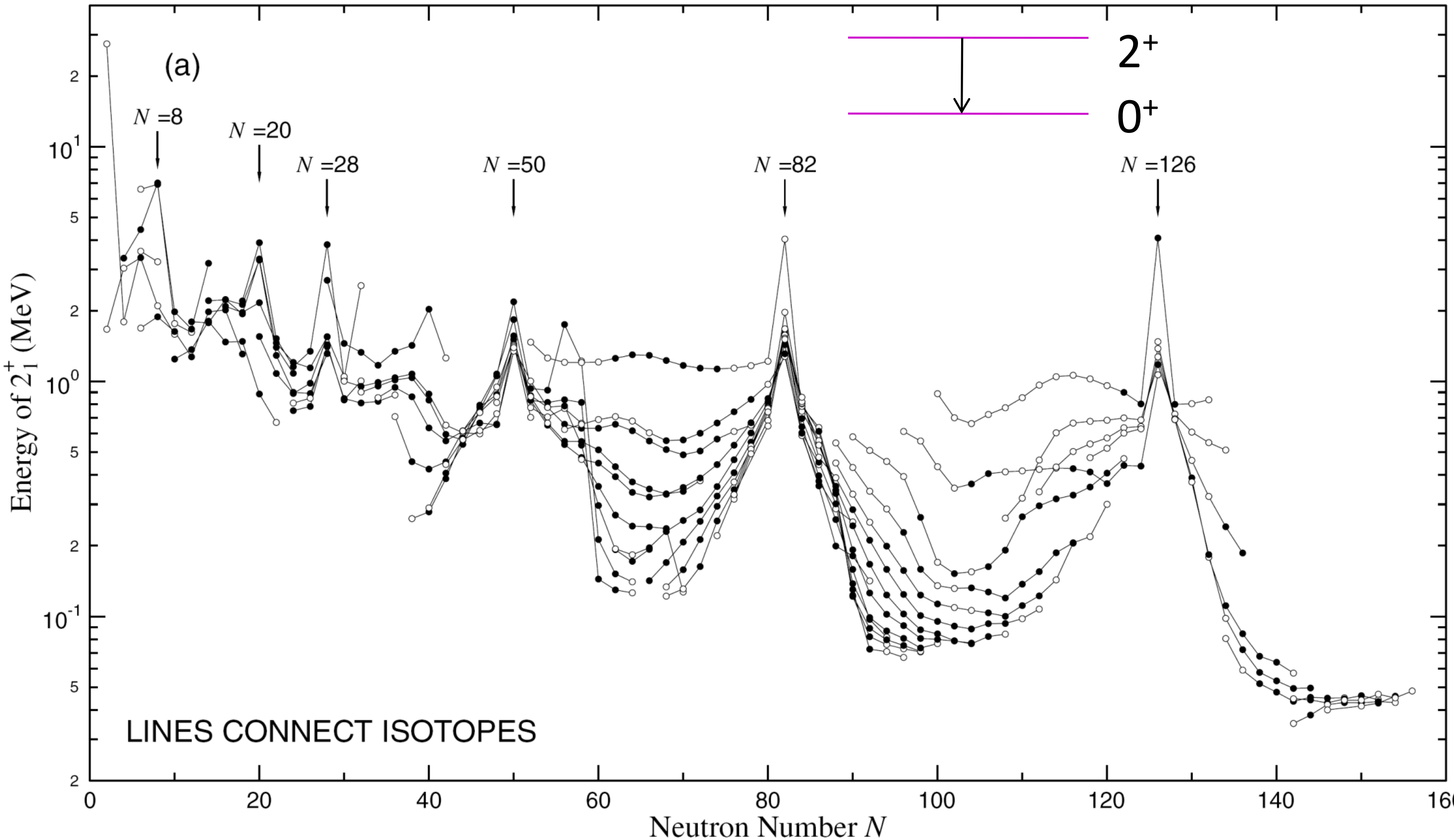


Event-by-event Doppler reconstructed using the rare isotope's velocity and the emission angle of the γ-ray

