

EXPERIMENTS WITH RADIOACTIVE BEAMS

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Menu

Production methods

Types of experiments

Reaction experiments

Issues and solutions

Looking towards the future

Production methods

- ▶ ISOL technique (Isotope Separation On-Line)
 - ▶ Fragment stable target nuclei with high intensity light beam (usually protons)
 - ▶ Radioactive nuclei (fragments) thermally diffuse out of target, are ionized, filtered and accelerated into a beam

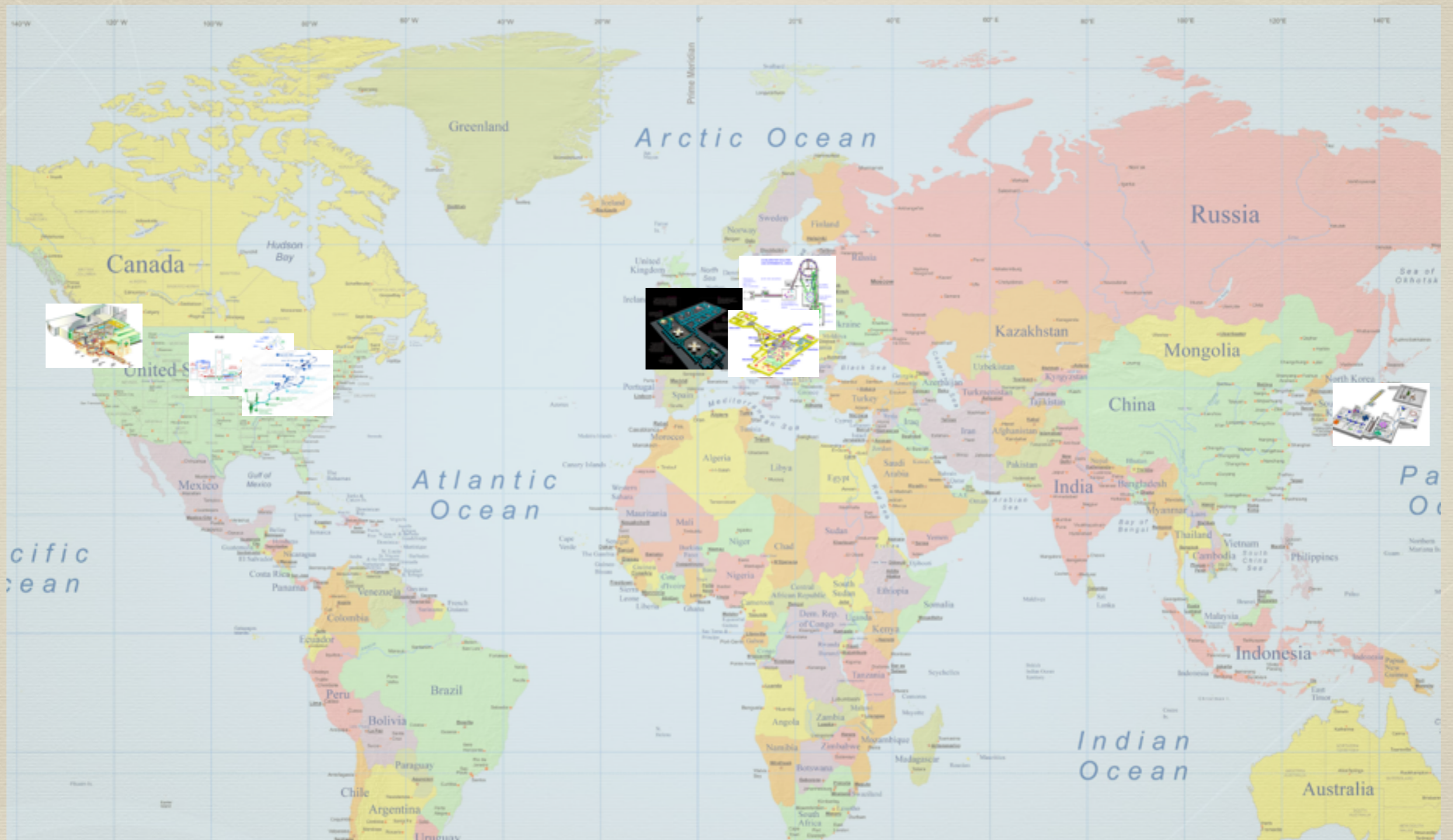
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- ▶ Projectile fragmentation technique
 - ▶ Fragment high intensity stable beam nuclei on light target (usually Beryllium or Carbon)
 - ▶ Radioactive nuclei (fragments) fly out of target, are filtered and shaped into a beam

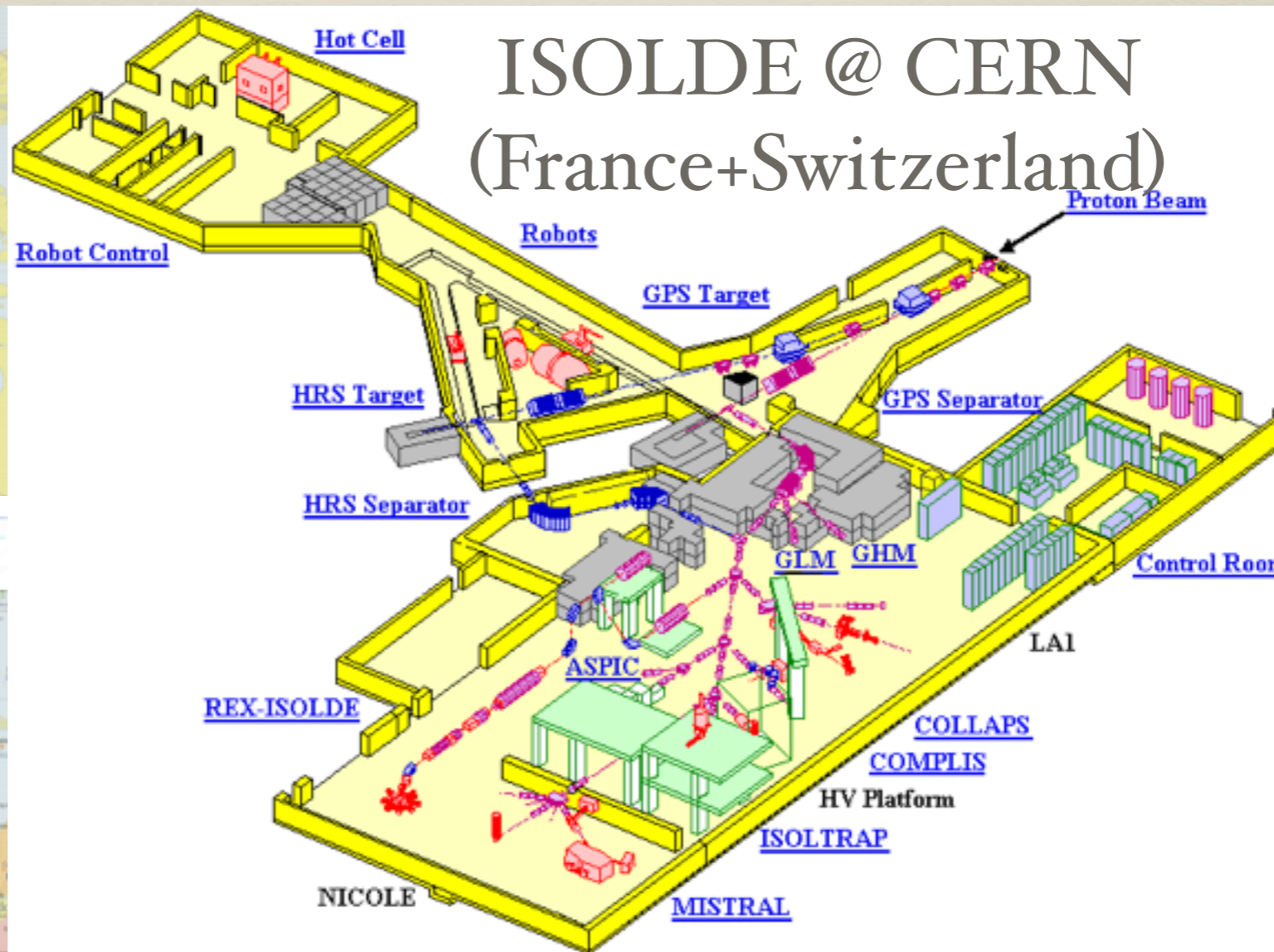
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- ▶ Other methods
 - ▶ Transfer reactions, fission source, etc...

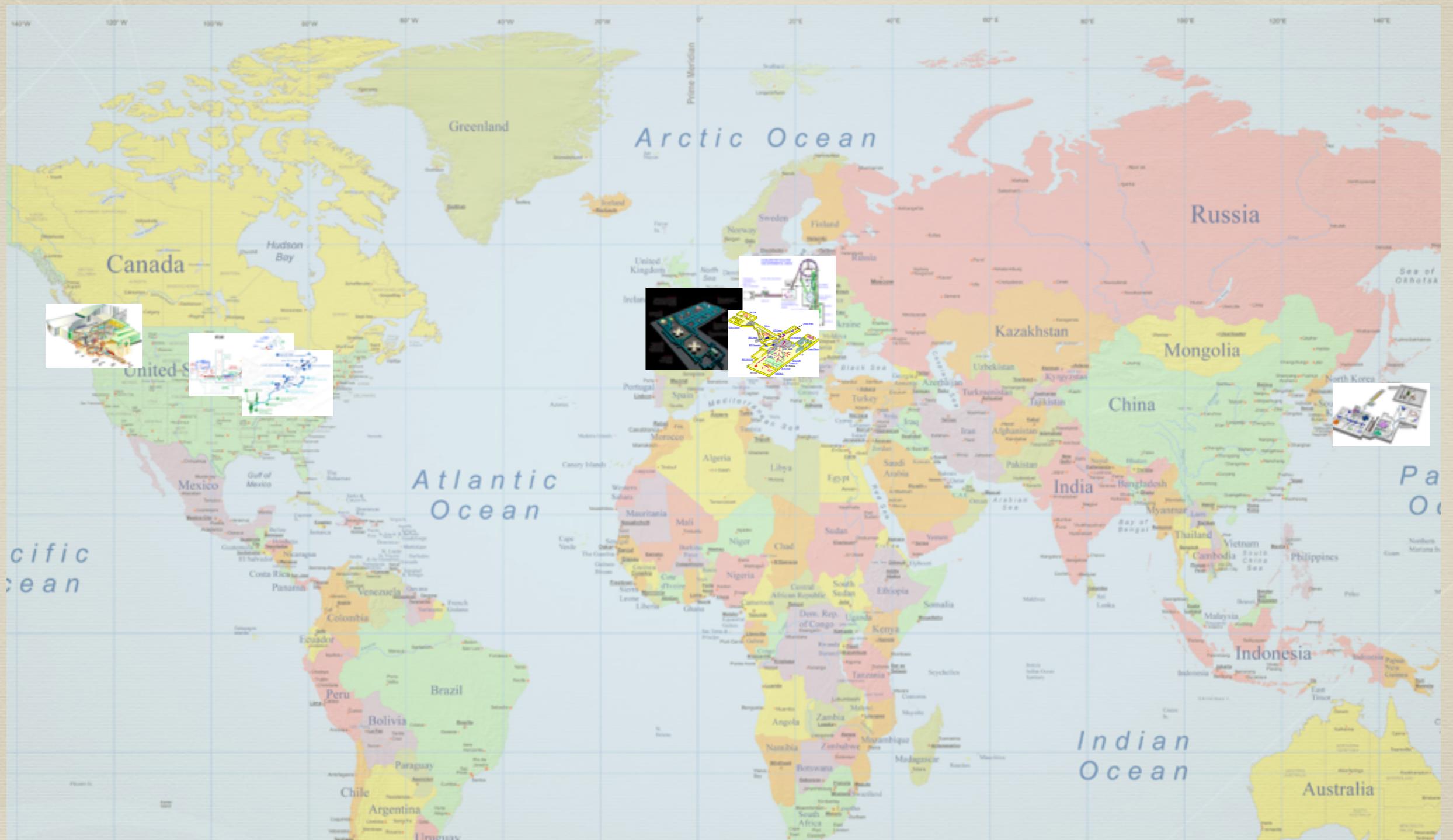
Main facilities in the world



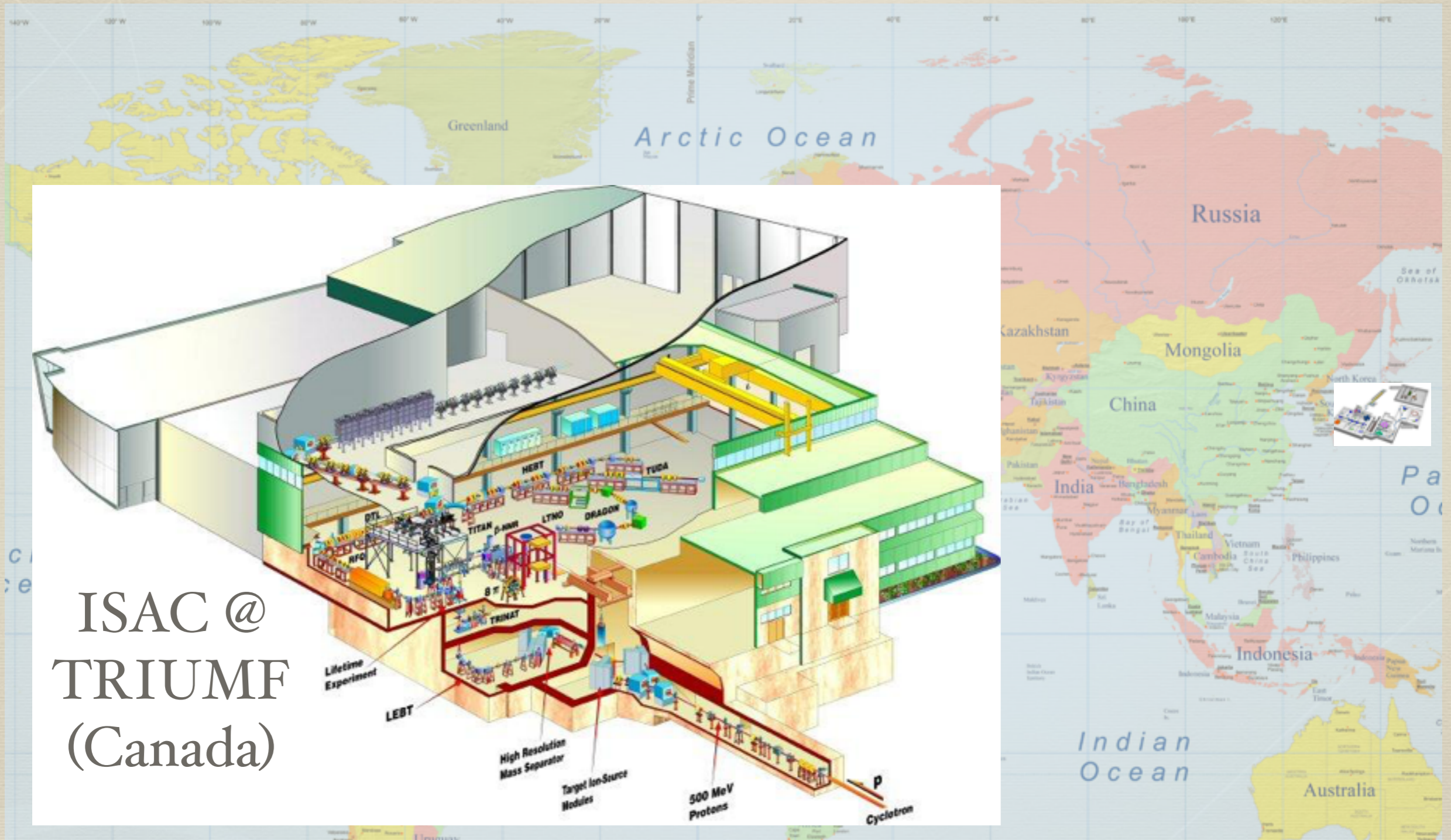
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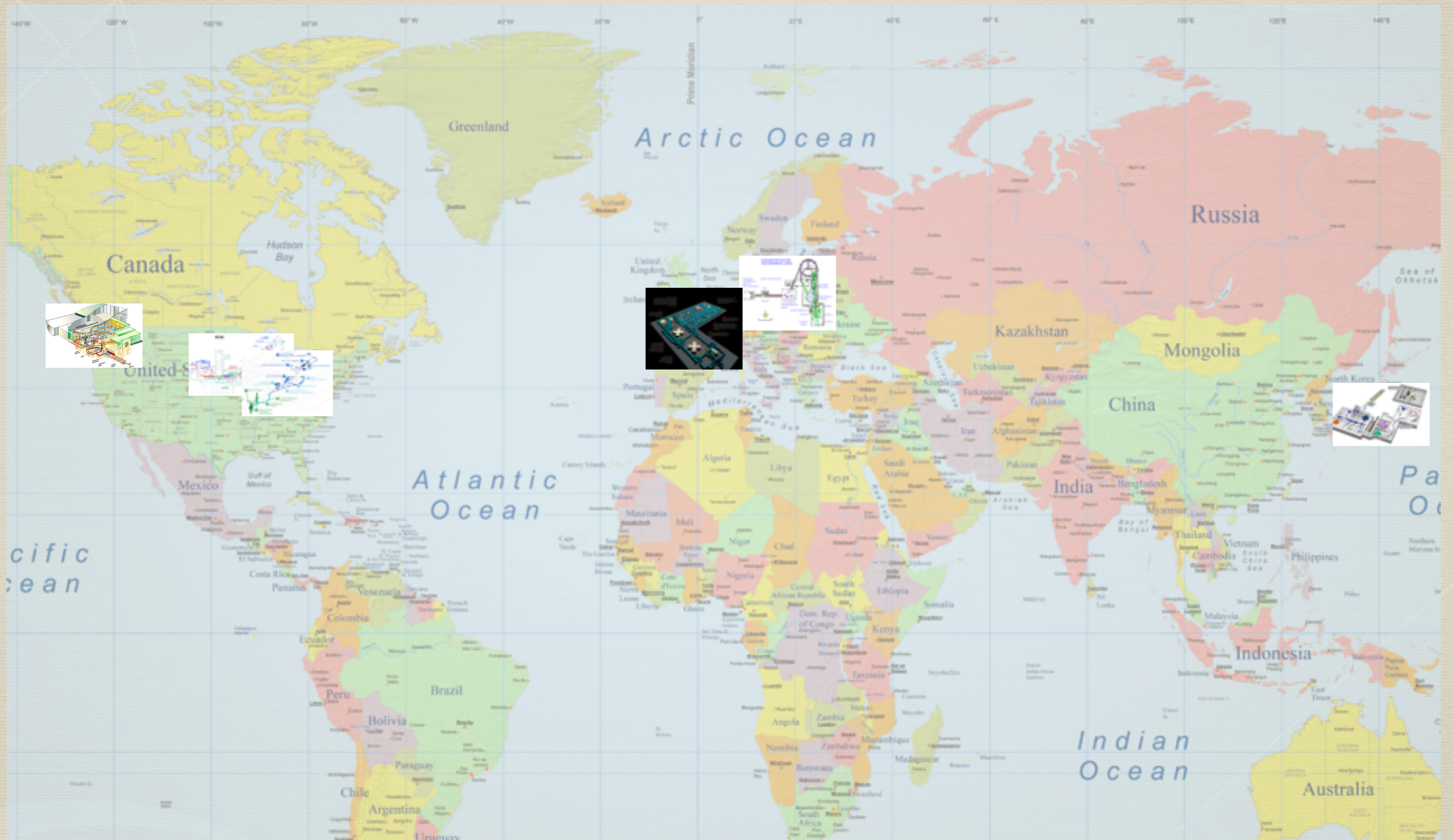


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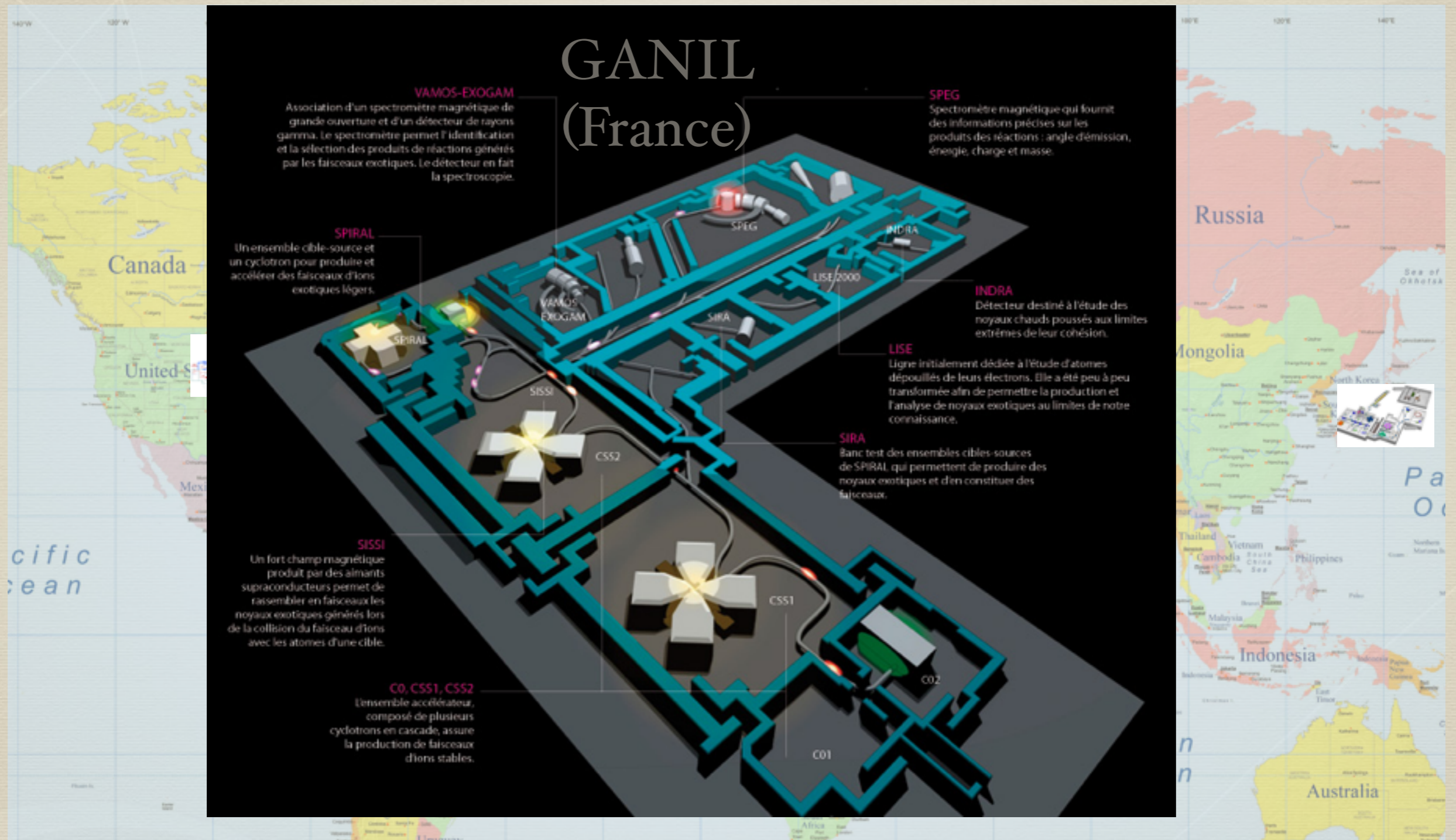


ISAC @
TRIUMF
(Canada)

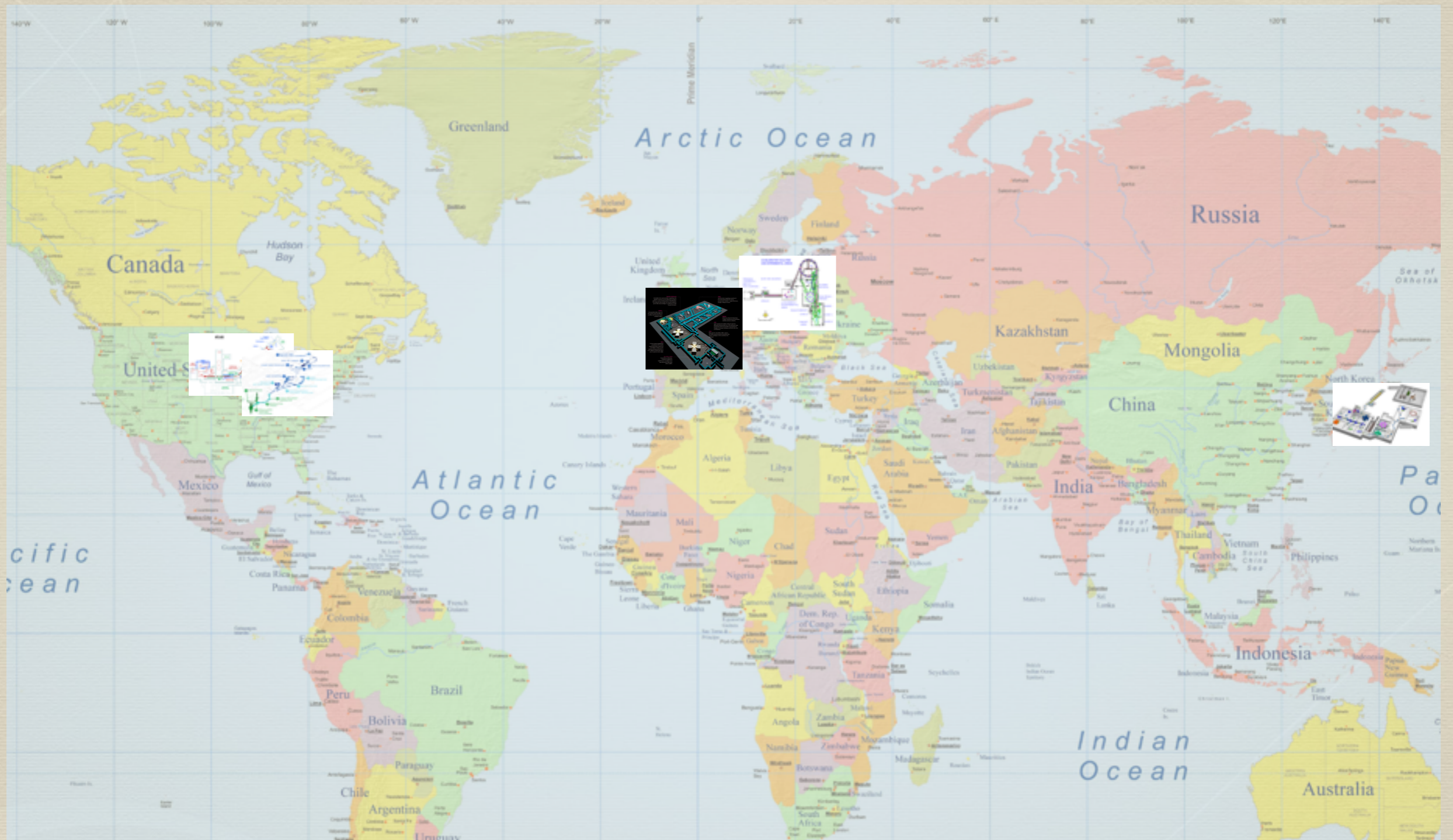
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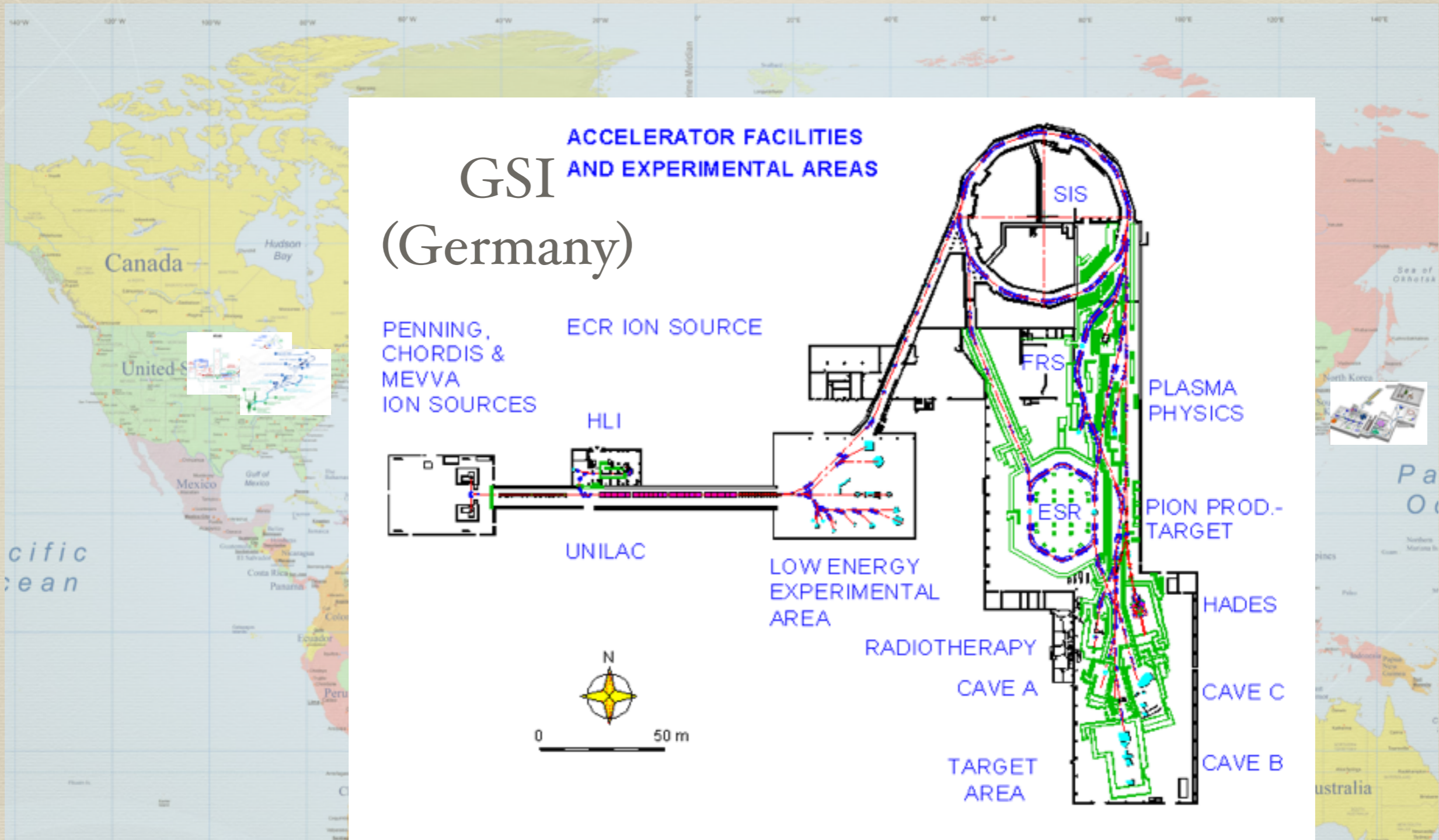
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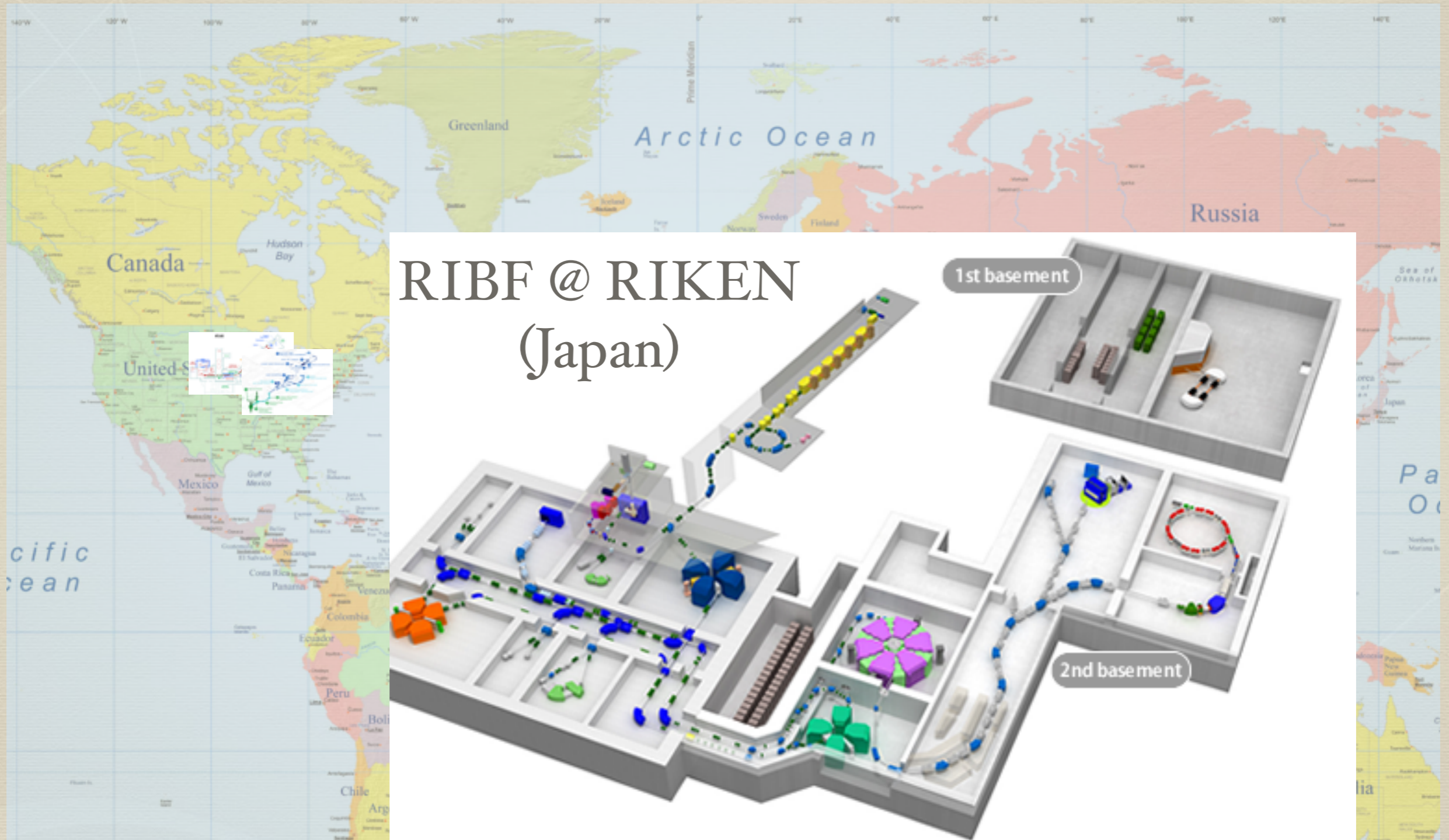
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RIBF @ RIKEN
(Japan)