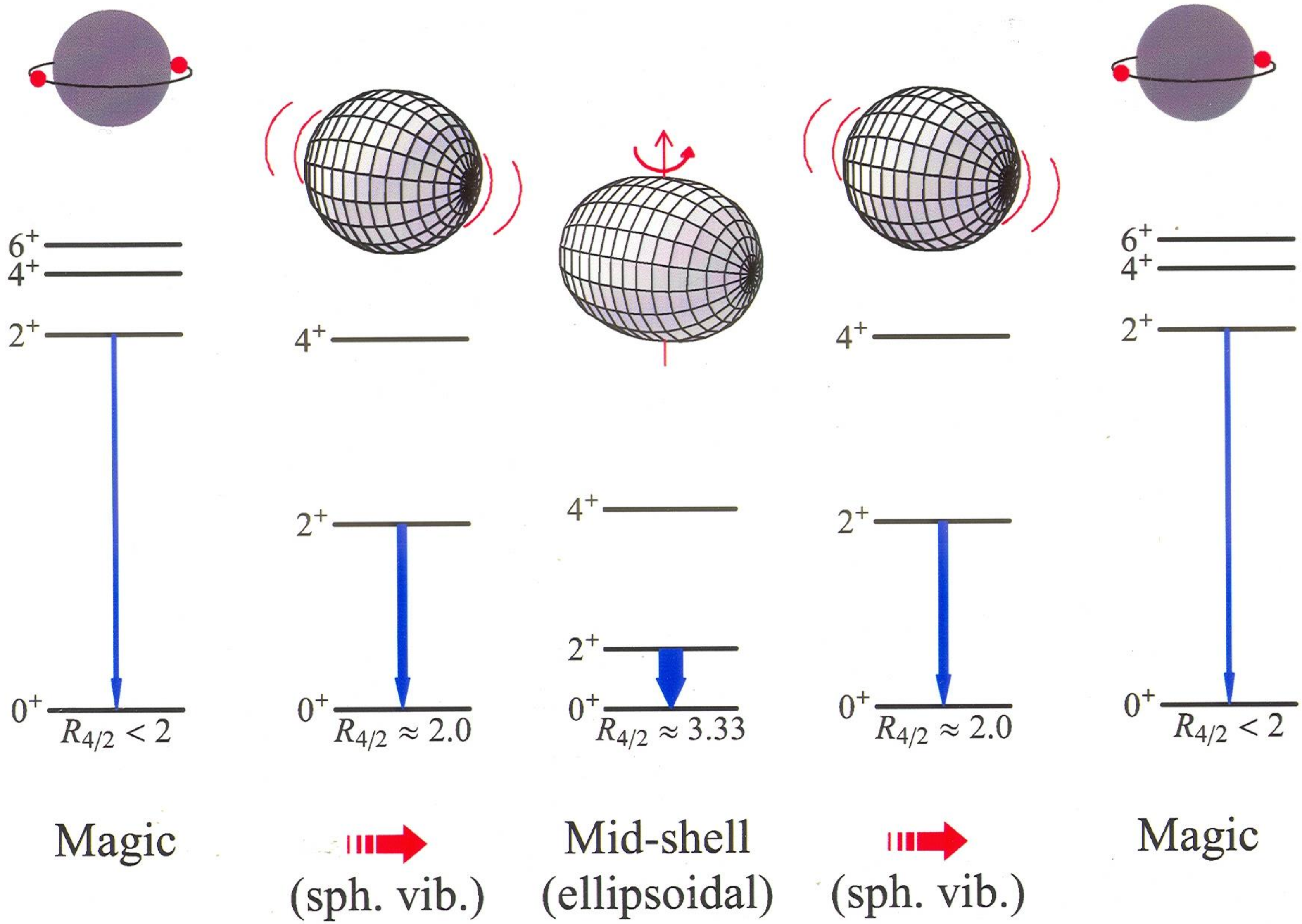
Collective excitations

Even-even nuclei: 2^+_1 state energy as an indicator of shell structure

S. Raman et al., Atomic Data & Nuclear Data Tables 78, 1



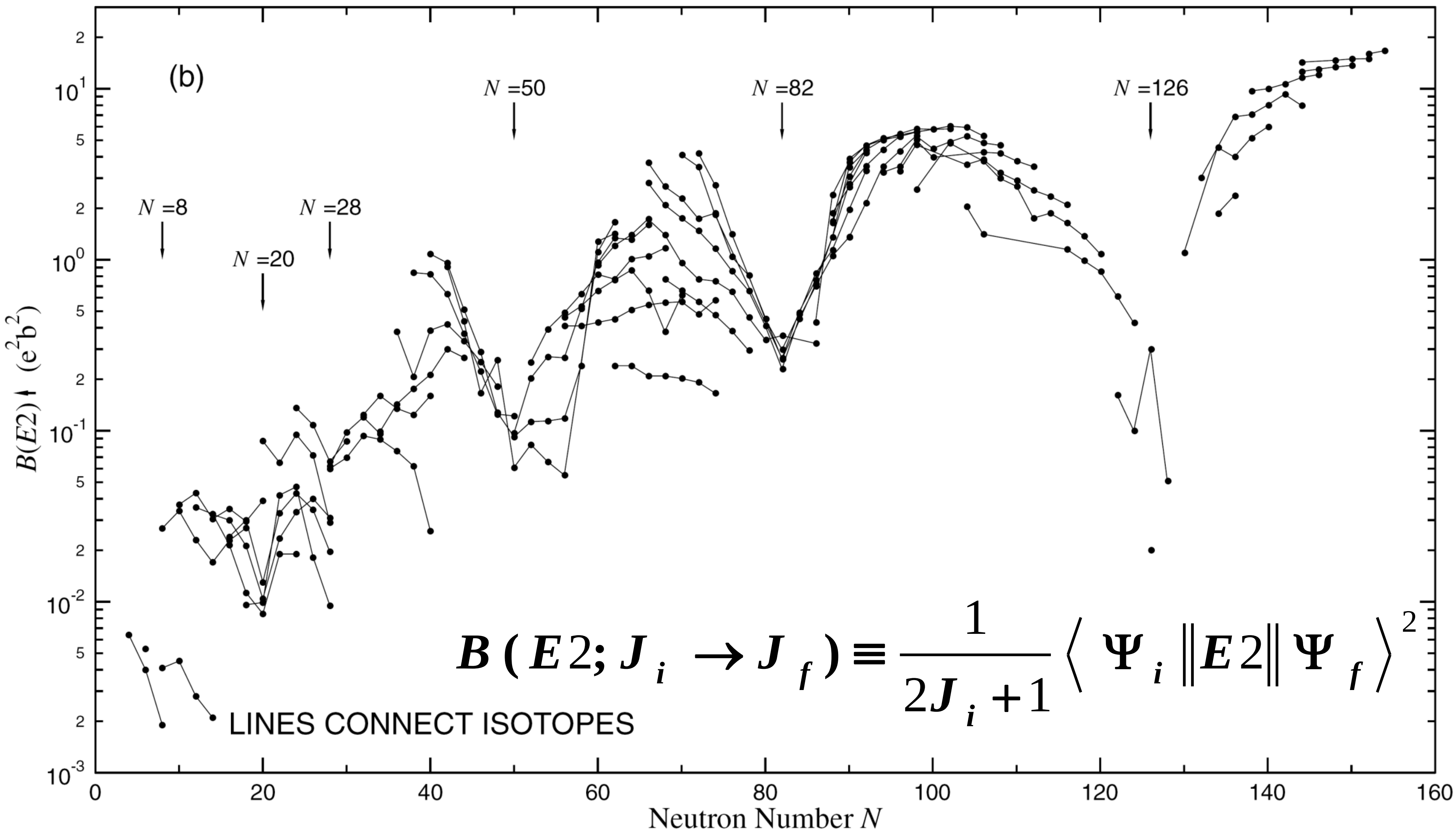
Even-even nuclei: 2^+ states are typically the first excited state on top of 0^+ ground states



Adapted from Rick Casten

Even-even nuclei: 2^+_1 excitation strength as an indicator of shell structure

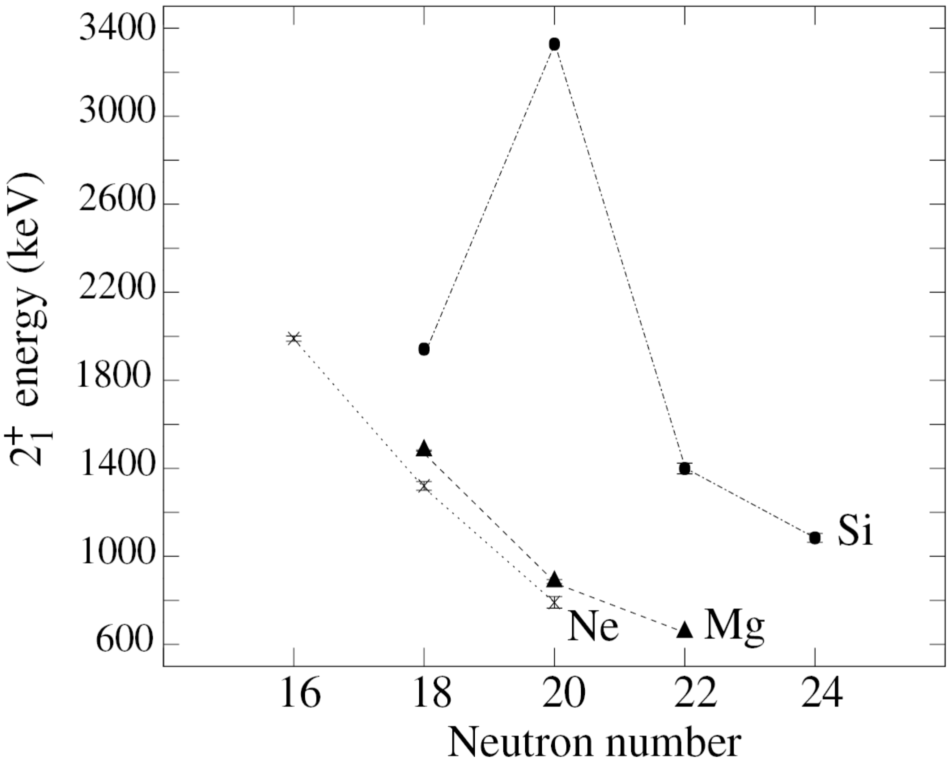
S. Raman et al., Atomic Data & Nuclear Data Tables 78, 1



$$B(E2; J_i \rightarrow J_f) \equiv \frac{1}{2J_i+1} \left\langle \Psi_i \| E2 \| \Psi_f \right\rangle^2$$

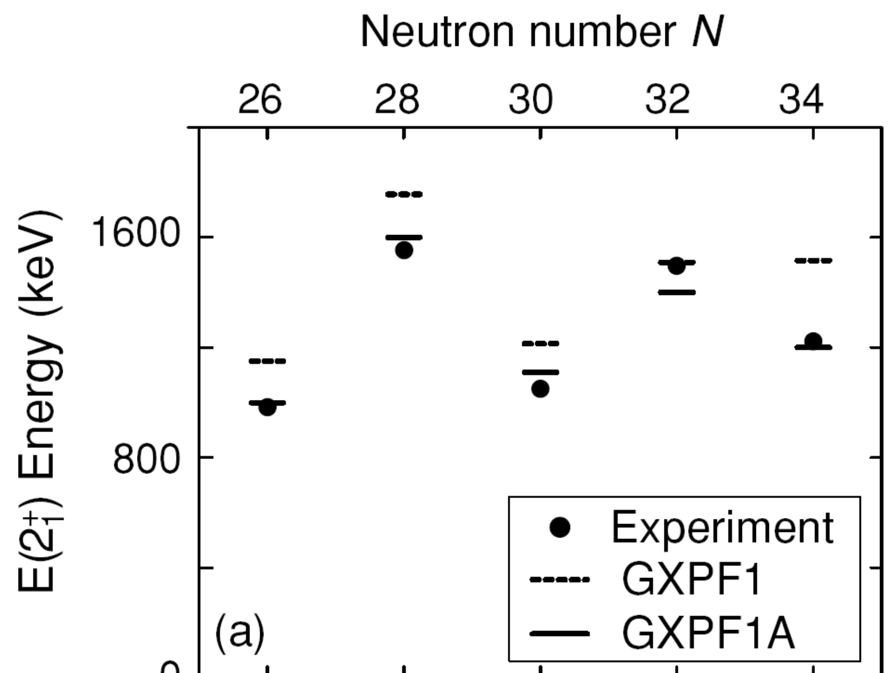
Examples of changes in shell structure

A. Gade and T. Glasmacher, Prog. In Part. and Nucl. Phys. 60, 161 (2008)



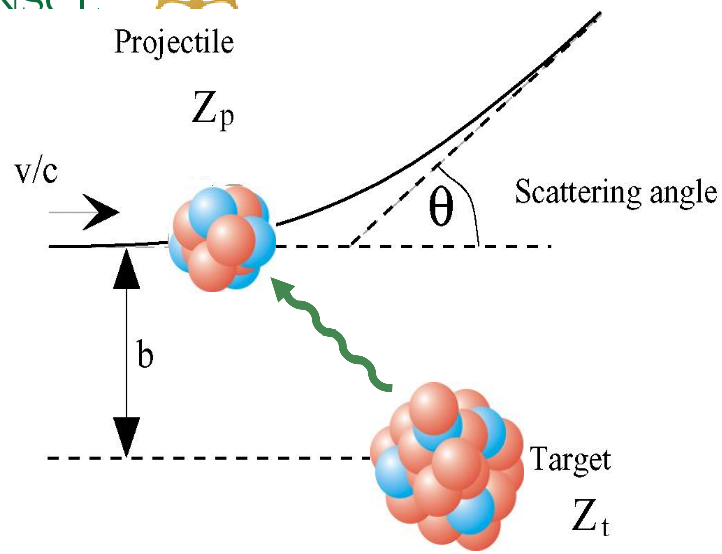
N=20 is not a good shell closure anymore in Mg and Ne isotopes