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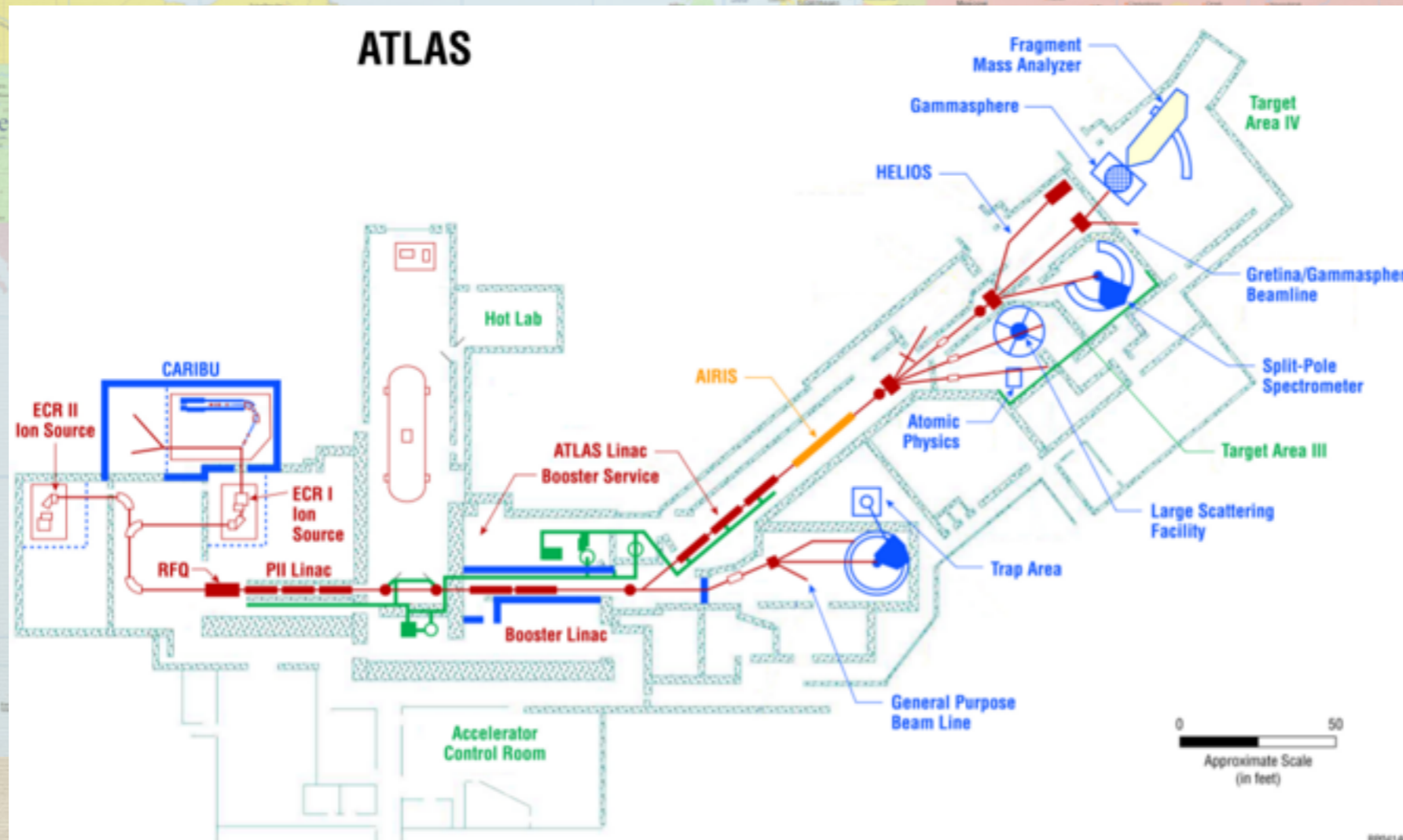
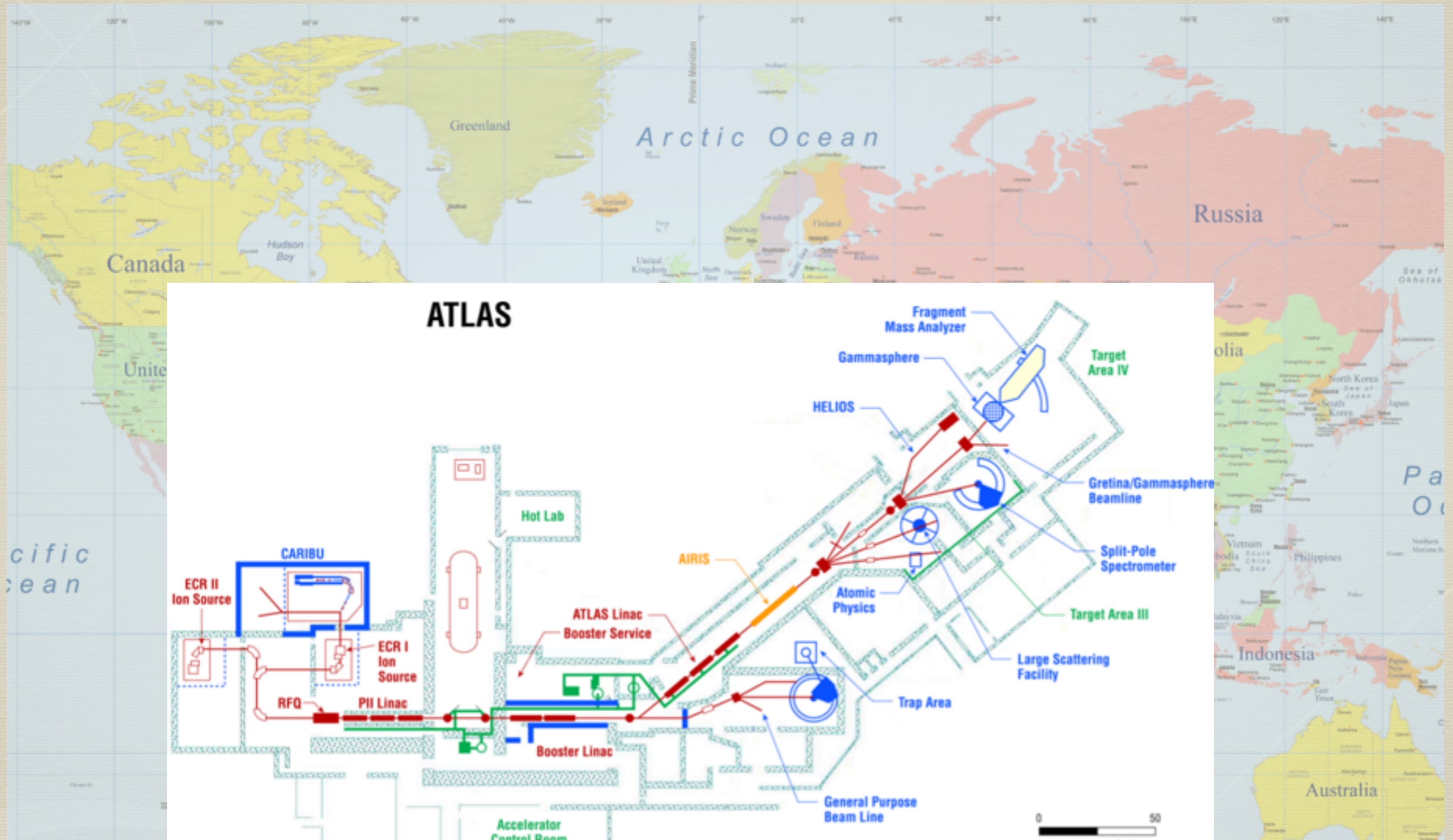
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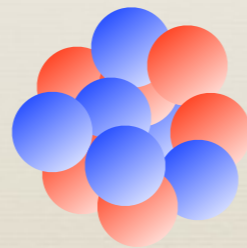


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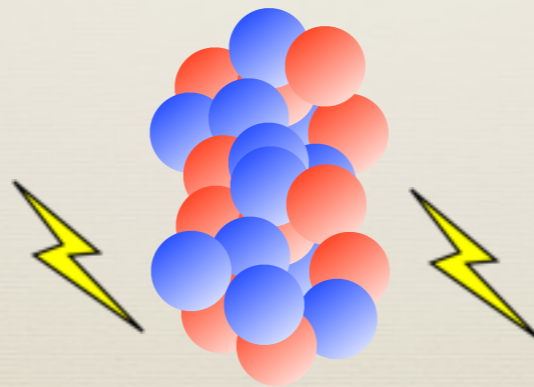
Projectile fragmentation

- ▶ Randomly cut stable nuclei into fragments
 - ▶ High energy (50 - 1 GeV/u or more)
 - ▶ “Spectator” nucleons form projectile fragment
 - ▶ Projectile fragments carry most of momentum
 - ▶ High efficiency of collecting them at forward angles
 - ▶ Thick targets to increase probability of nuclear reaction



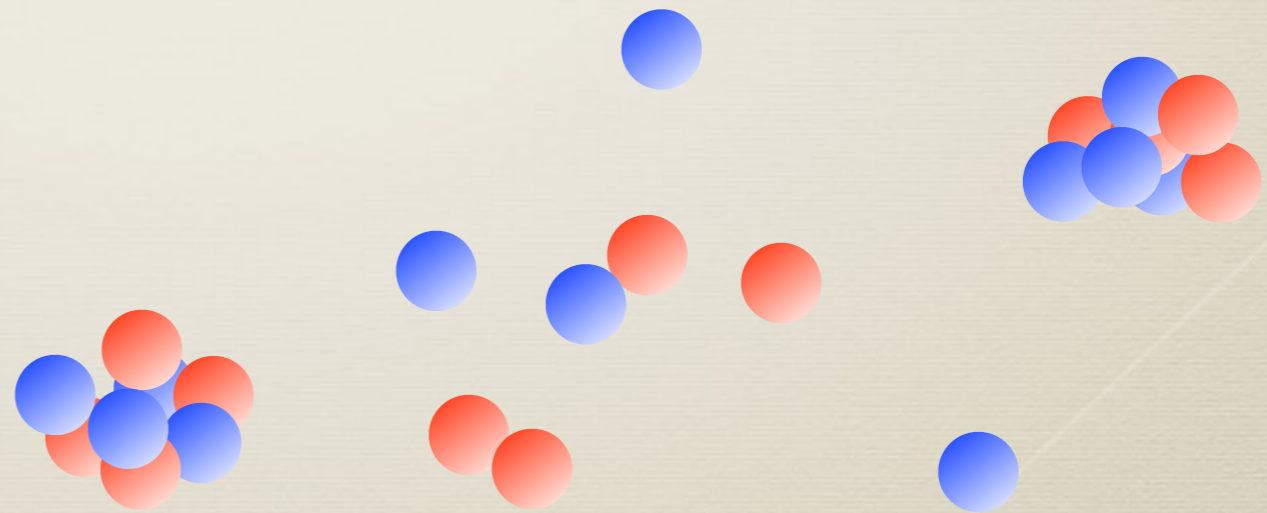
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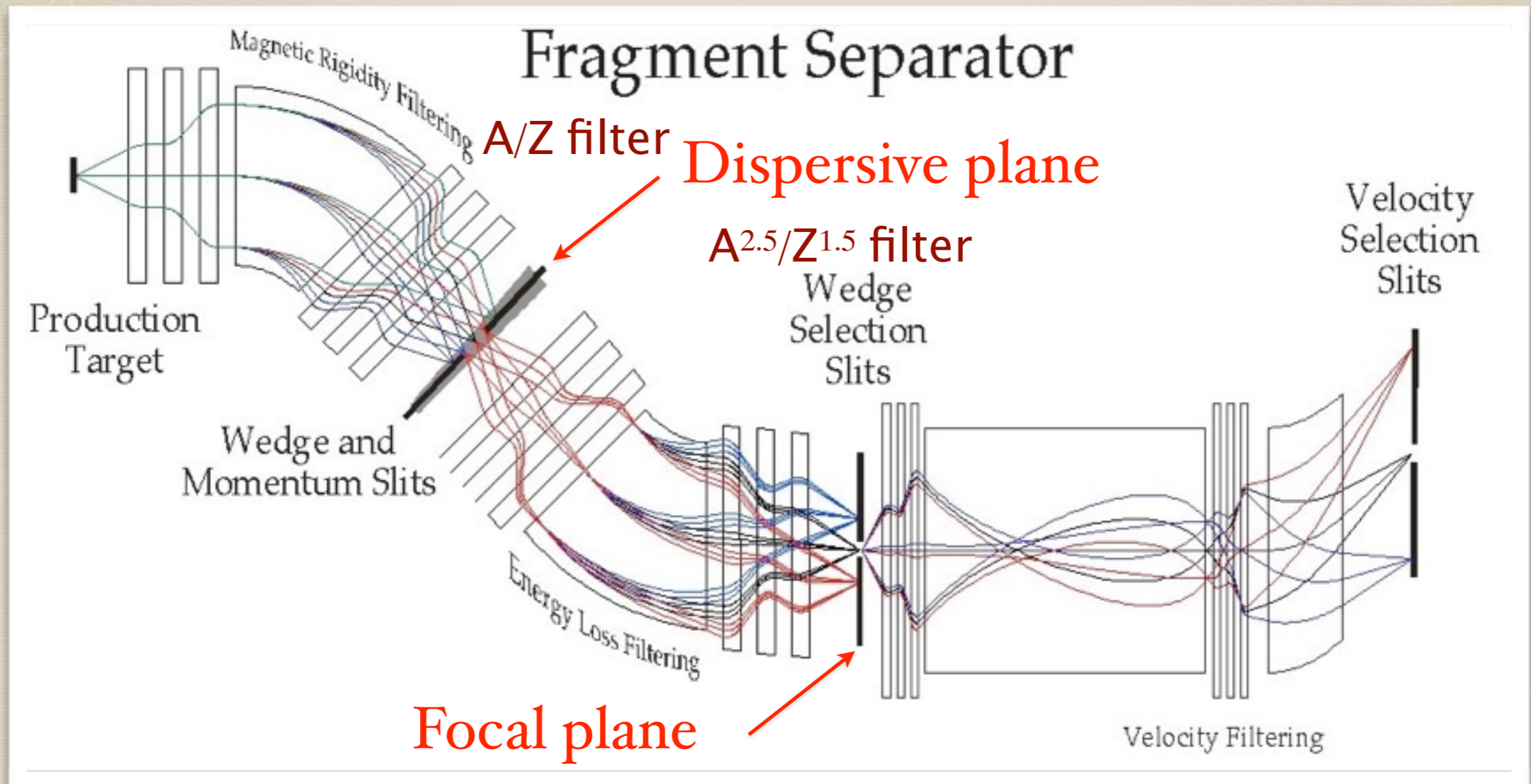


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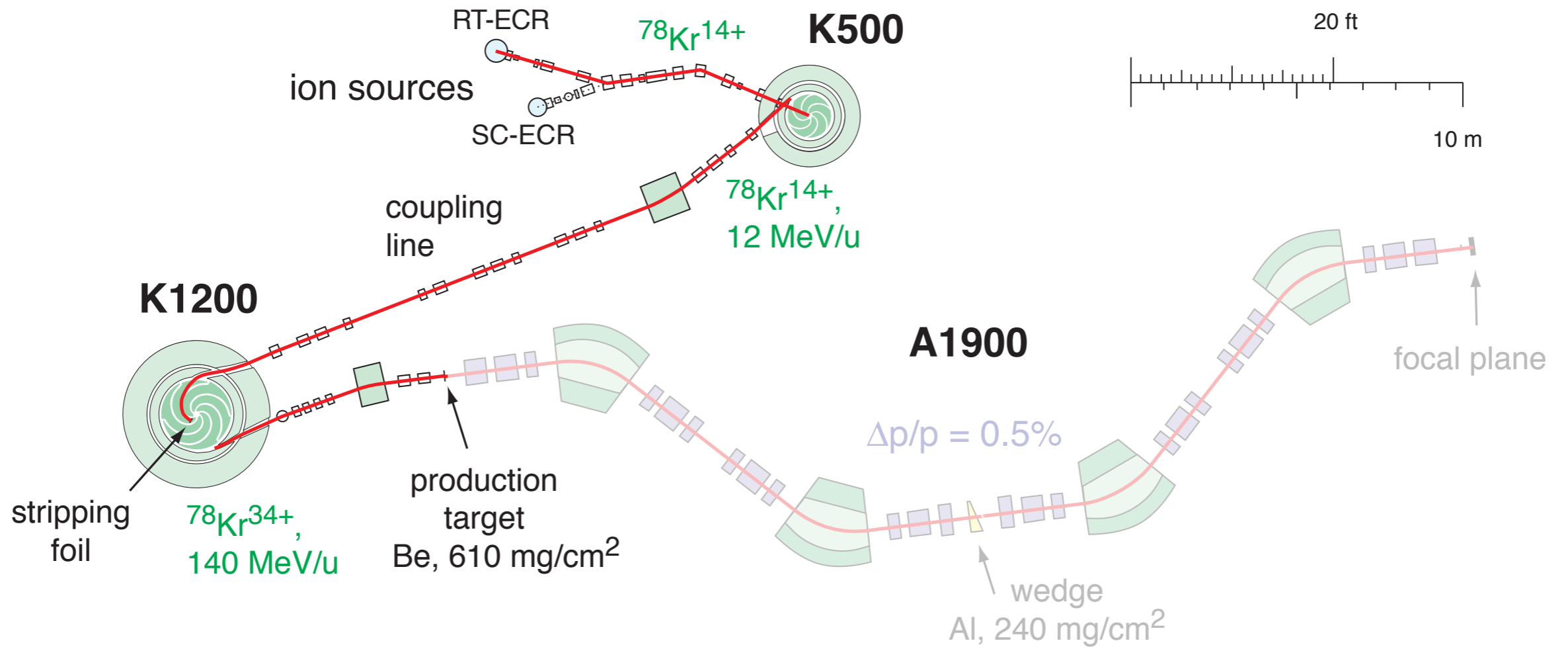


Fragment separator

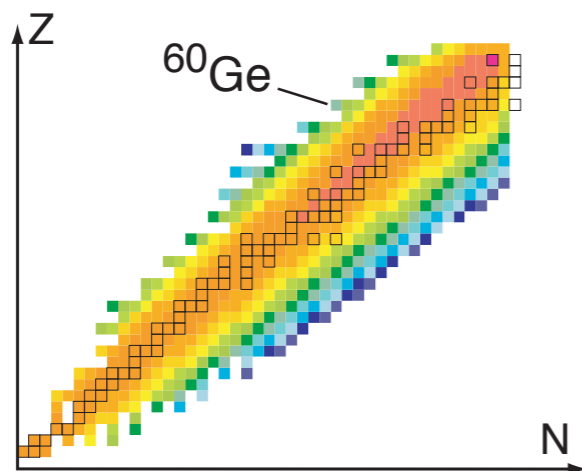


- ▶ Select forward focused fragments produced from projectile fragmentation reactions
- ▶ Use various selection criteria: magnetic rigidity ($B\rho$), energy loss (wedge), velocity (Wien filter, RF Separator)

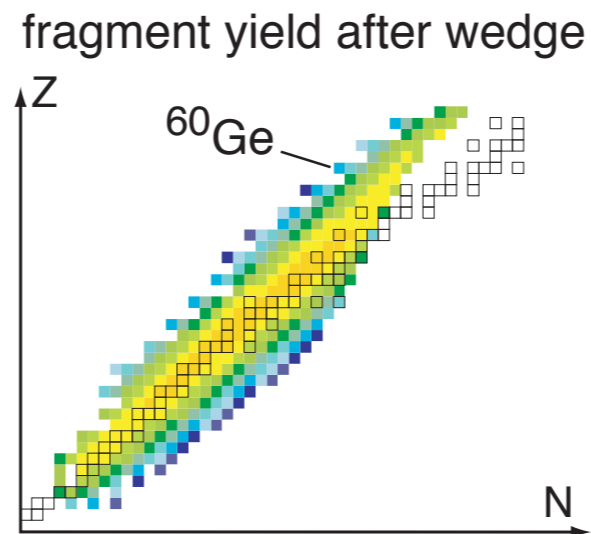
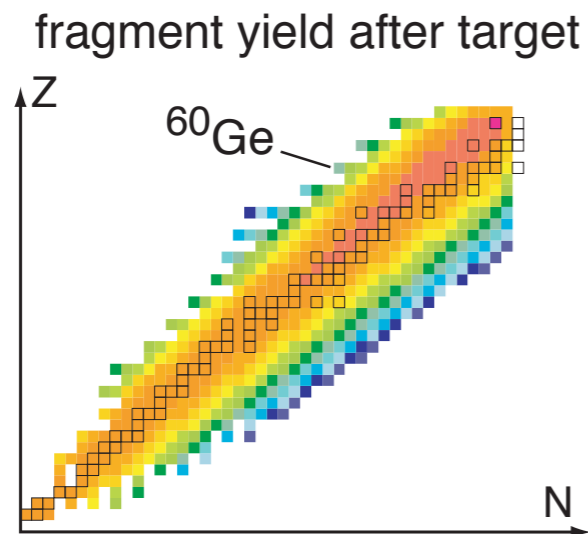
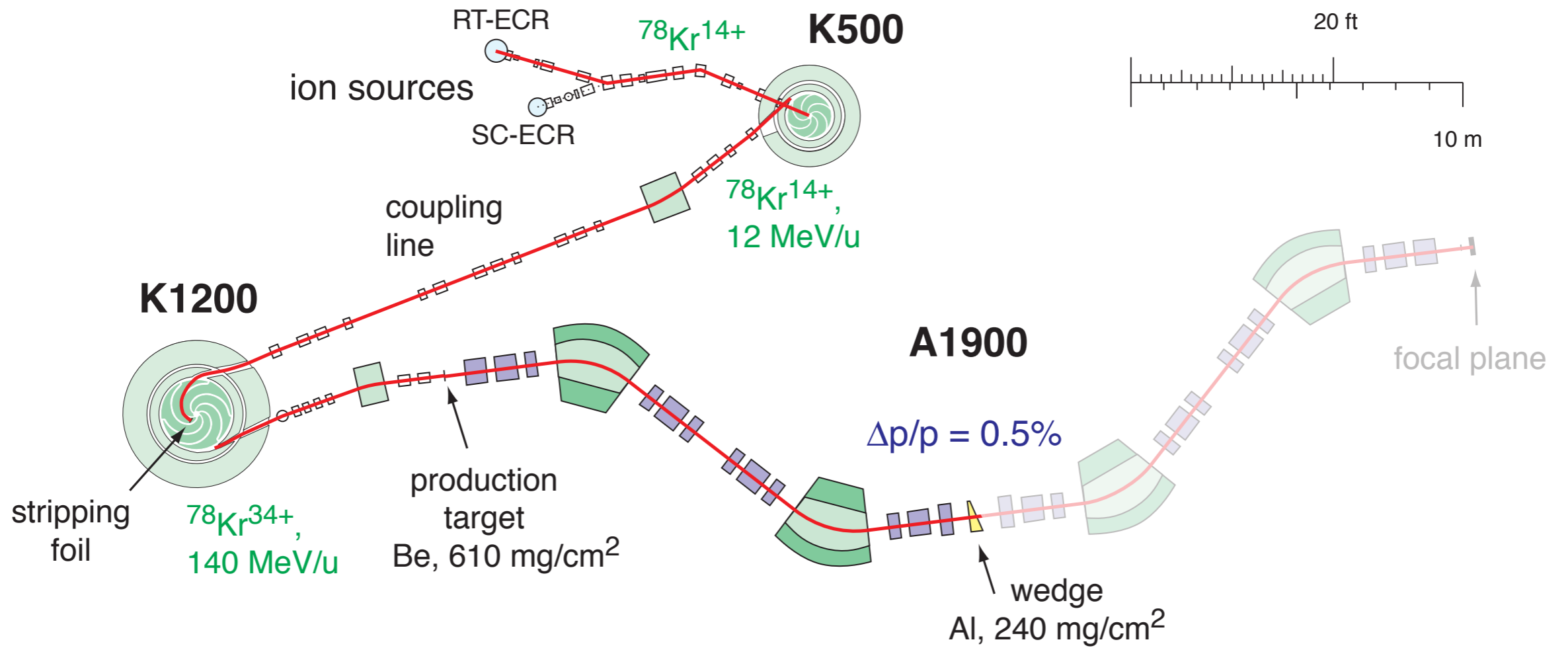
Experimental Setup at the NSCL



fragment yield after target

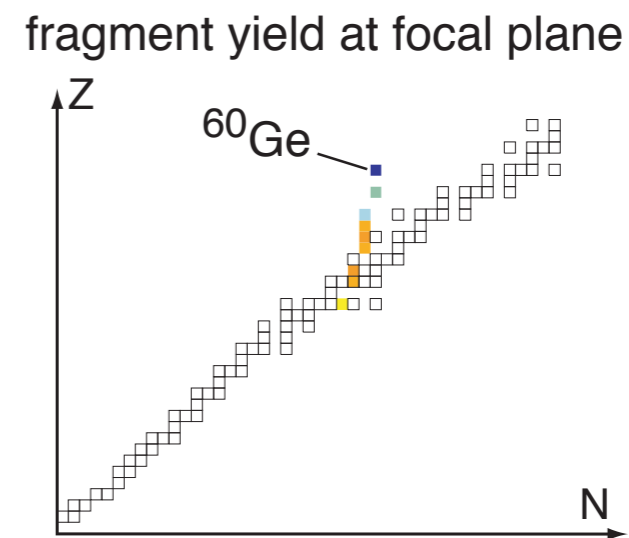
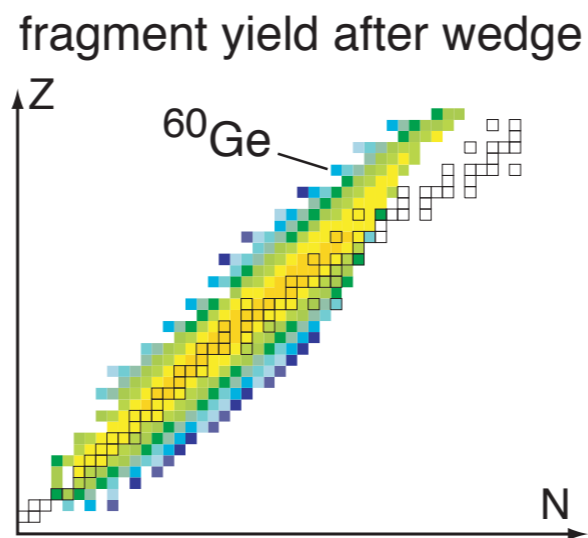
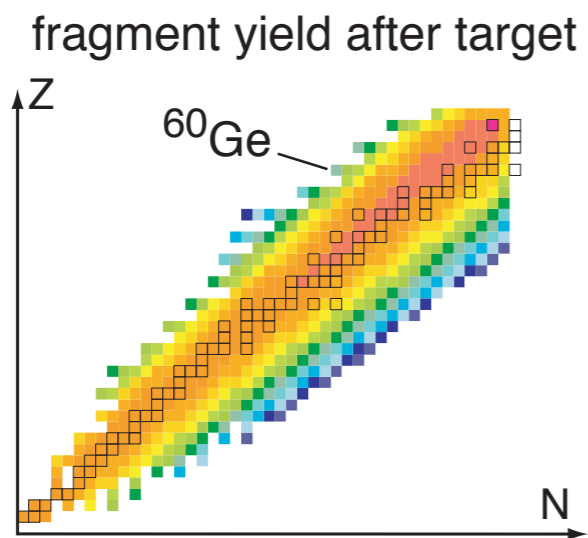
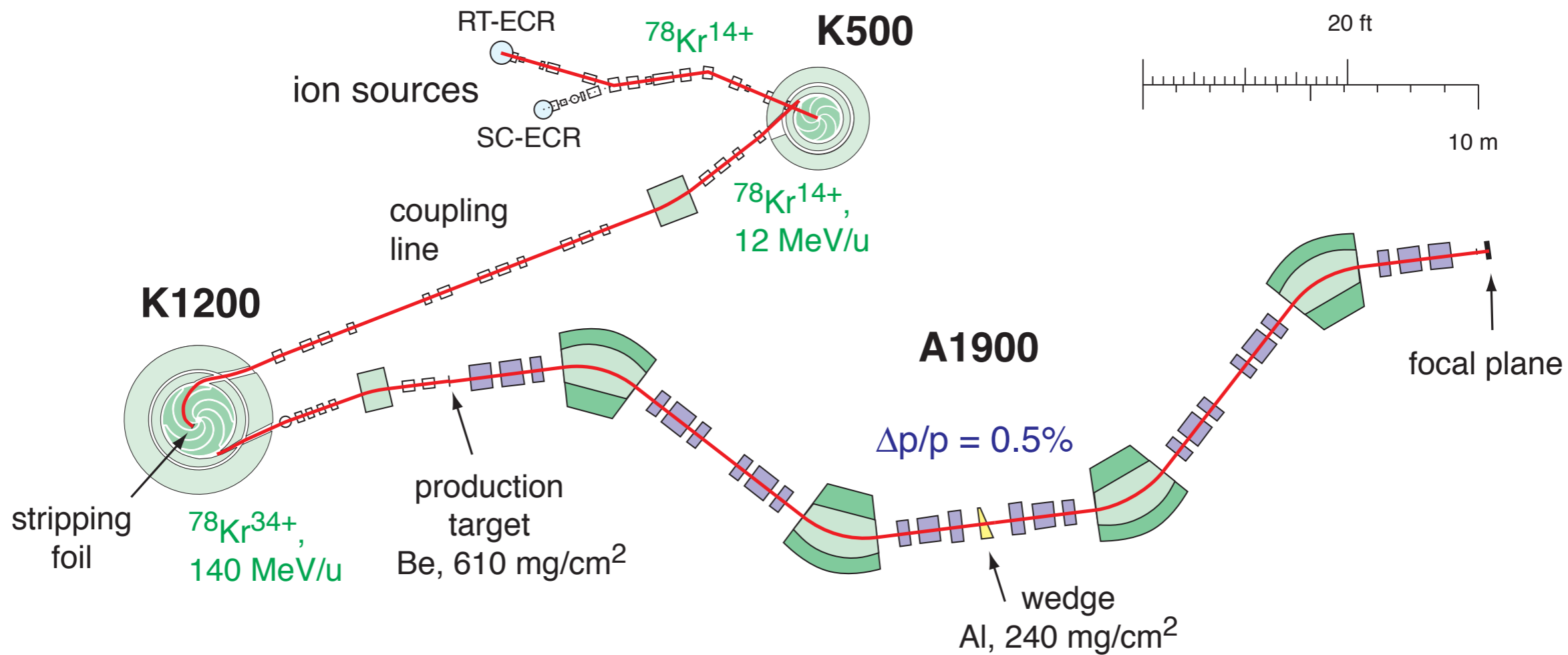


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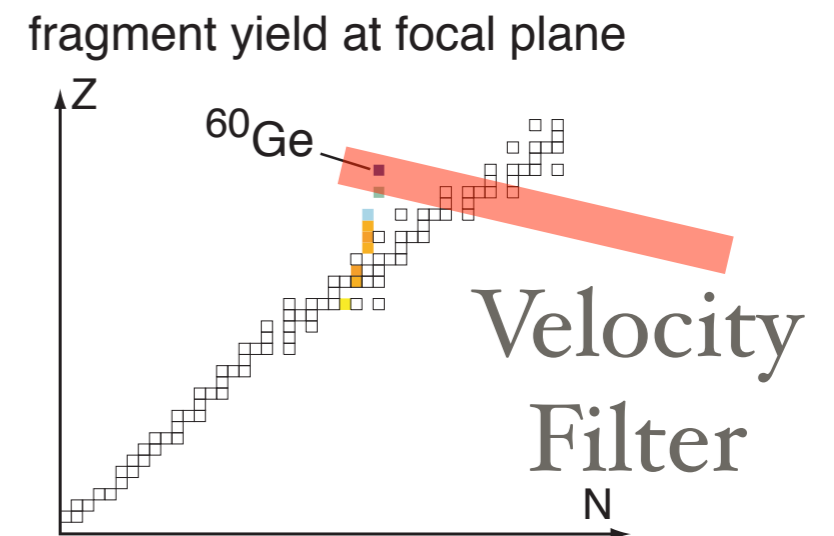
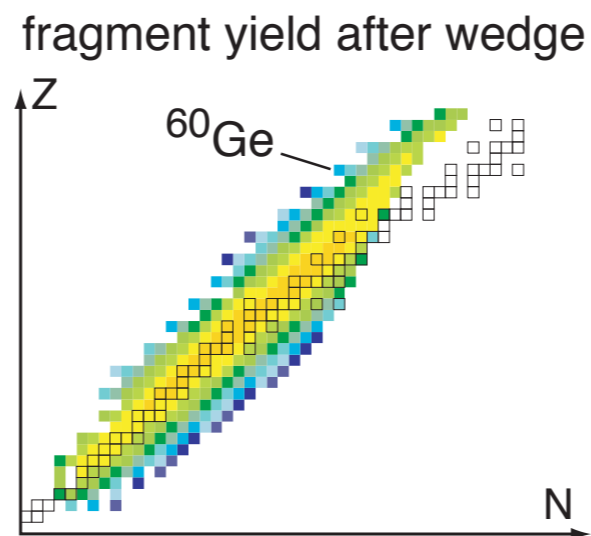
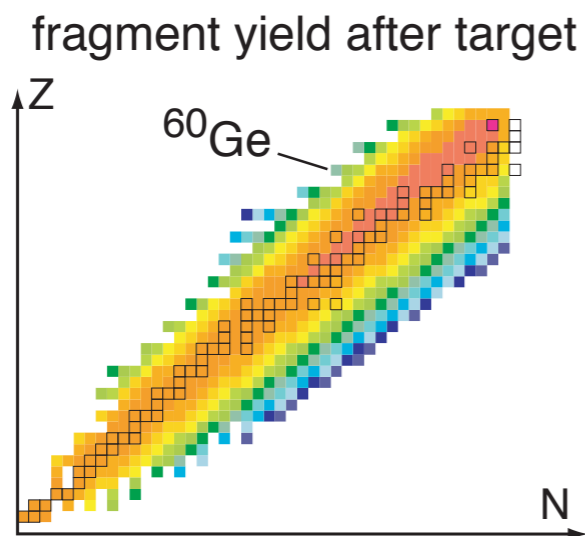
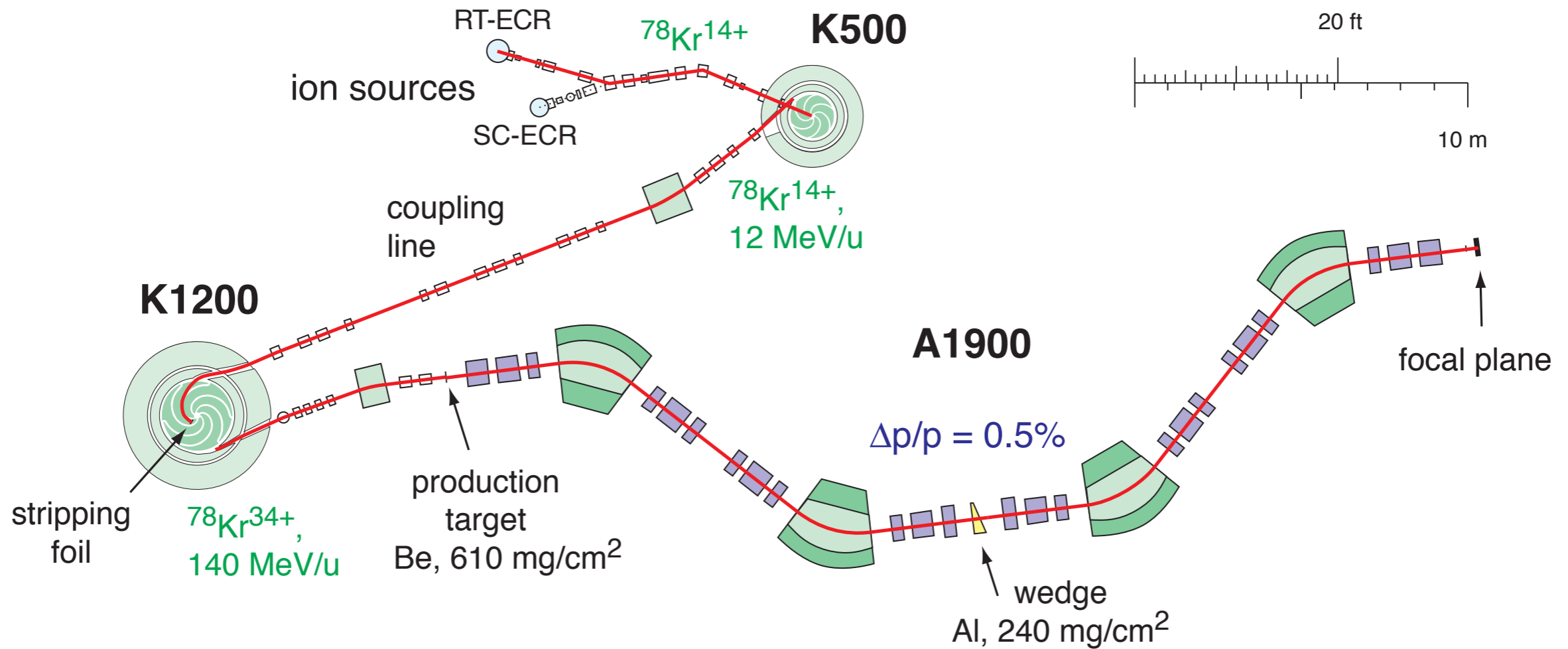
A/Z filter

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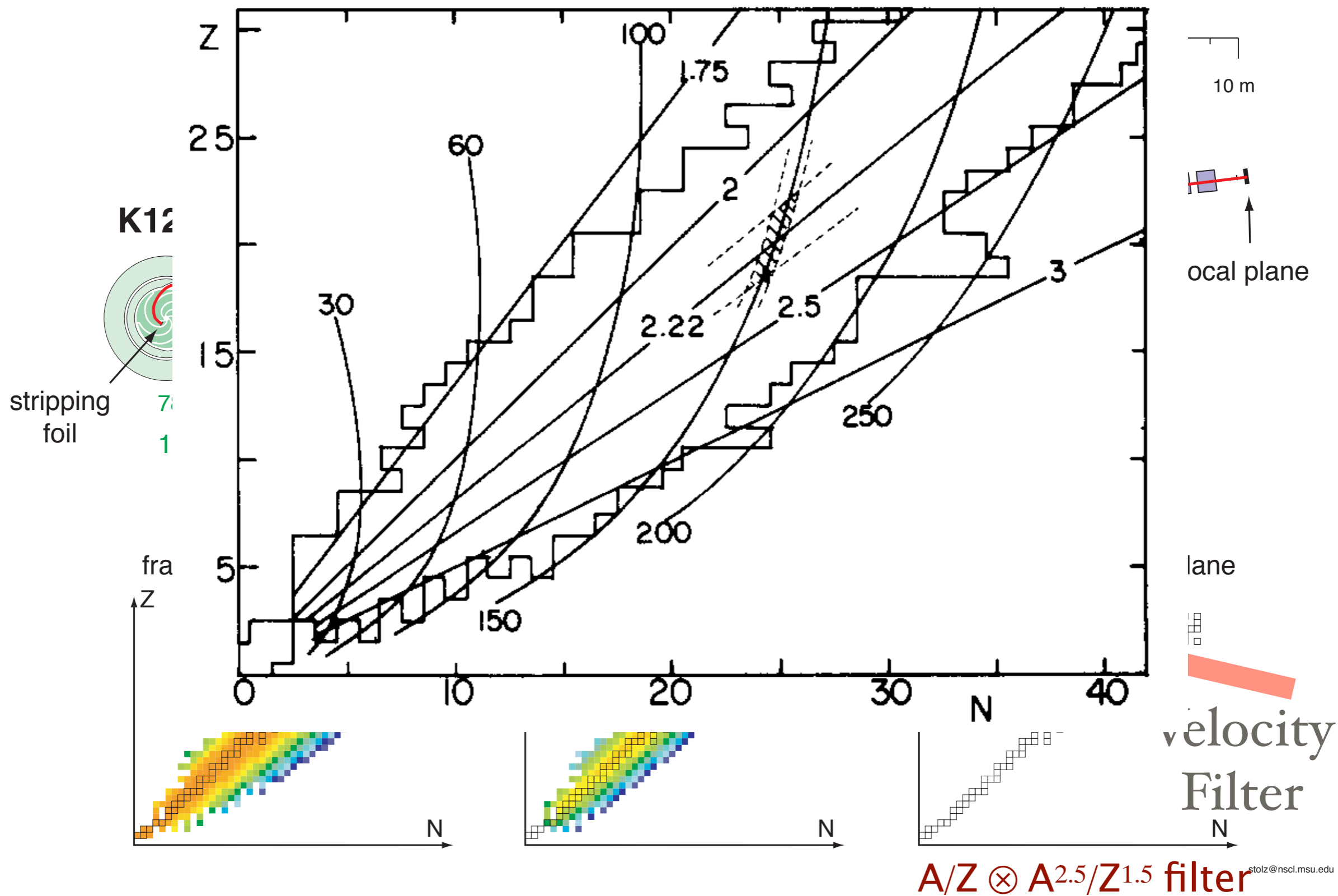


$A/Z \otimes A^{2.5}/Z^{1.5}$ filter

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 - ▶ Dynamical and statistical properties, but also (and in fact mostly) static properties

Nuclear reactions

Impact
Parameter



A vertical double-headed arrow is positioned to the right of the text 'Impact Parameter', indicating the vertical distance between the centers of the two overlapping spheres.