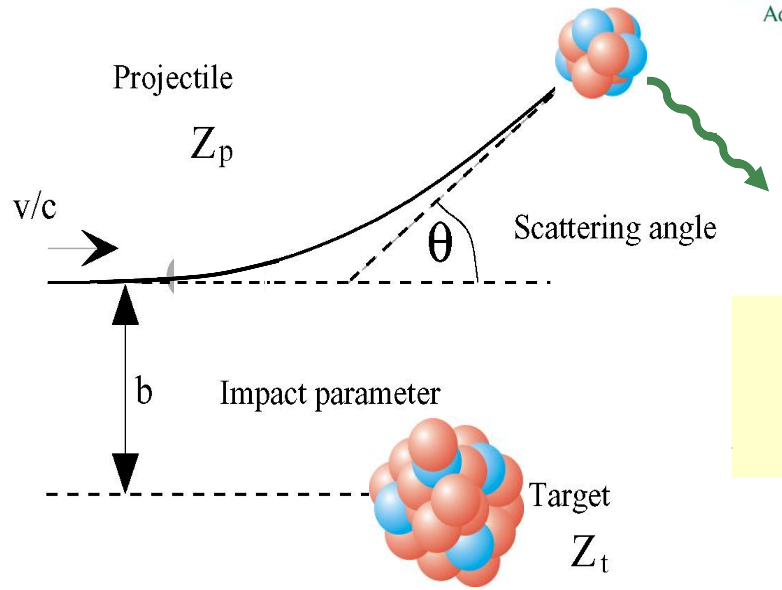
N=32 is a new magic number in the Ti isotopes

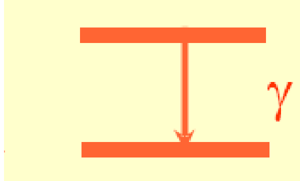


D.-C. Dinca et al., PRC 71, 041302 (2005)

Low-energy projectile Coulomb excitation







Exchange of virtual photons mediates excitation

Measure de-excitation γ-rays

Beam energies at the Coulomb barrier (SPIRAL):
E_x, B(σλ) excitation strength, band structures ($0^+ \rightarrow 2^+ \rightarrow 4^+ \rightarrow 6^+$)

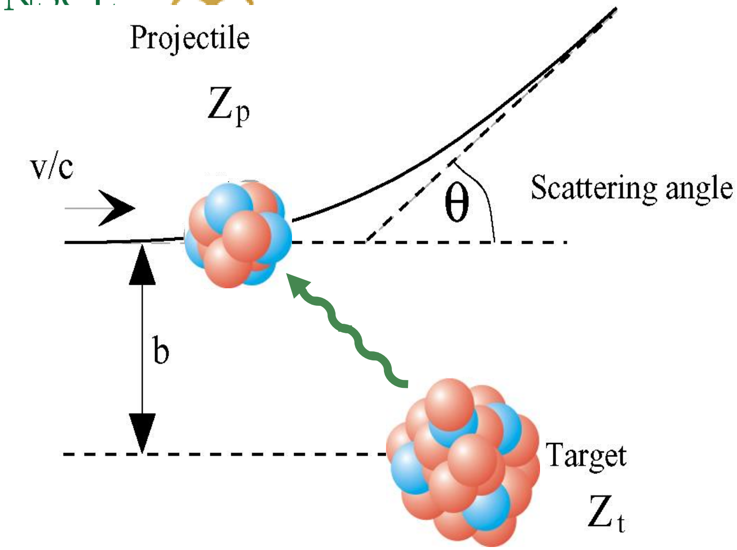
$$V_C(MeV) = \frac{1.44 \times Z_1 \times Z_2}{r(fm)}$$

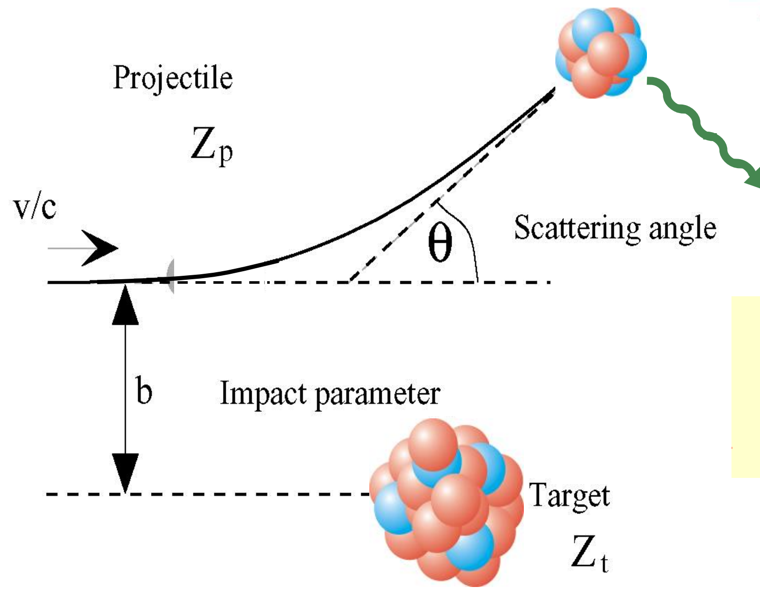
Beam energies well below the Coulomb barrier (ISOLDE, HRIBF):
Usually only the first 2^+ state accessible

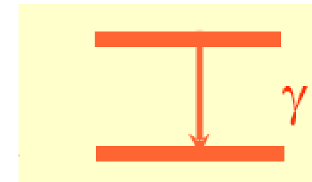
$$r(fm) \sim 1.2(A_1^{1/3} + A_2^{1/3})$$

D. Cline, Annu. Rev. Part. Sci. 36, 683 (1986)

High-energy projectile Coulomb excitation







Exchange of virtual photons mediates excitation

Measure de-excitation γ-rays

Intermediate and relativistic energies (NSCL, RIKEN, GANIL, GSI): $E(2^+_1)$, $B(E2,0^+ \rightarrow 2^+_1)$ excitation strength, two-step to 4^+ heavily suppressed (short interaction time at high beam energies)

BUT: the collision between target and projectile happens above the Coulomb barrier for every target-projectile combination

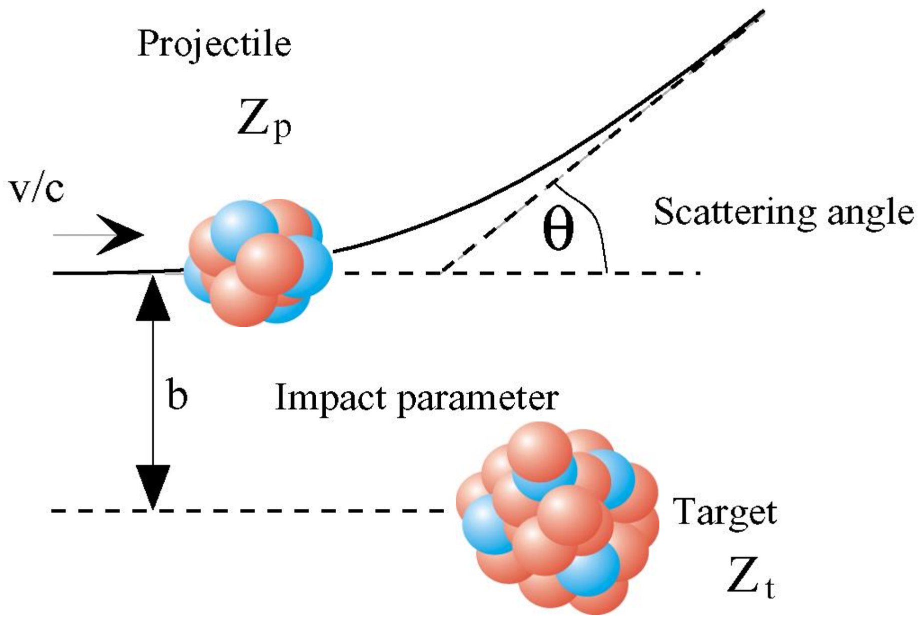
How can this still be Coulomb excitation?