Nuclear reactions

- ▶ “Simple” or “Direct” reactions
 - ▶ Usually peripheral collision (large impact parameter)
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 - ▶ Probe single-particle properties of nuclei



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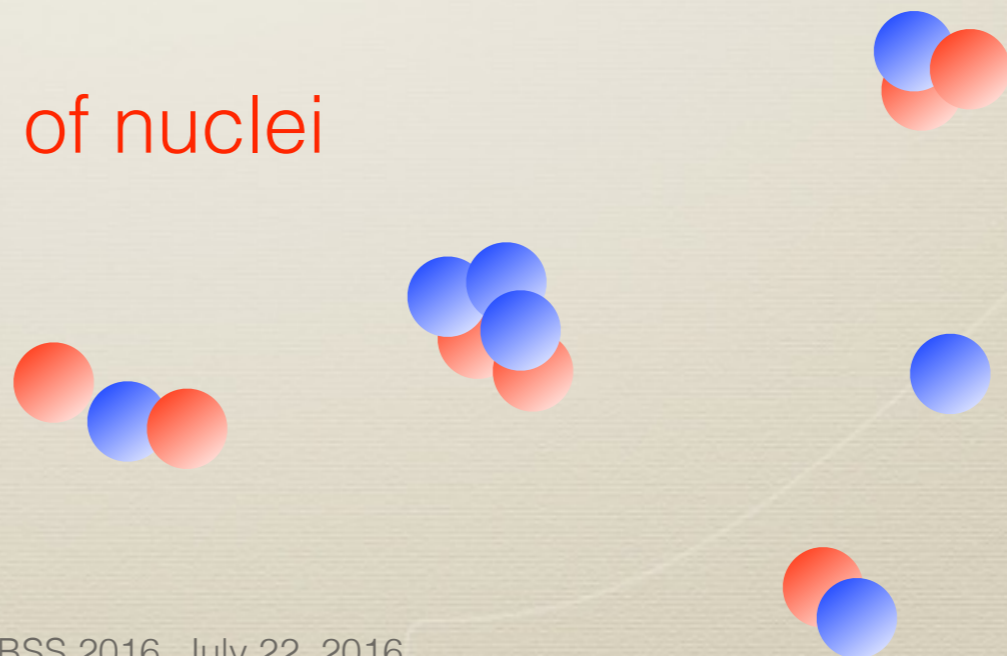
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Impact
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Direct reactions are tools

- ▶ ... to study Nuclear Structure and Astrophysics

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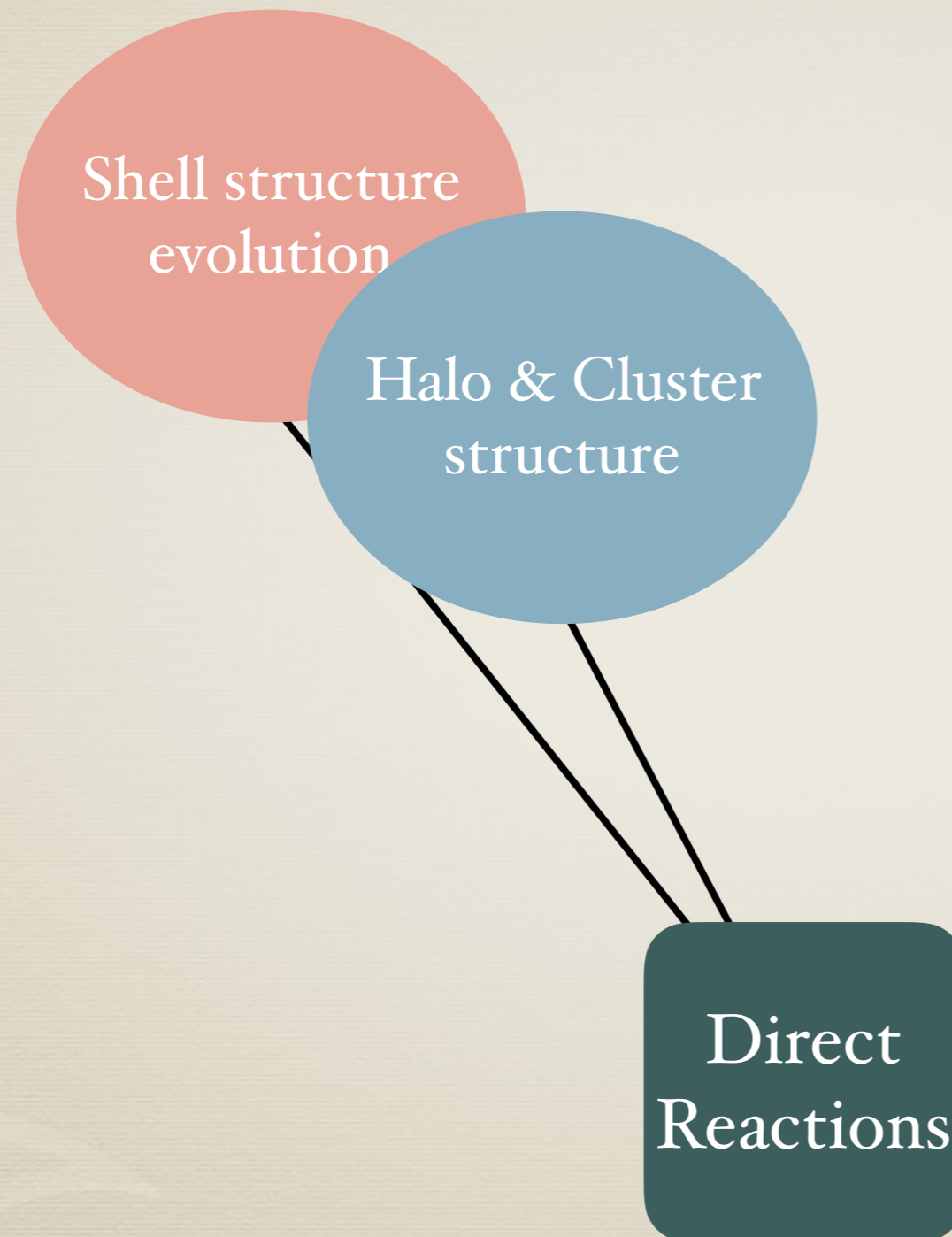
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Shell structure
evolution

Direct
Reactions

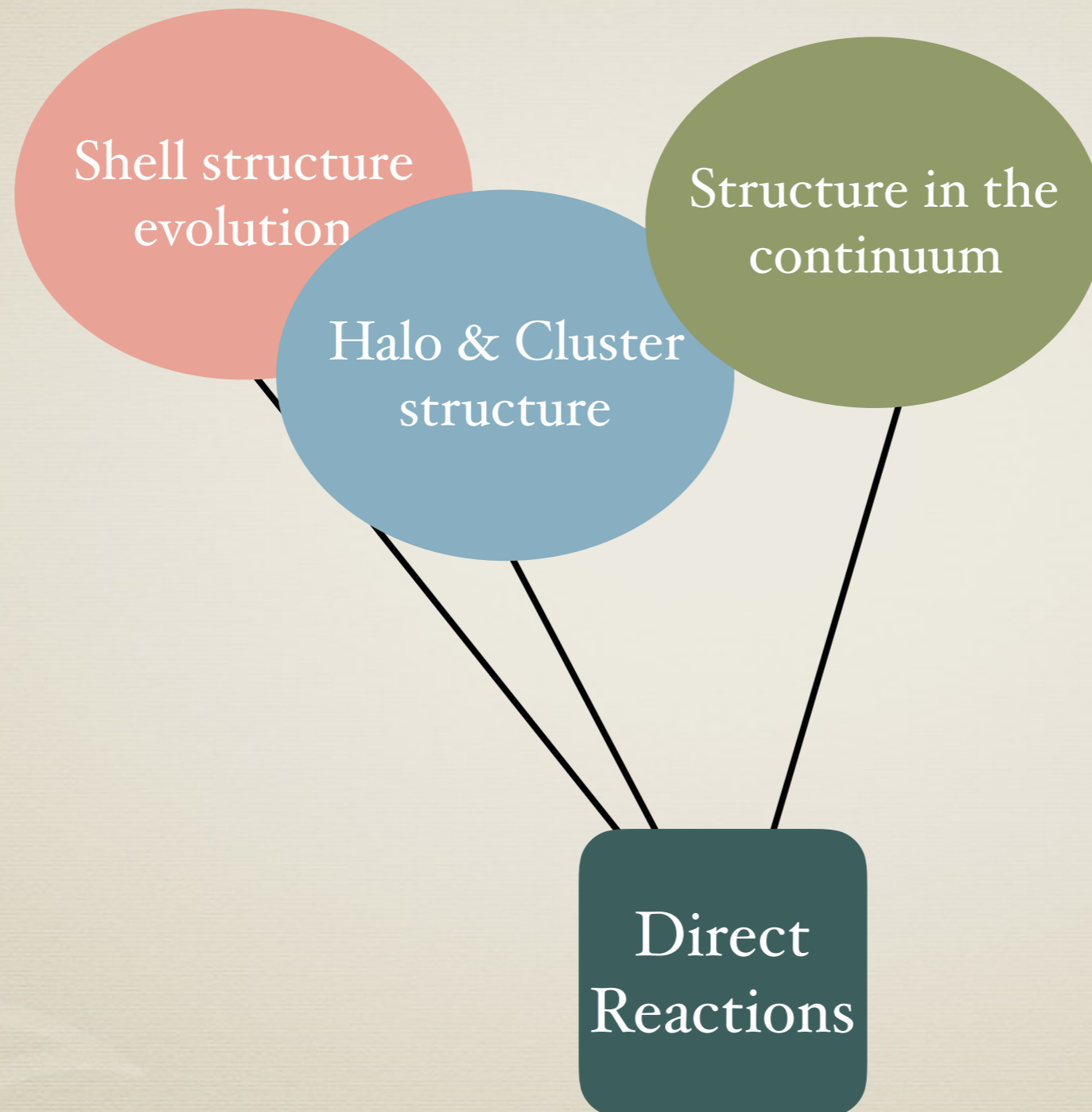
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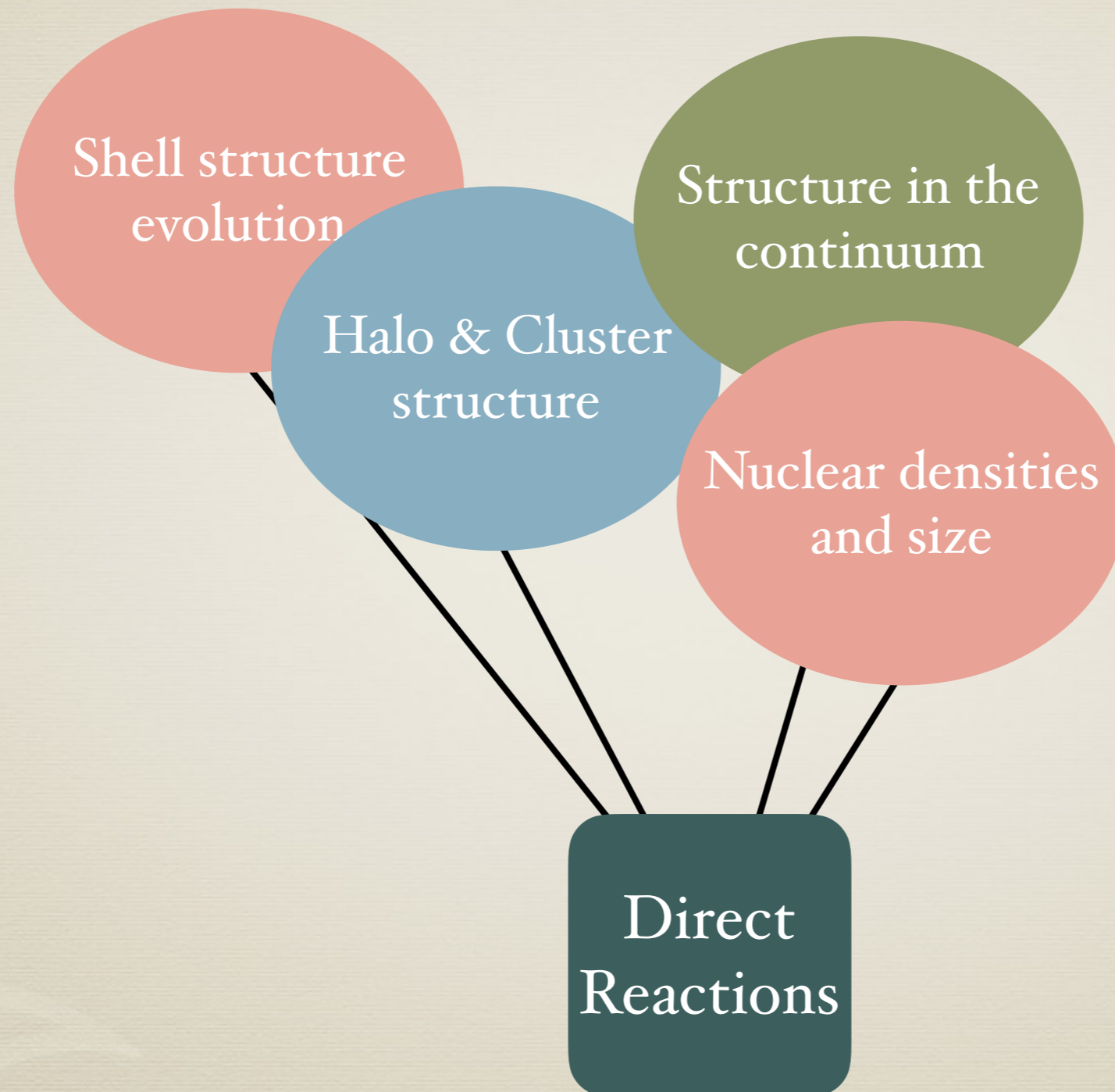
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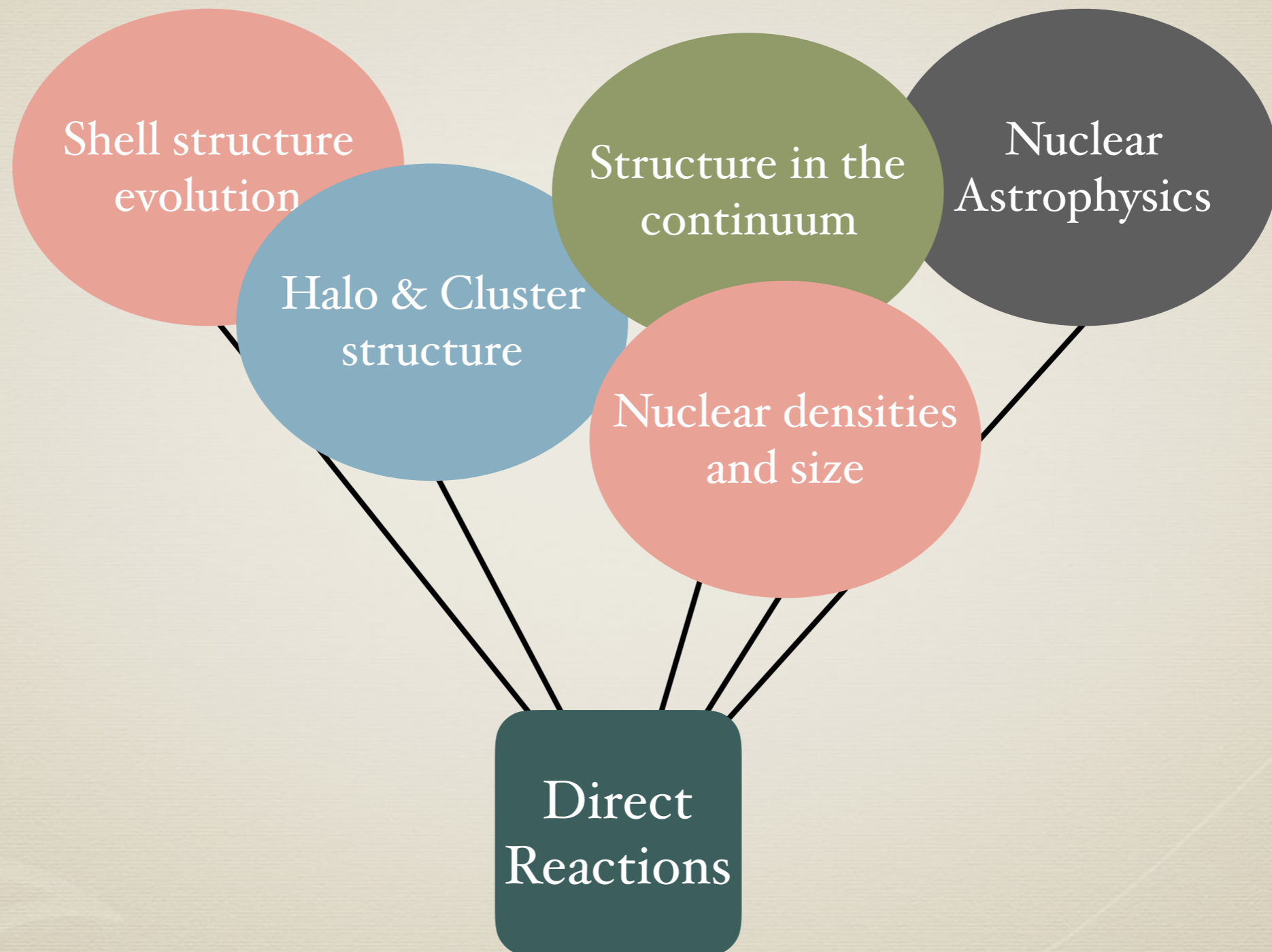
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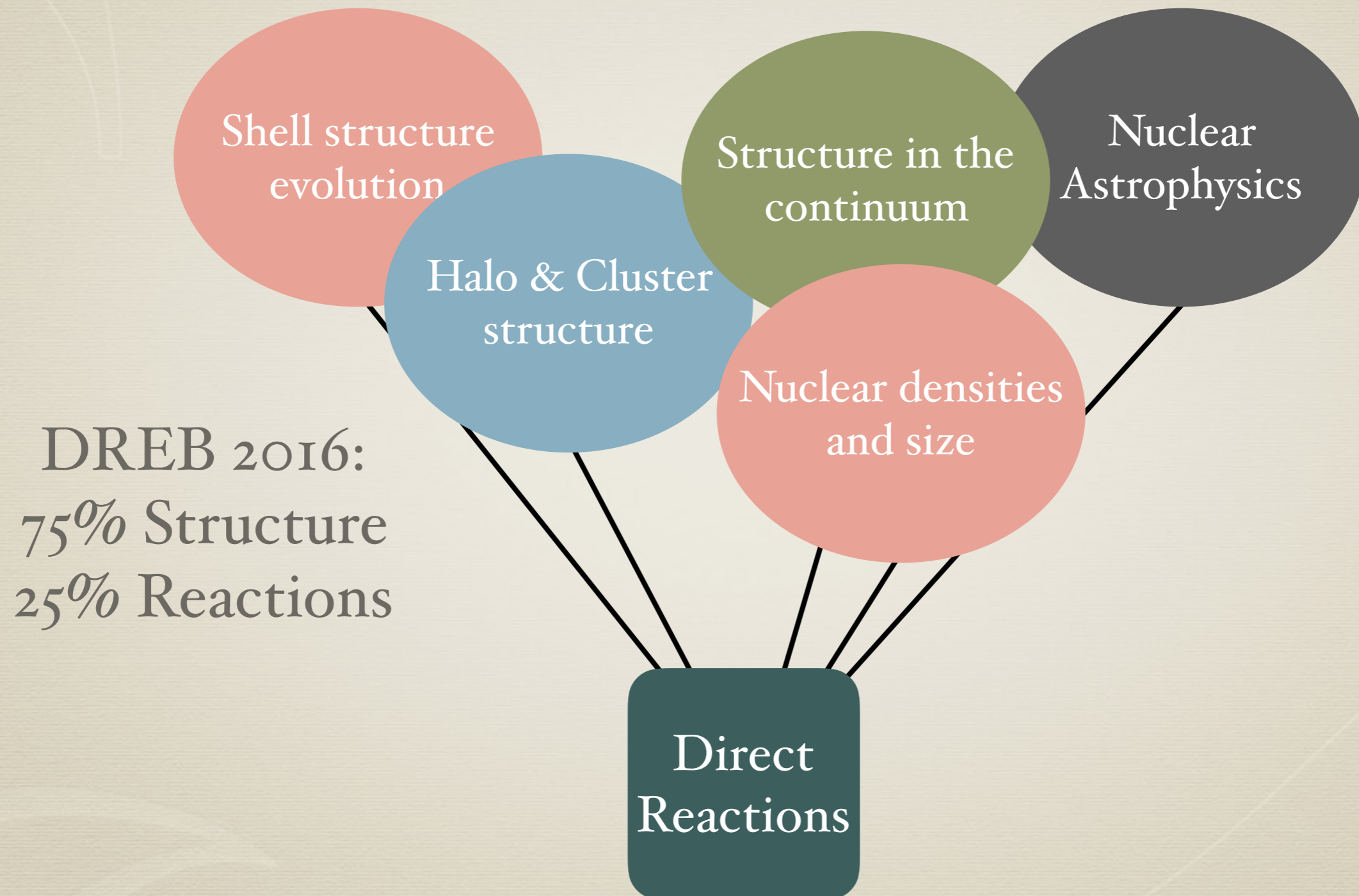
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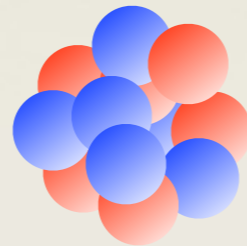
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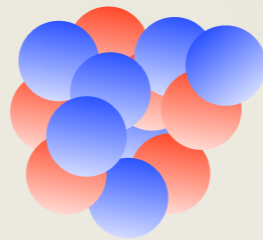
Not so simple...