T. Glasmacher, Annu. Rev. Part. Sci. 48, 1 (1998)

How can it be Coulomb excitation at energies above the Coulomb barrier ?!

At NSCL, RIKEN, GSI … the collision between target and projectile happens above the Coulomb barrier for every target-projectile combination

But: electromagnetic interaction dominates for b > R_{int}



For given v/c:

impact parameter b=b(θ)

$$b_{\min} = \frac{a}{\gamma} \cot(\theta_{\max}^{\mathrm{cm}}/2)$$

$$a = \frac{Z_p Z_t e^2}{\mu v^2}$$

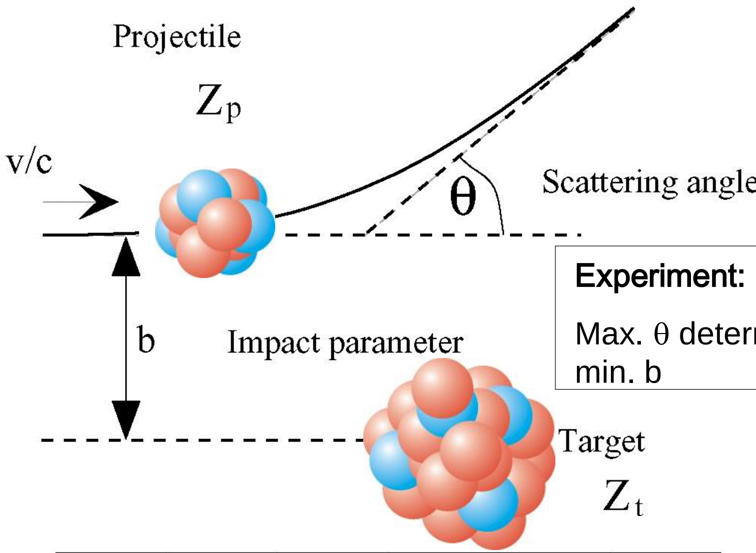
Experiment:

Maximum scattering angle determines minimum b. Restrict analysis to events at the most forward scattering angles so that b(θ) >R_{int}

T. Glasmacher, Annu. Rev. Part. Sci. 48, 1 (1998)

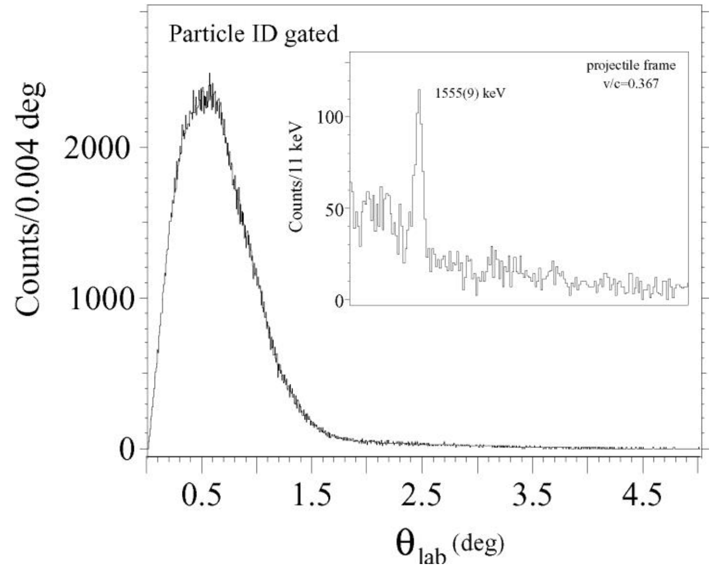
Intermediate-energy Coulomb excitation
Example: $^{46}Ar + ^{197}Au$

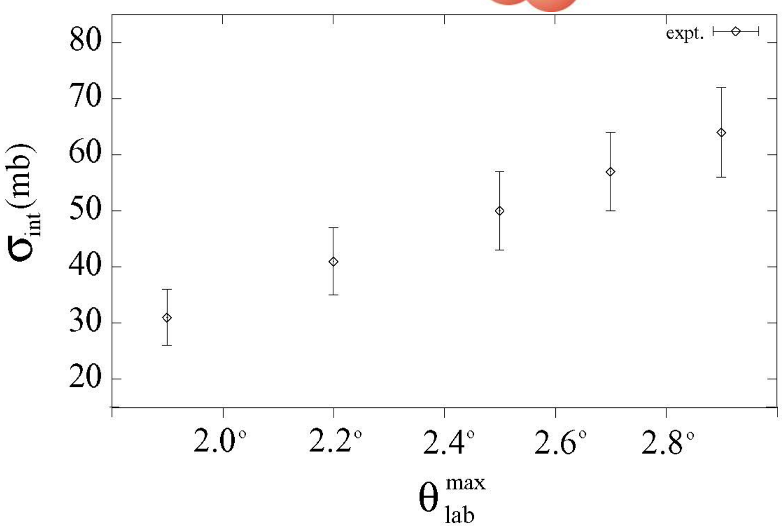
A. Gade *et al.*, PRC 68, 014302 (2003)

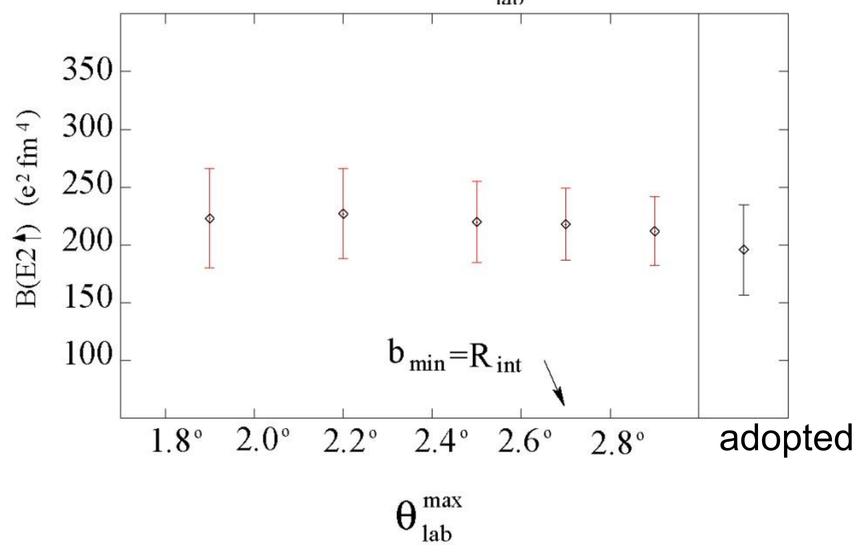


Experiment:

Max. θ determines min. b



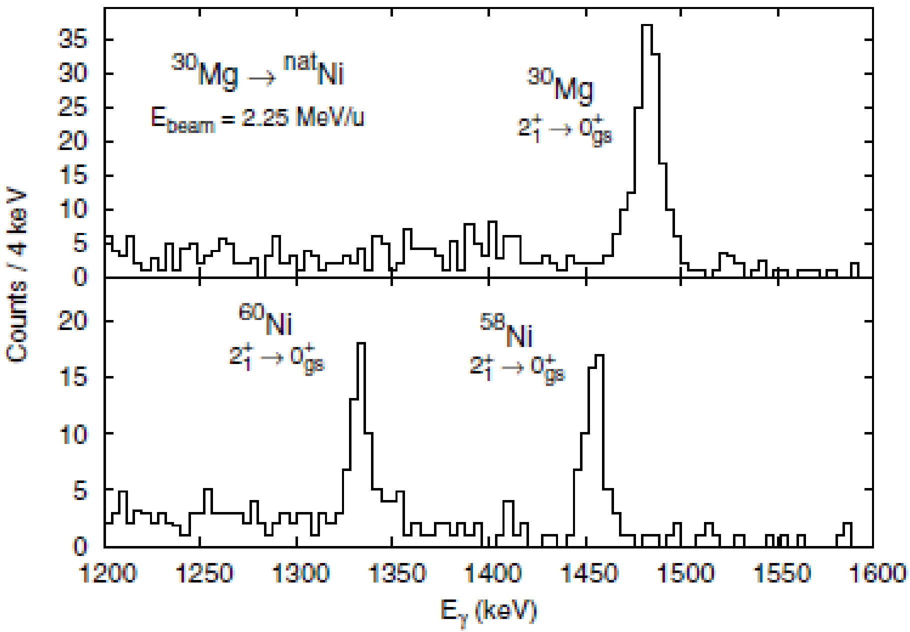


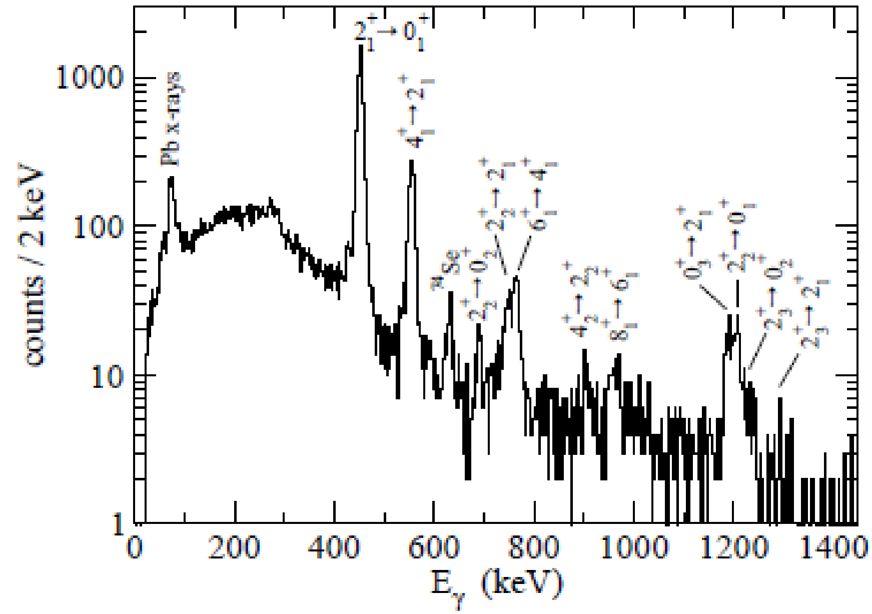


A. Winther and K. Alder, NPA 319, 518 (1979)

Low-energy Coulomb excitation
Example: ^{30}Mg + ^{58,60}Ni and ^{78}Kr + ^{208}Pb

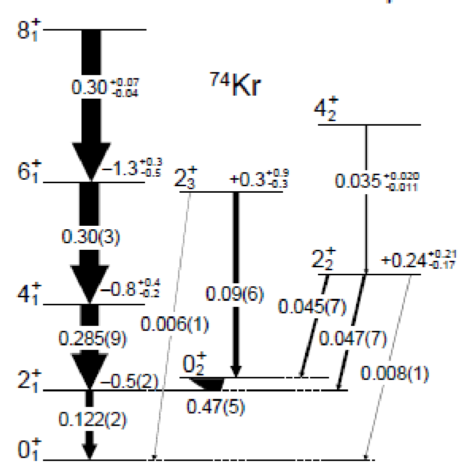
O. Niedermaier *et al.*, PRL 94, 172501 (2005)





^{30}Mg at 2.25 MeV/nucleon on natural Ni target (1.0 mg/cm^2)
From REX-ISOLDE at CERN
γ-ray detection with MINIBALL.
Particle detection with CD-shaped double-sided Si strip detector

$$\frac{\sigma_{CE}(^{30}\mathrm{Mg})}{\sigma_{CE}(^{58,60}\mathrm{Ni})} = \frac{\epsilon_\gamma(^{58,60}\mathrm{Ni})}{\epsilon_\gamma(^{30}\mathrm{Mg})} \frac{W_\gamma(^{58,60}\mathrm{Ni})}{W_\gamma(^{30}\mathrm{Mg})} \frac{N_\gamma(^{30}\mathrm{Mg})}{N_\gamma(^{58,60}\mathrm{Ni})},$$

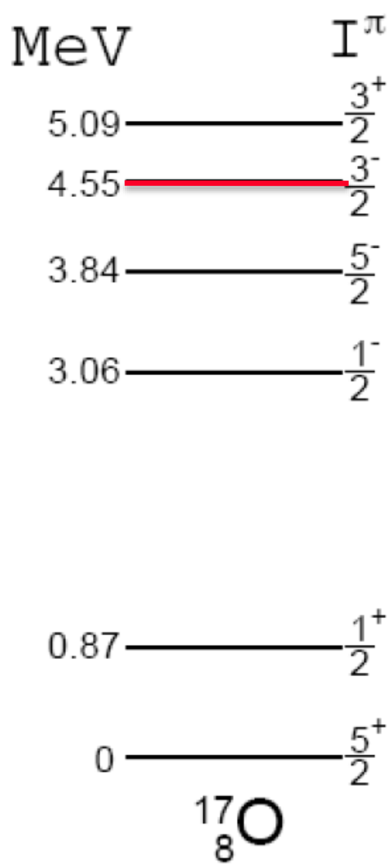


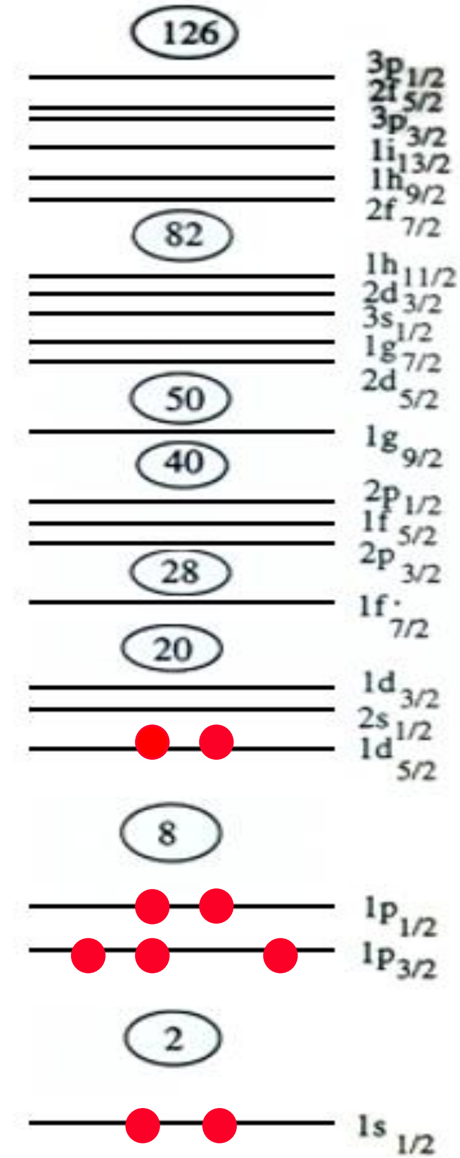
^{74}Kr multistep Coulex at 4.7 MeV/u on 1mg/cm^2 ^{208}Pb target at GANIL.
Data analysis done in a χ^2 minimization with a coupled channels code (GOSIA)

E. Clement *et al.*, PRC 75, 054314 (2005)

Single-particle states

Excited states in nuclei with one nucleon outside a magic number

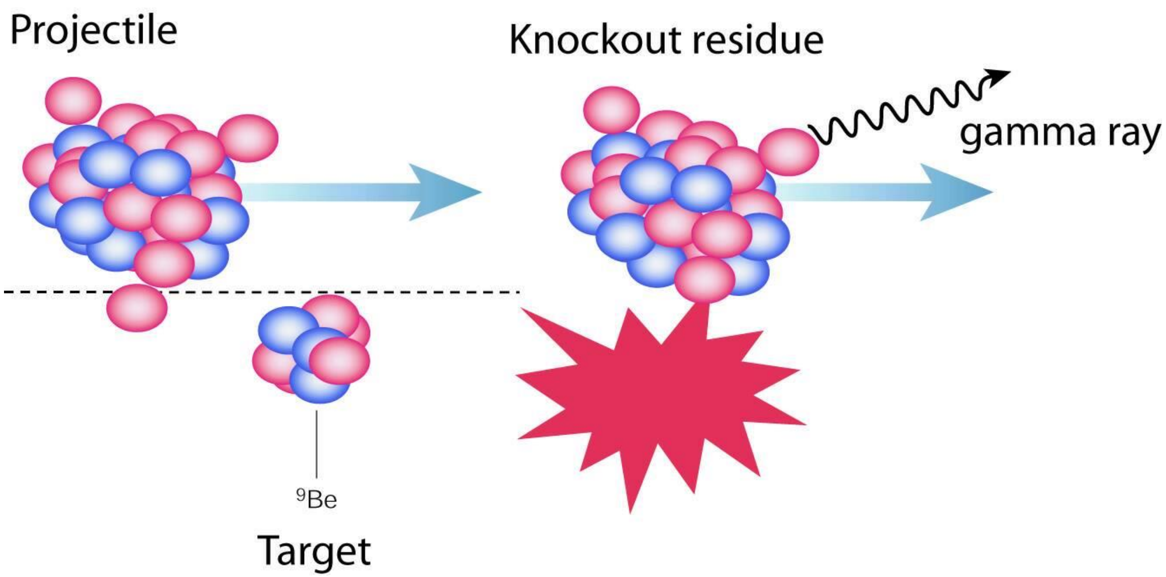


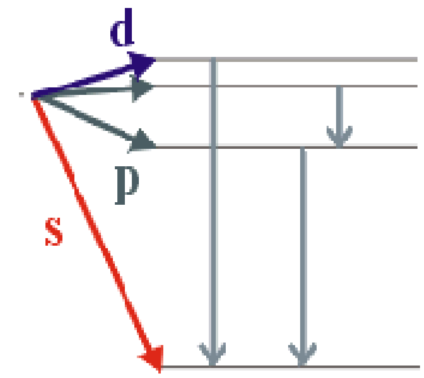


One-nucleon knockout
A direct reaction

- **more than 50 MeV/nucleon:**

  Straight-line trajectories



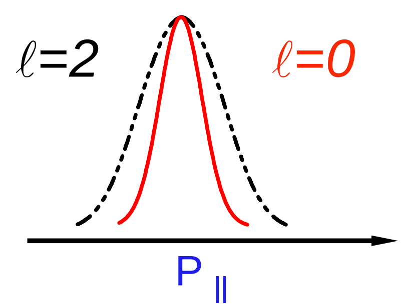


residue moment distribution
→ ℓ-value of knocked-out n

P.G. Hansen, PRL 77, 1016 (1996)