- ▶ Pictures and words are necessary, but they confine and restrict the imagination



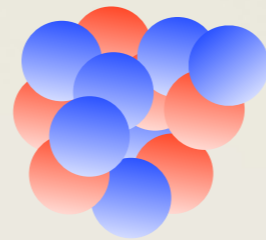
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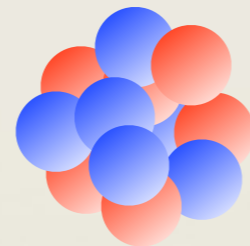
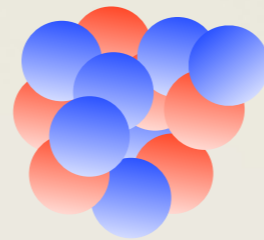
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Angular momentum
transfer

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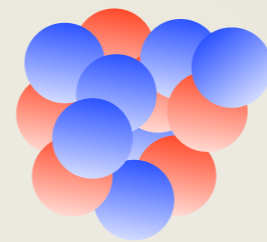
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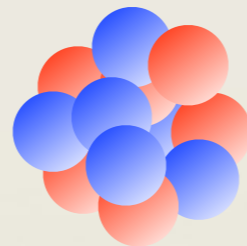
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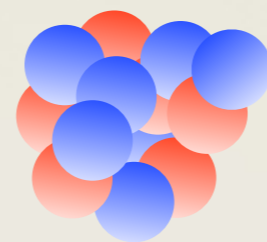


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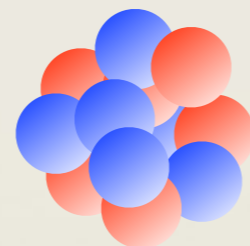


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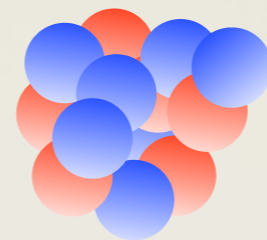


Elastic scattering below the Coulomb barrier

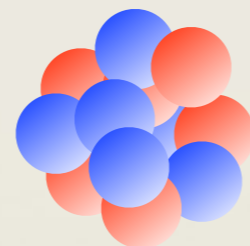


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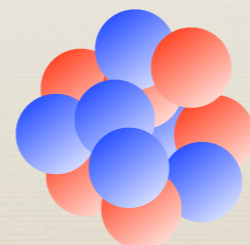
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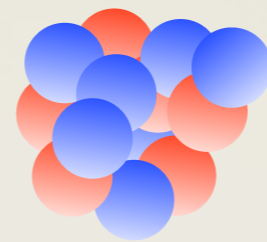


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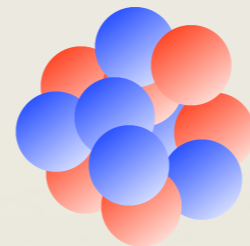


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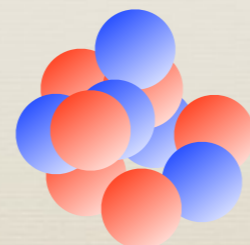
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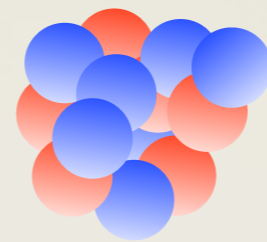


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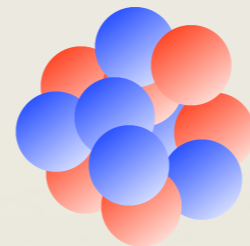


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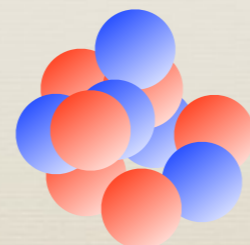
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Quasi-free (p,pn) scattering



A large palette of tools

- ▶ Covering a very wide range of beam energies
 - ▶ From < 1 MeV/u up to > 1 GeV/u