$$\sigma(nl^{\pi}) = C^2S(j, nl^{\pi})\ \sigma_{sp}(j, S_n)$$

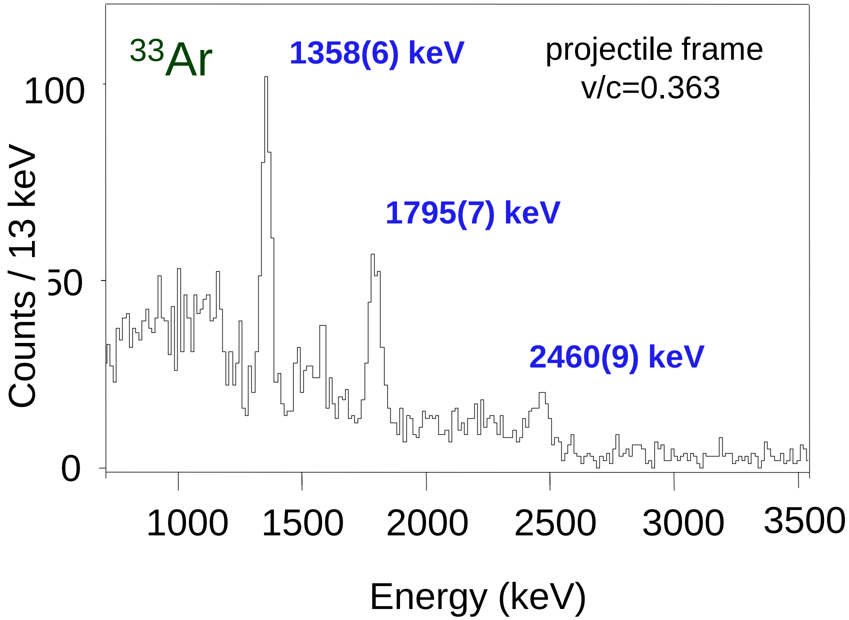
\# nucleons in orbit — reaction cross section

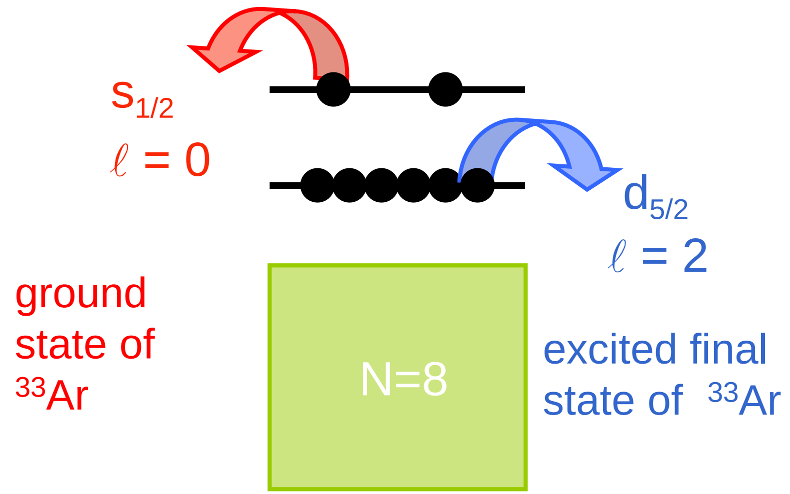


P.G. Hansen and B.M. Sherrill, NPA 693 ,133 (2001).
P.G. Hansen and J. A. Tostevin, Annu. Rev. of Nucl. and Part. Sci. 53, 219 (2003).

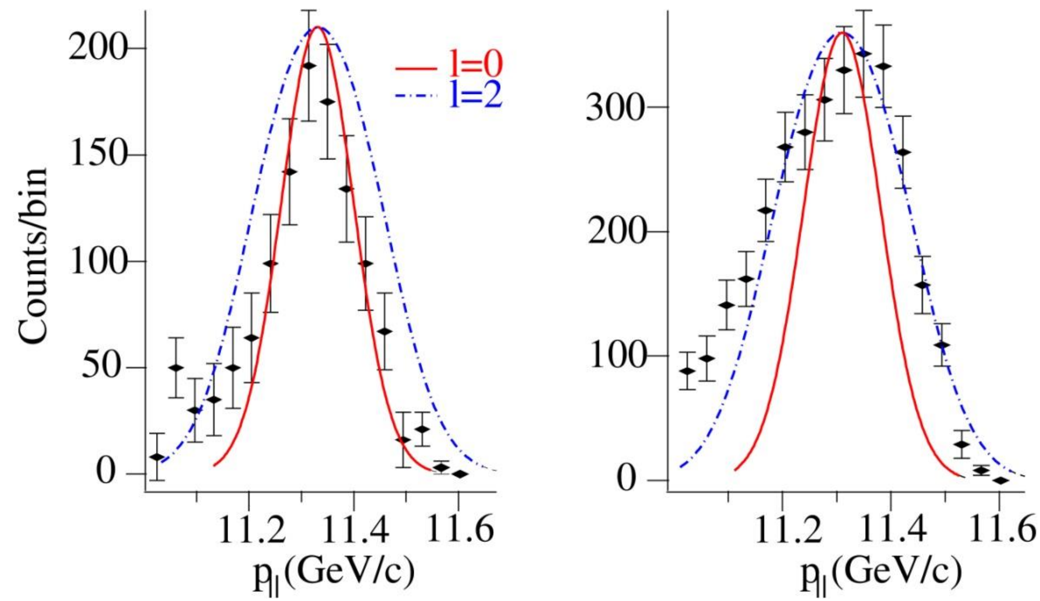
Spectroscopy in one-nucleon knockout

Example: $^{9}Be(^{34}Ar,^{33}Ar)X$



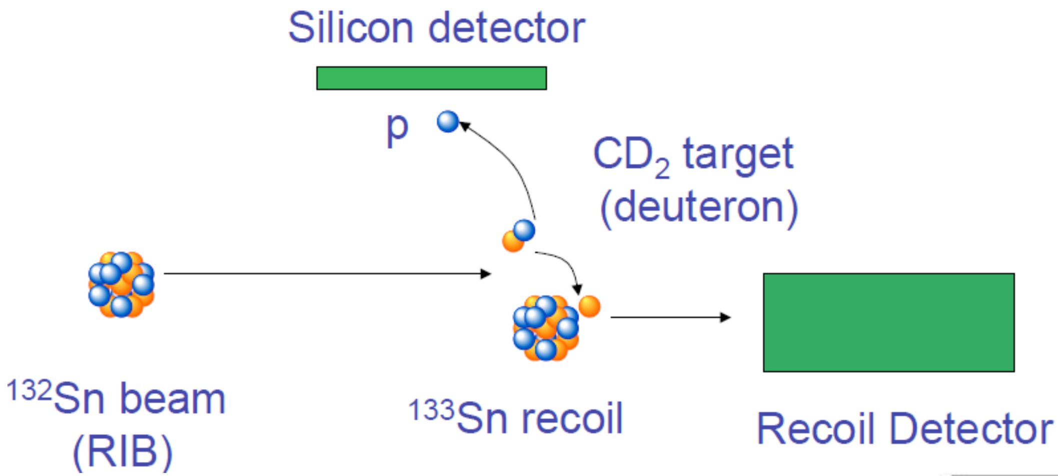


	BR (%)	σ_{exp} (mb)	C^2S_{exp}
$1/2^+$	30.2(46)	4.7(9)	0.38(6)
$3/2^+$	20.2(44)	3.2(8)	0.36(9)
$5/2^+$	31.7(31)	4.9(7)	0.56(8)
$(5/2^+)$	17.9(30)	2.8(6)	>0.34(7)



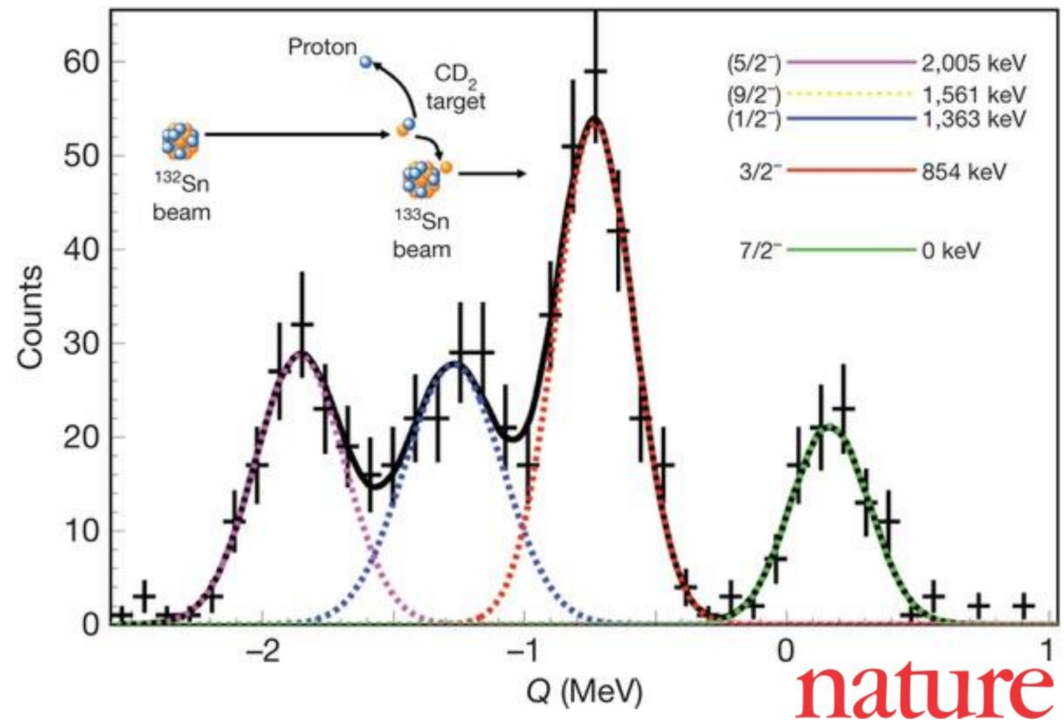
A. Gade *et al.*, PRC 69 034311 (2004).

Low-energy transfer reactions – $^{132}Sn(d,p)^{133}Sn$ at HRIBF



Q-value spectrum for the $^{132}Sn(d,p)^{133}Sn$ reaction at 54° in the centre of mass.

- 4.77 MeV/u ^{132}Sn produced and accelerated at HRIBF bombarded a 160μg/cm^2 CD_2 target. Exit-channel proton detection with ORRUBA Si strip detectors under 69-107° polar angles



KL Jones *et al.* *Nature* **465**, 454-457 (2010) doi:10.1038/nature09048

Lifetimes of excited states

Can provide information on collective and single-particle degrees of freedom

Lifetimes of excited states

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Lifetimes of excited 2^+ states in even-even nuclei: picosecond range

$$\tau_\gamma = 40.81 \times 10^{13} E^{-5} [B(E2)\uparrow / e^2 b^2]^{-1}$$

Some excited states live much longer: Isomers

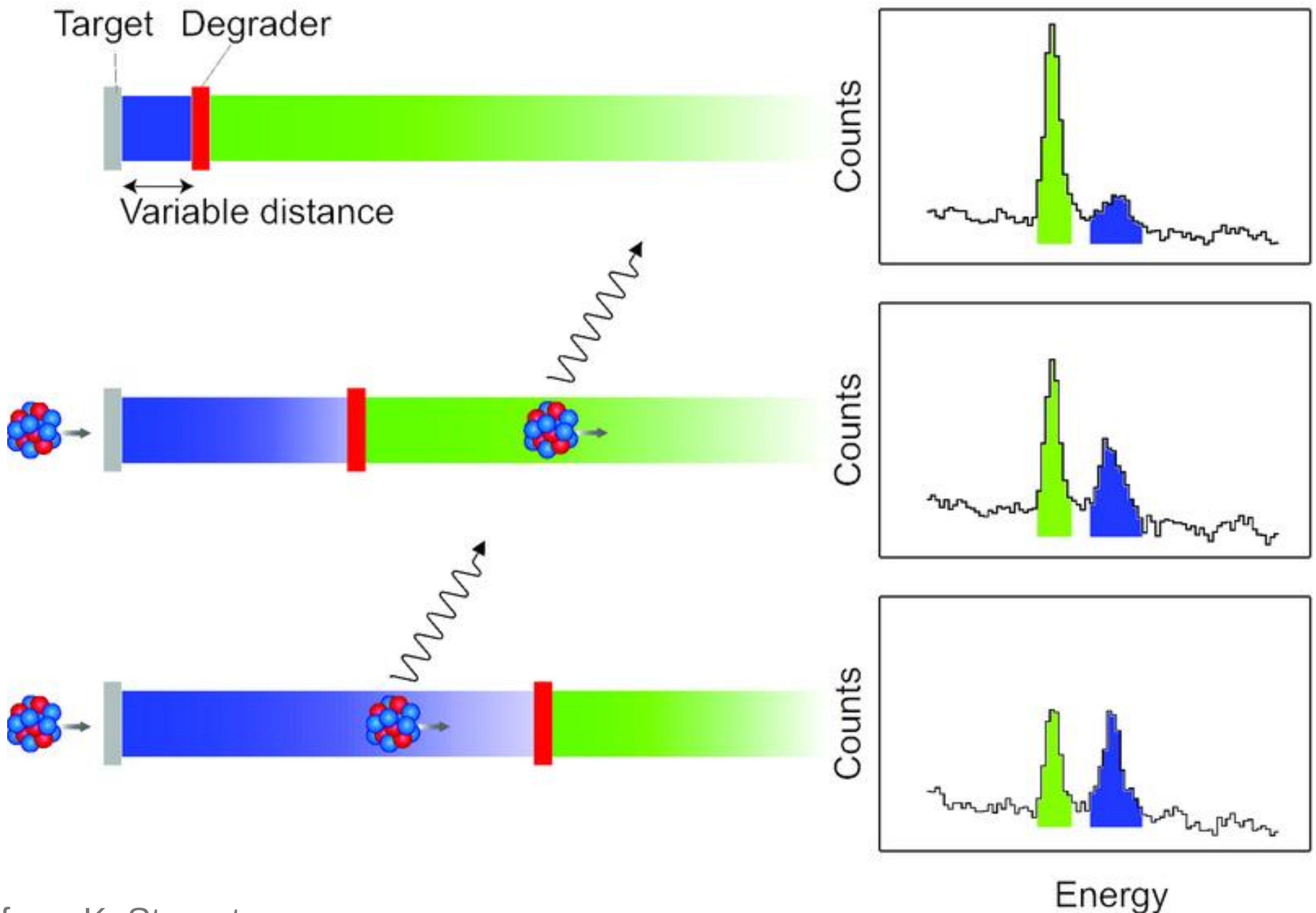
Table I: Examples of extreme isomers

Nuclide	Half-life	Spin (ħ)	Energy	Attribute
^{12}Be	~500 ns	0	2.2 MeV	low mass
^{94}Ag	300 ms	21	6 MeV	proton decay
^{152}Er	11 ns	~36	13 MeV	high spin and energy
^{180}Ta	>10^{16} y	9	75 keV	long half-life
^{229}Th	~5 h	3/2	~7.6 eV	low energy
^{270}Ds	~6 ms	~10	~1 MeV	high mass

From P.M. Walker and J. J. Carroll, Nuclear Physics News 17, 11-15 (2007)