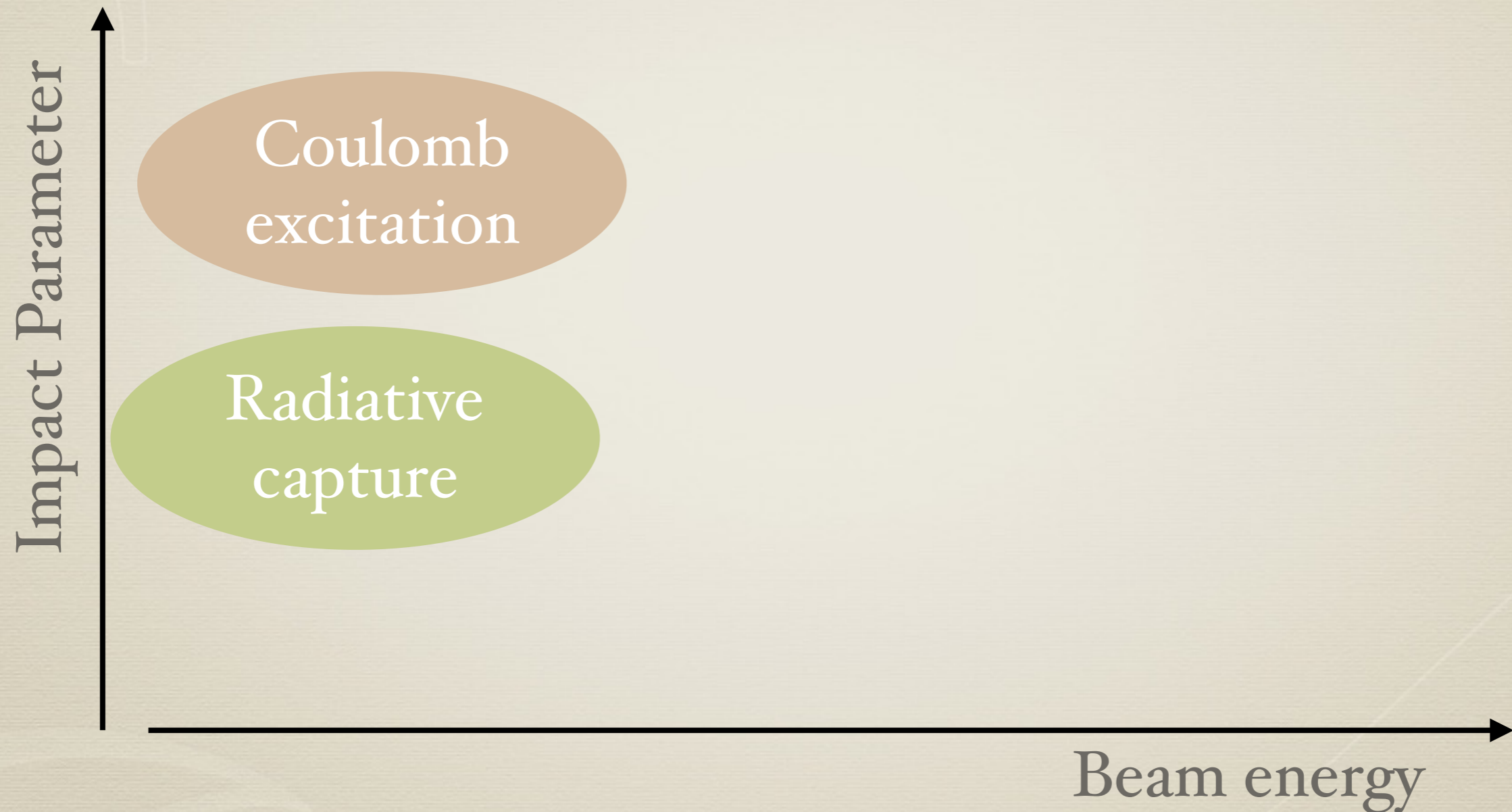
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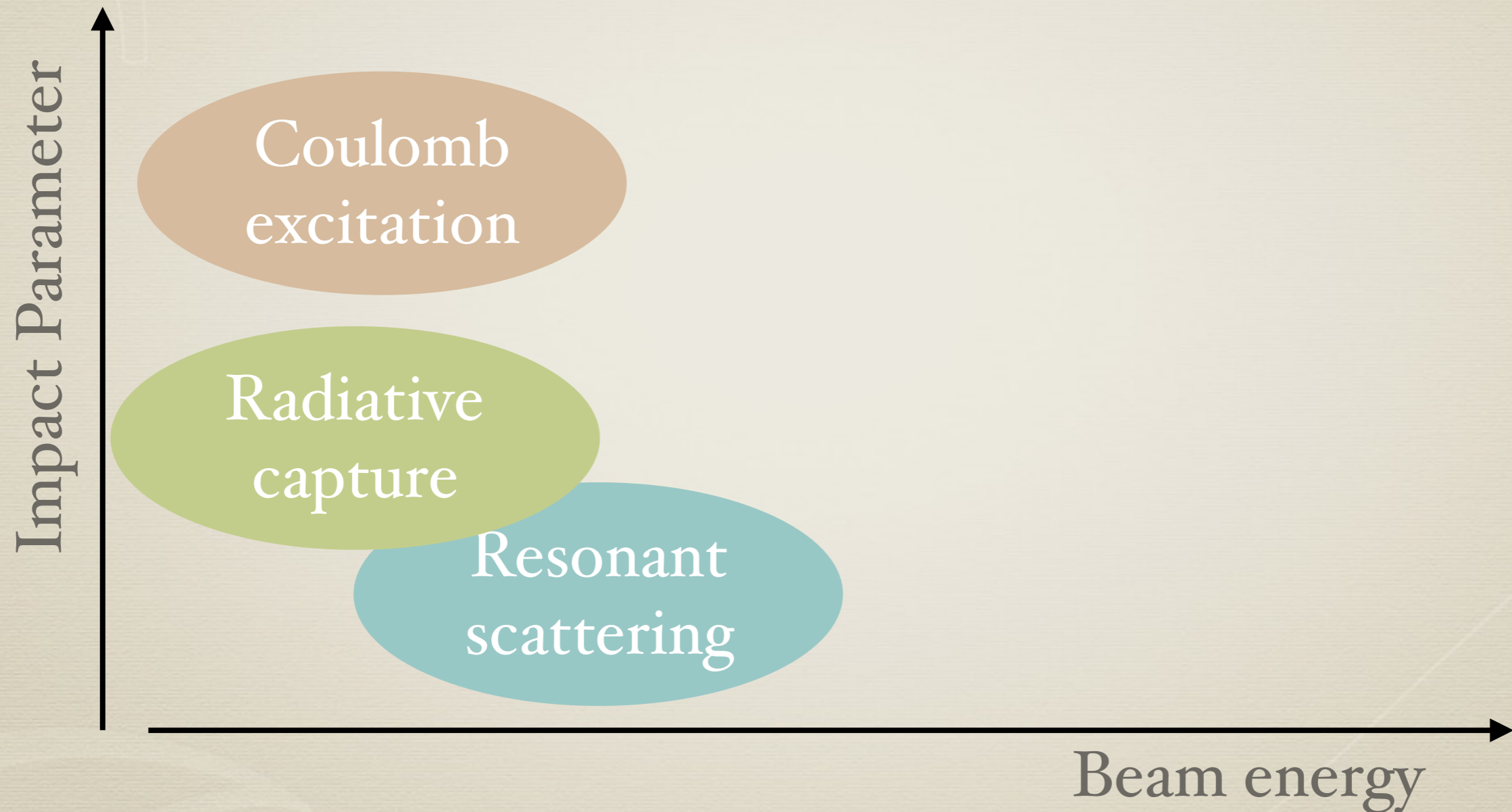
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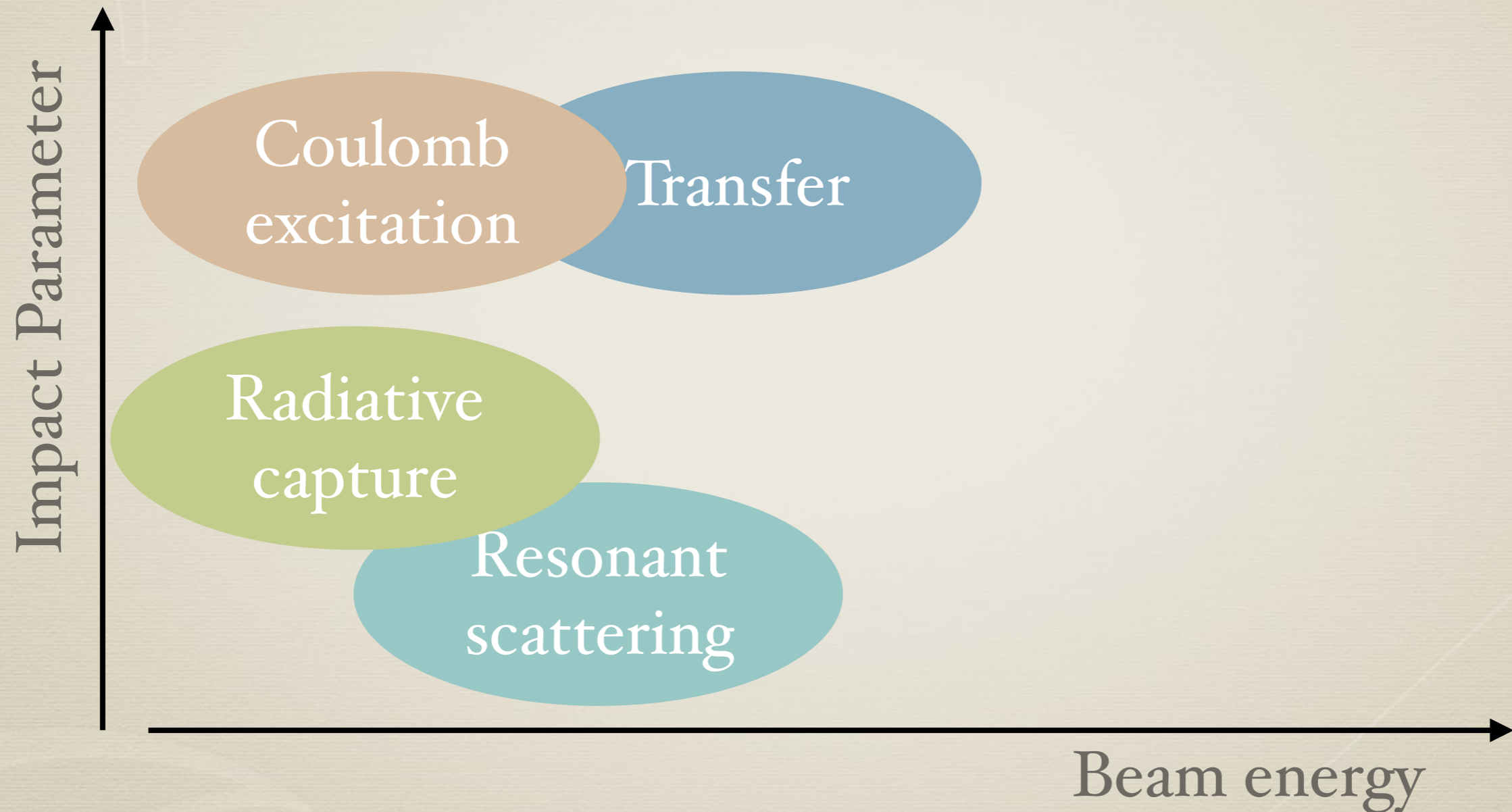
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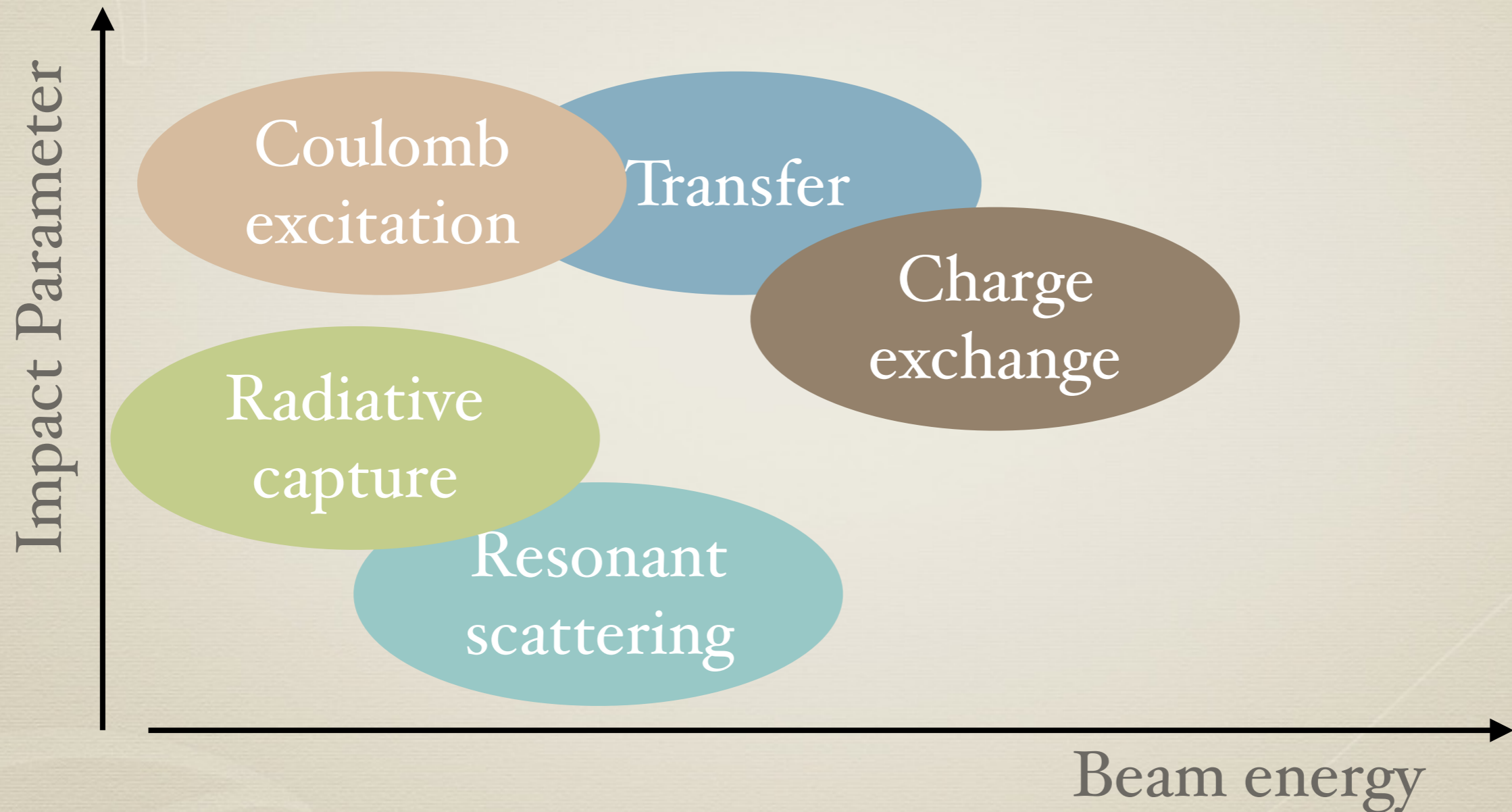
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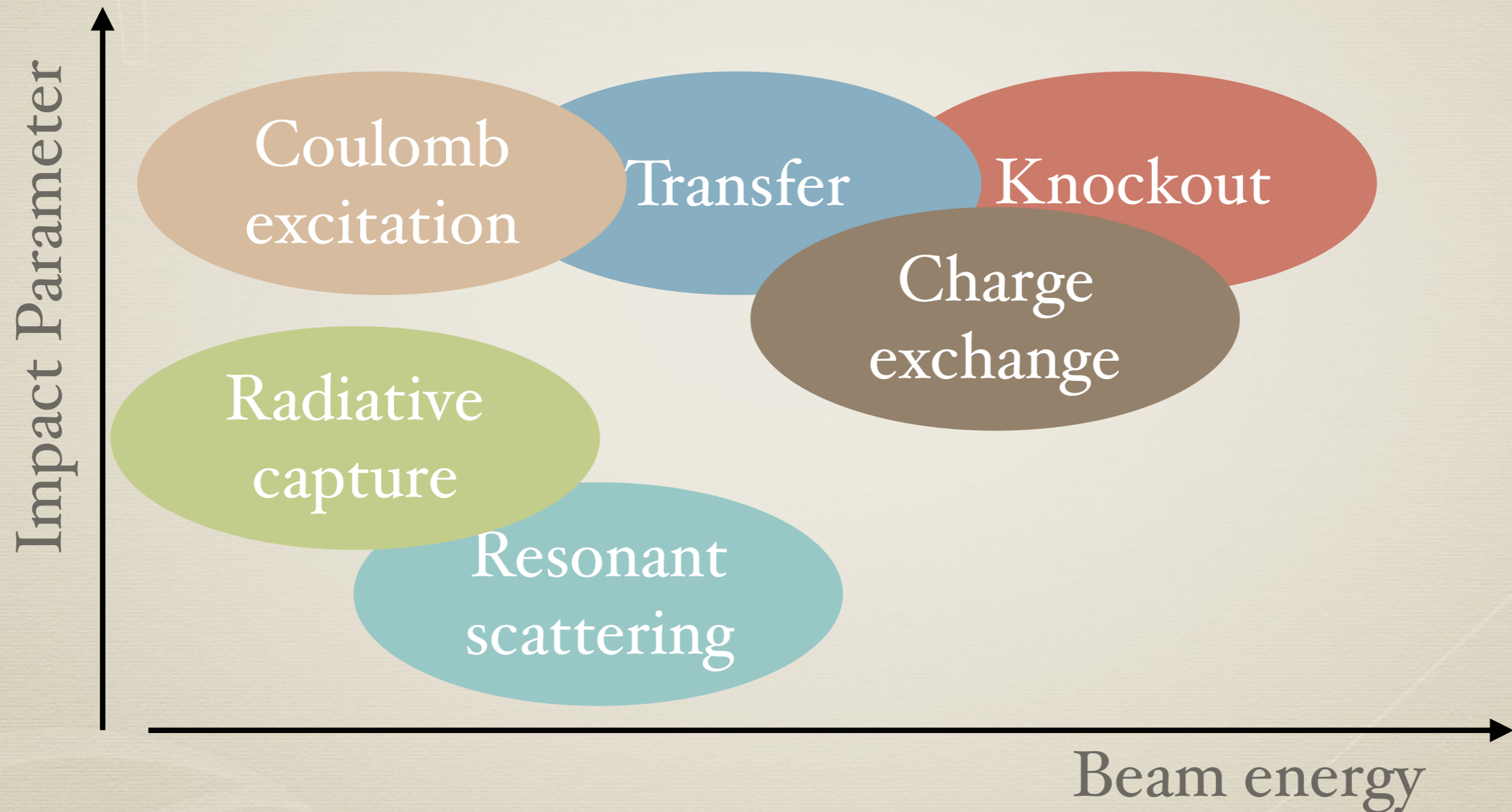
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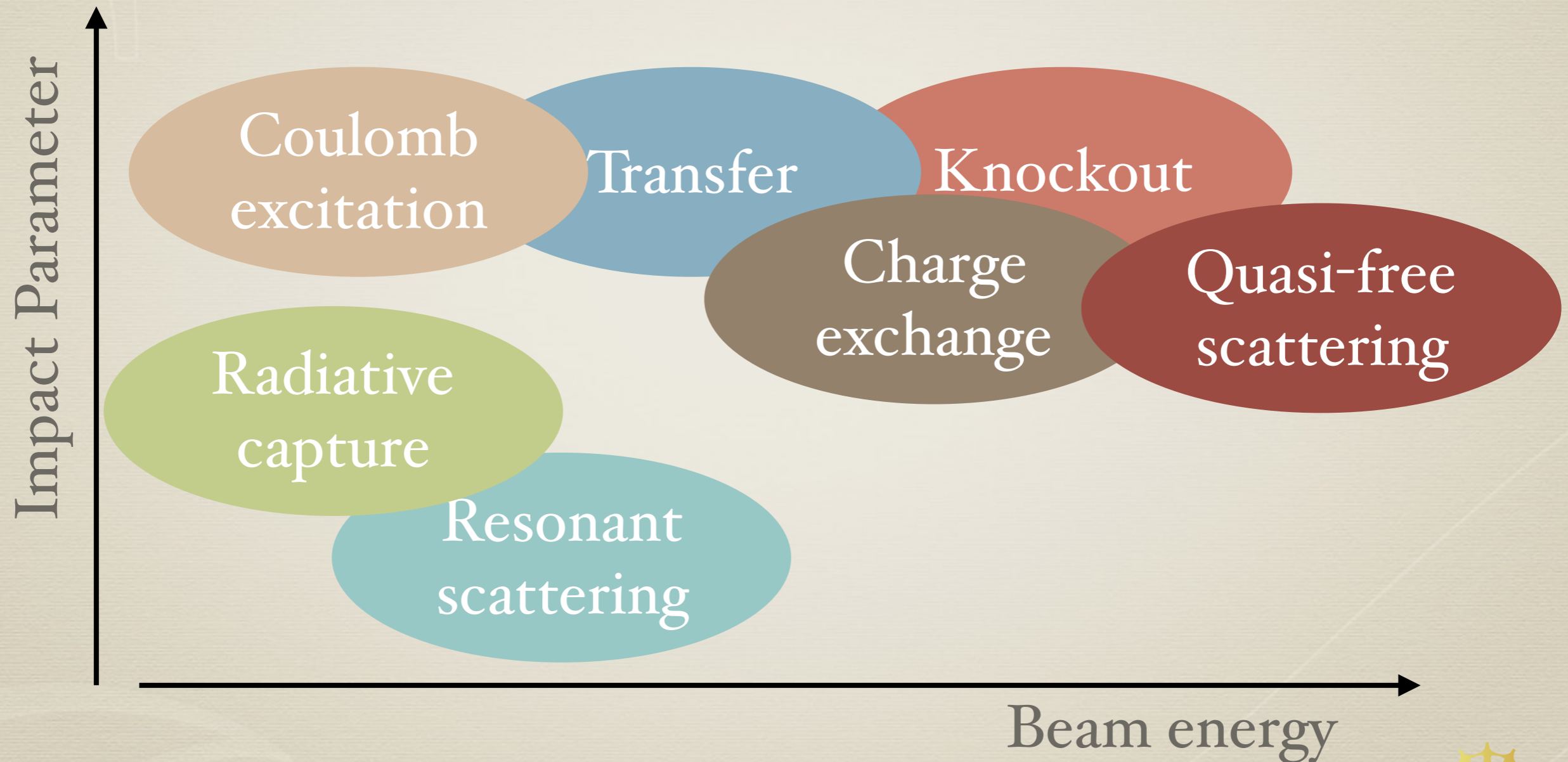
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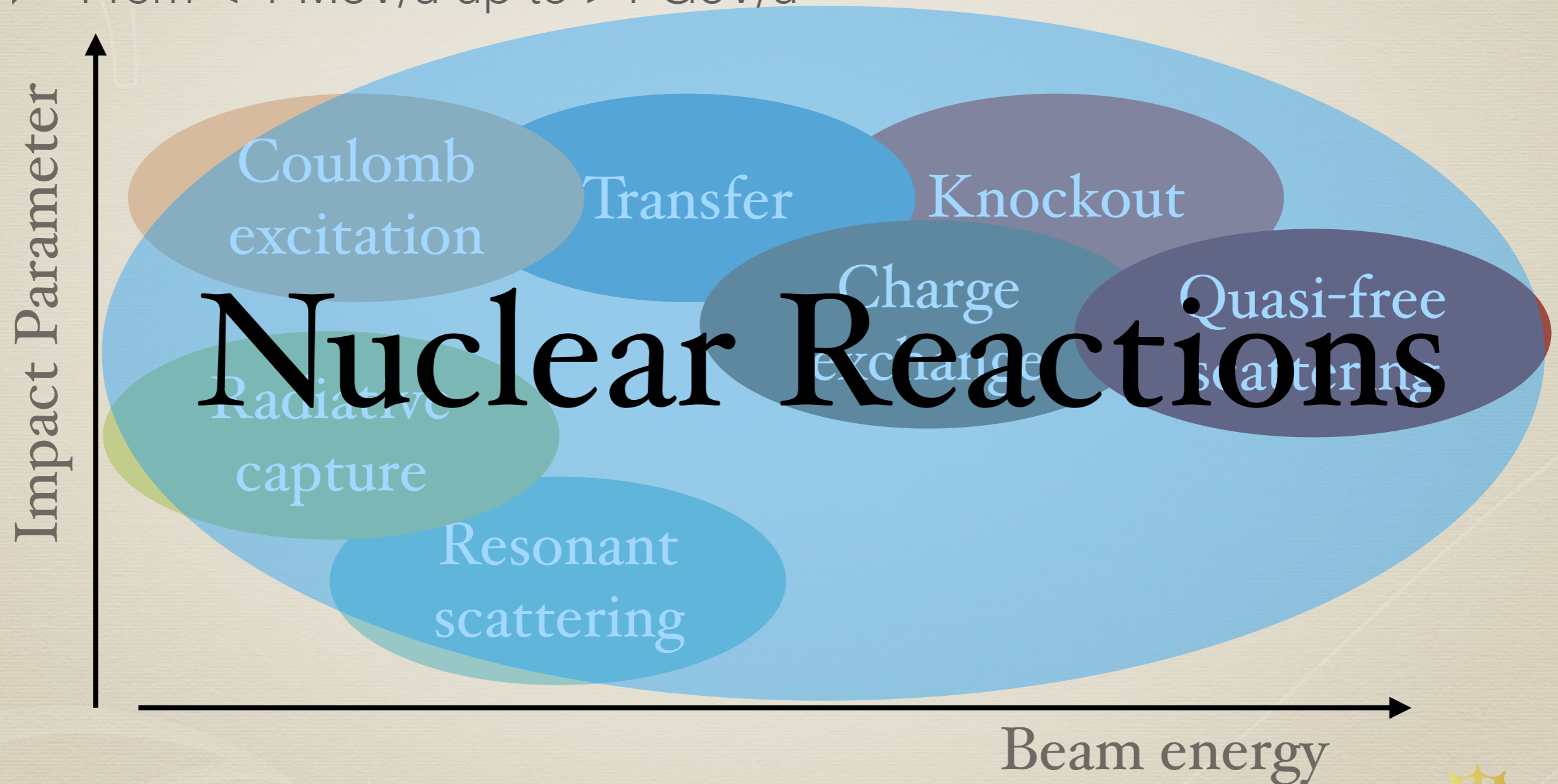
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Nuclear Physics Conundrum

- ▶ Using “simple” nuclear reaction to study nuclear structure

Observable:
cross section

Structure model:
spectroscopic factor

Reaction model:
single-particle
cross section

$$\sigma^{if} = \sum_{|J_f - J_i| \leq j \leq J_f + J_i} S_j^{if} \sigma_{sp}$$

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- ▶ We want to extract valuable structure information
- ▶ Challenge: need valid reaction model!

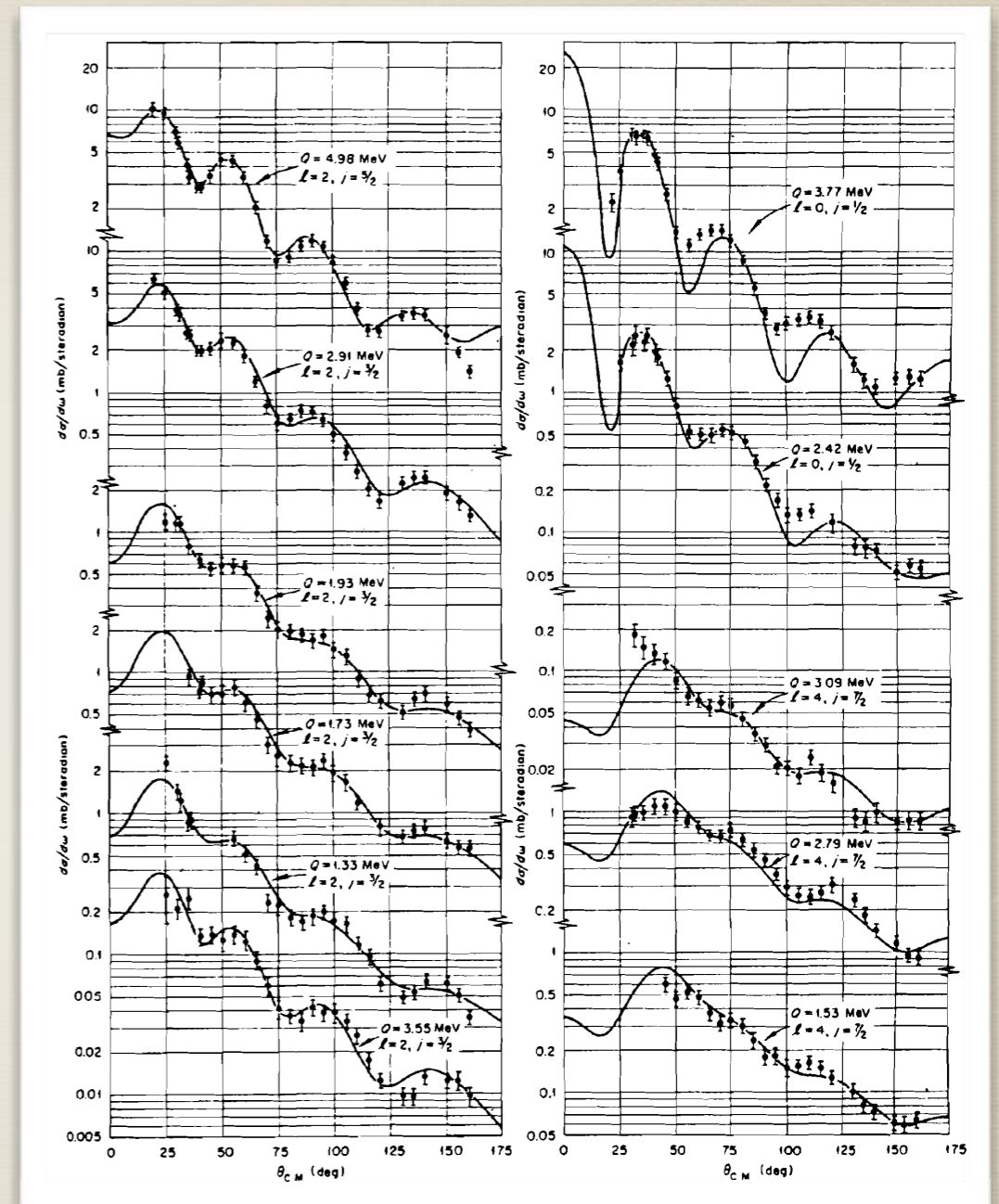
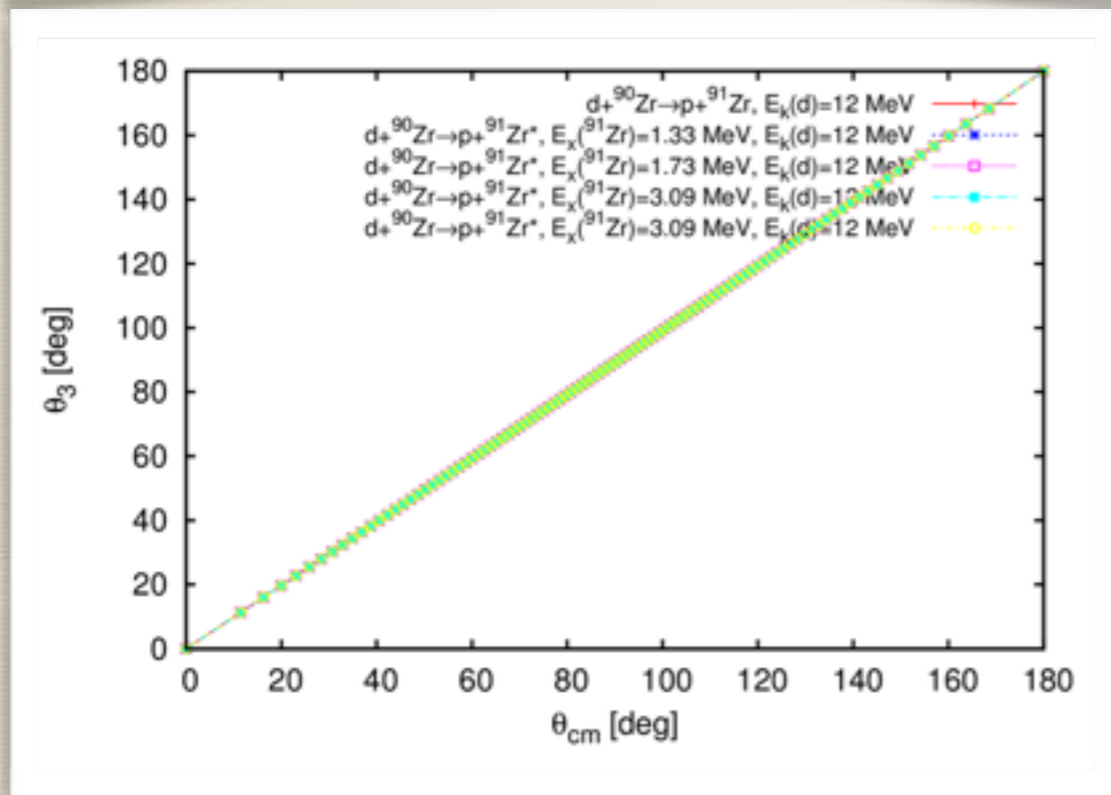
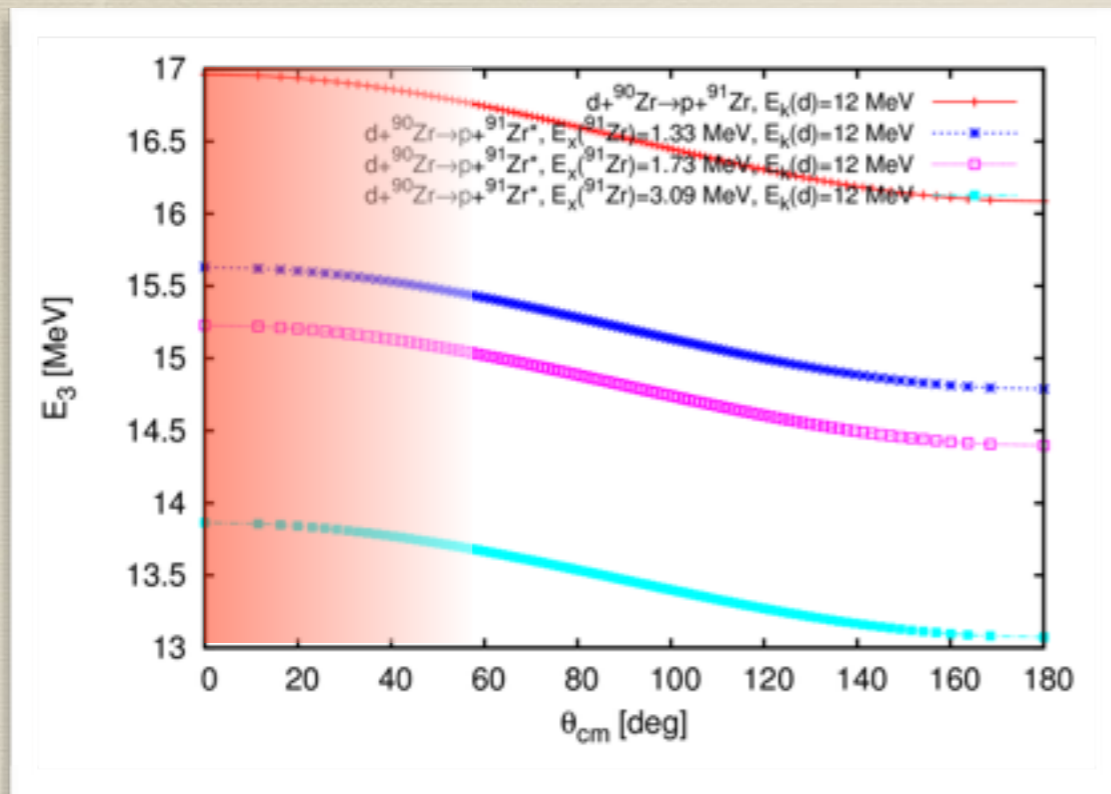
The good early days...