Plunger lifetime measurements

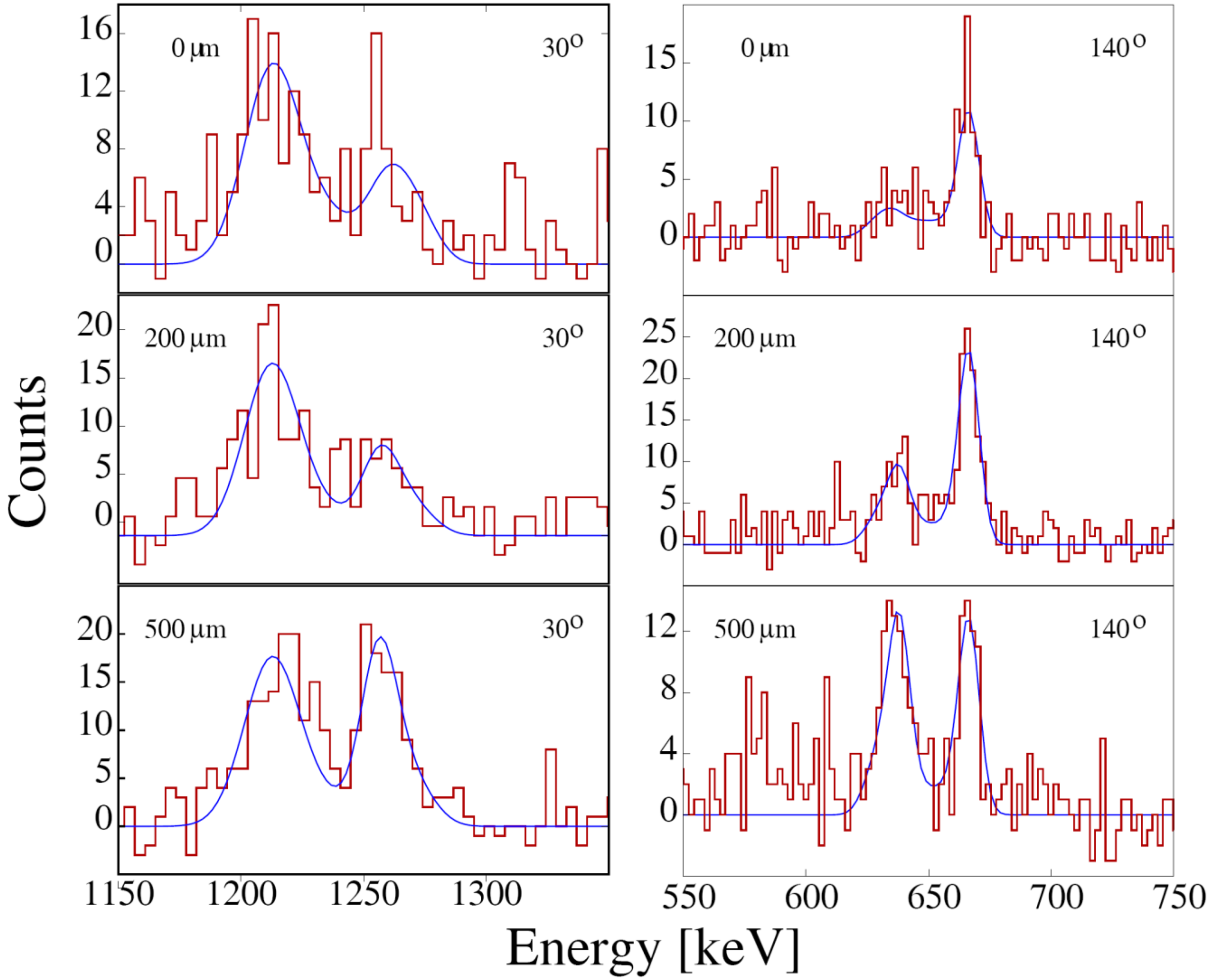
c=300 μm/ps **β ~ 0.3c** **10 ps ~ 1mm**



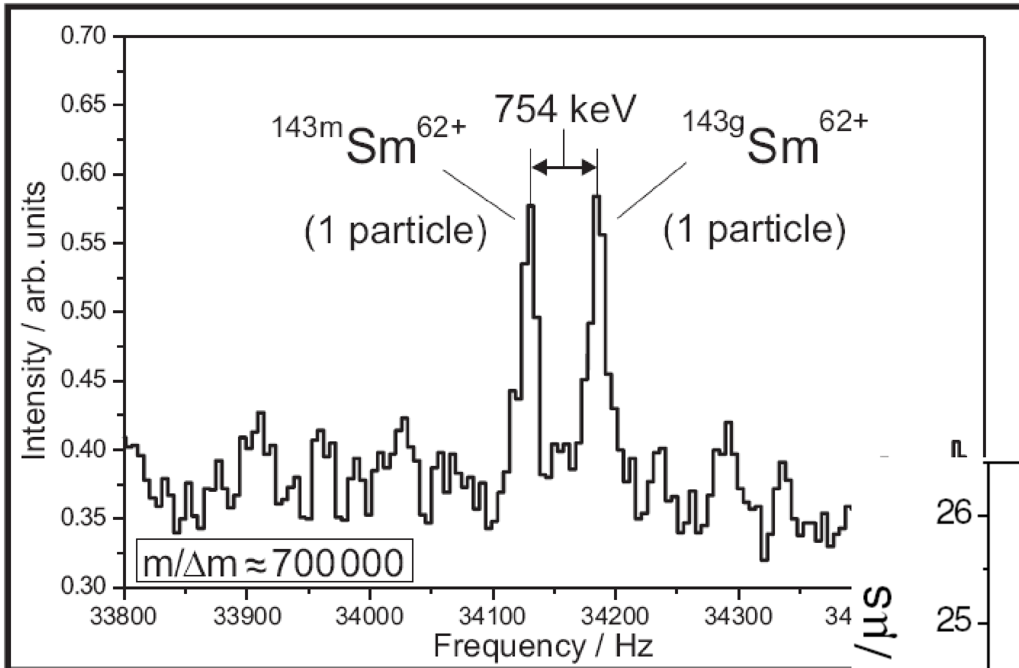
Adapted from K. Starosta

Line shapes and lifetimes
Example: ^{64}Ge $2^+_1 \rightarrow 0^+_1$

τ=3.2(5) [ps]

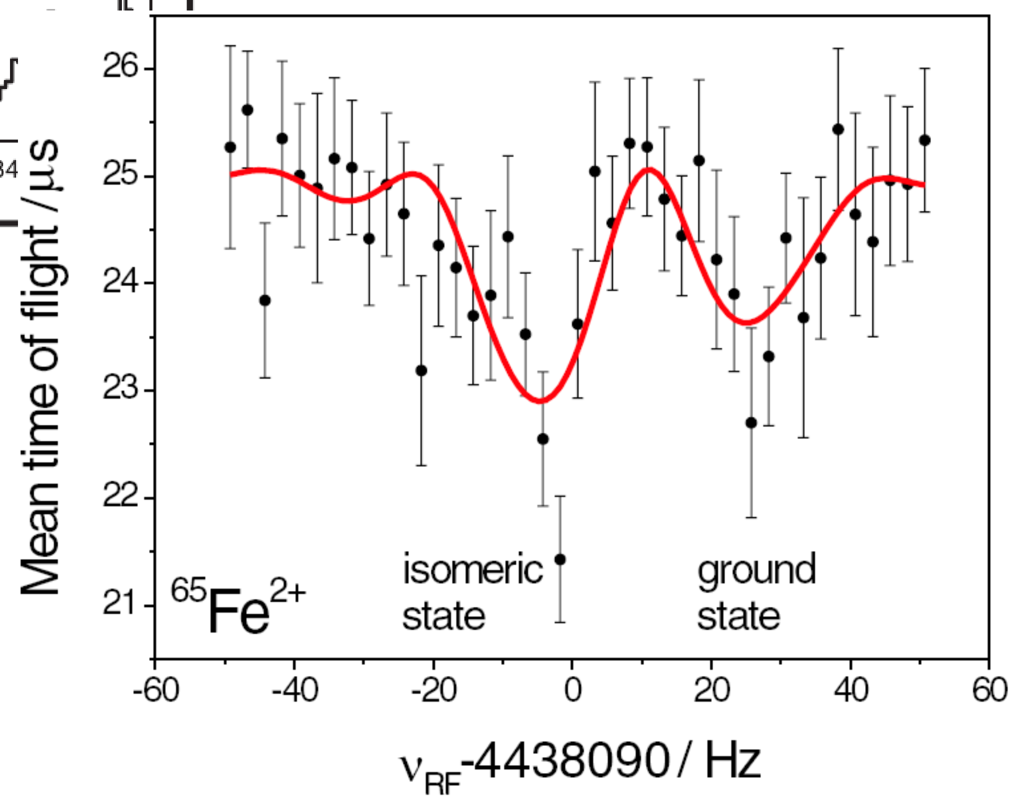


Long-lived excited states – isomers
Back to storage rings and penning traps



F. Bosch, Lect. Notes Phys. 651, 137(2004)

M. Block et al., PRL 100, 132501 (2008)



Isomers: decay hindered by nuclear structure (selection rules, energy, …) → long lifetime

Take away

- Excited states provide valuable information on the evolution of nuclear structure
  - Population of excited states in various schemes
- Reactions – powerful tools
  - Observables related to the collective degree of freedom
  - Single-particle structure from direct reactions
- Life-times of excited states
  - Different experimental approaches