- ▶ Use light nuclei (p, d, ^4He) on stable targets
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- ▶ Use direct reactions to study static properties of nuclei
 - ▶ Populate levels in heavy residue → level scheme
 - ▶ Measure energy and angular distribution of scattered particle → spins and parities of states
 - ▶ Measure cross sections and compare to reaction calculations → spectroscopic factors

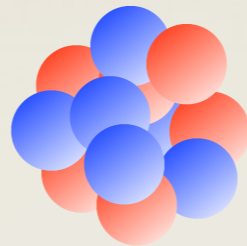
The $^{90}\text{Zr}(d,p)^{91}\text{Zr}$ reaction



$^{90}\text{Zr}(d,p)^{91}\text{Zr}$ @ 6 MeV/u to 1g, 2d, 3s orbitals
 H. Feshbach, Nuclear Reactions, Wiley & Sons
 (1992)

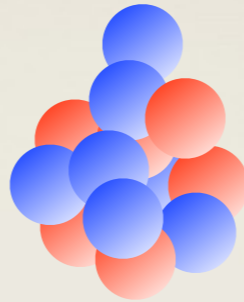
Issue n°1: inverse kinematics!

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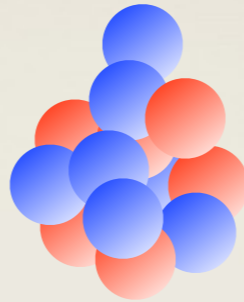
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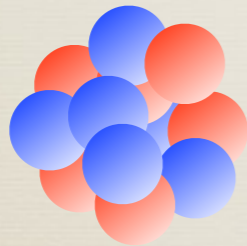


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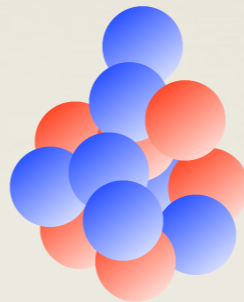


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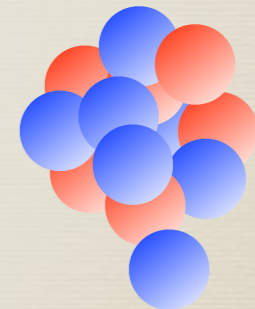


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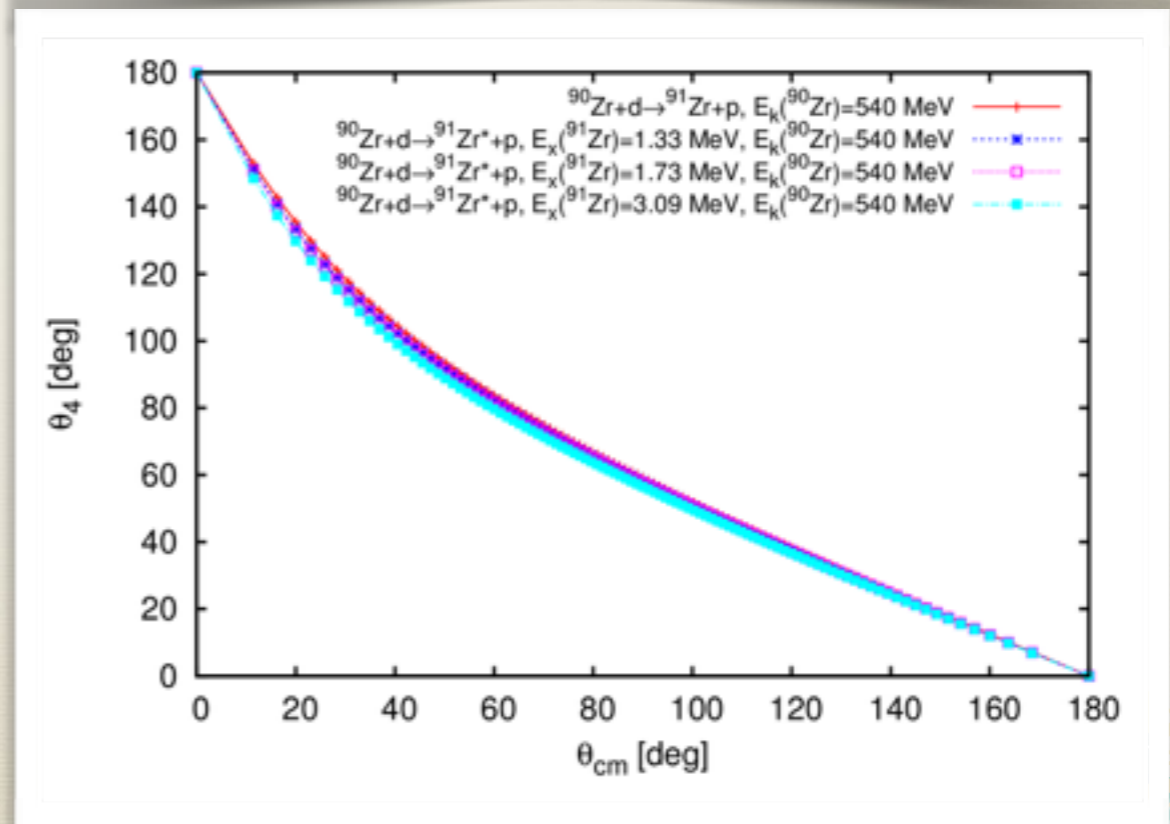
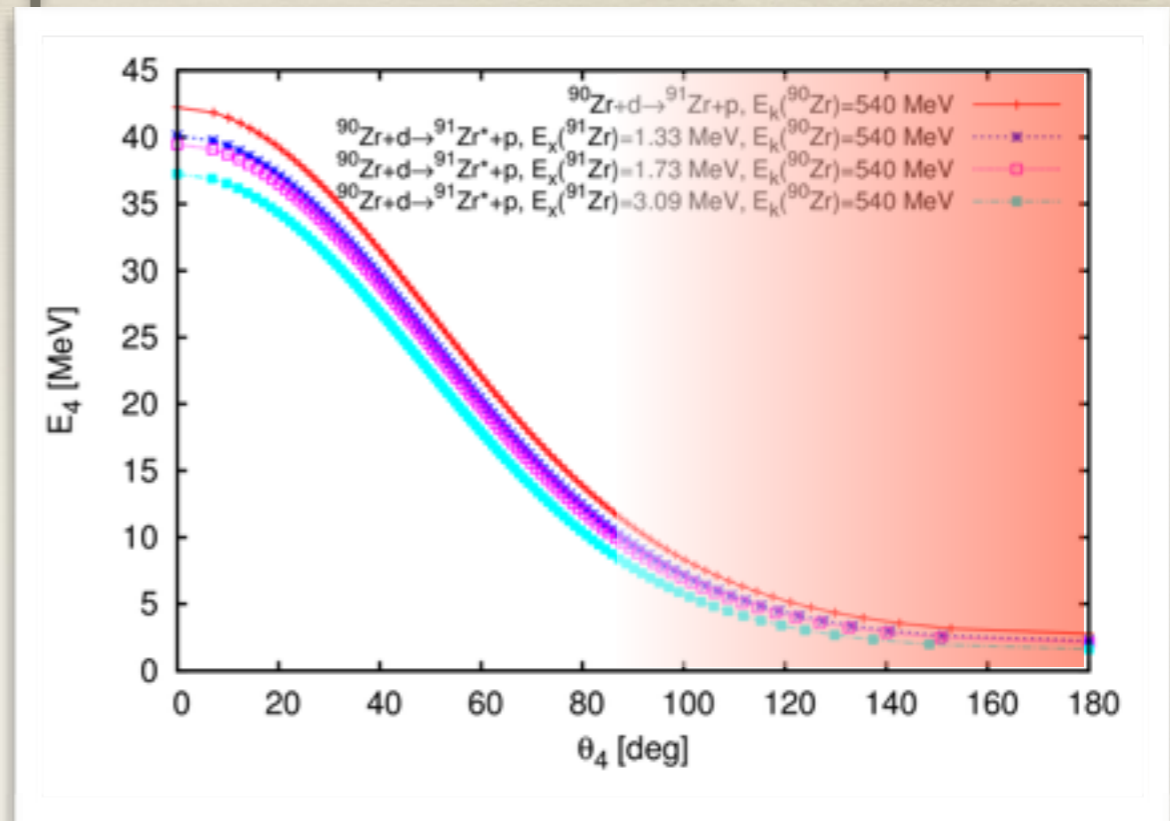
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Backwards!

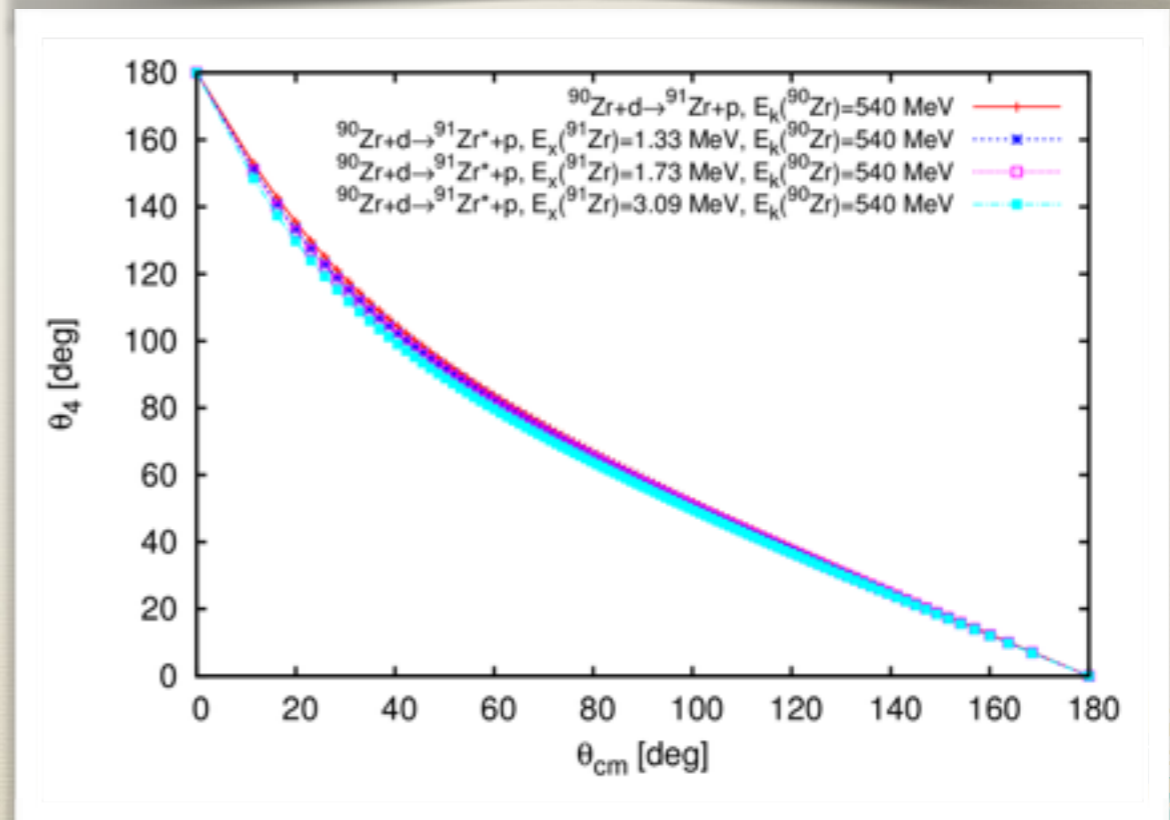
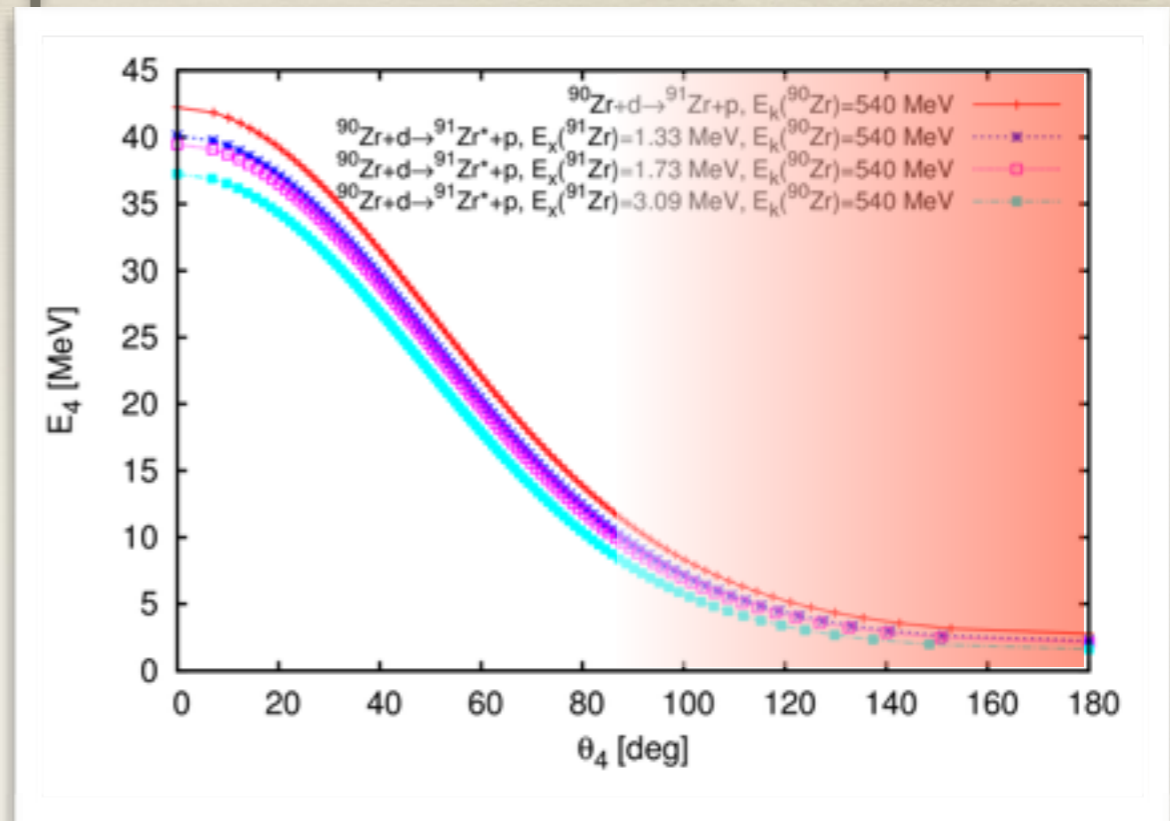
$d(^{90}\text{Zr}, ^{91}\text{Zr})p$ reaction

► What happened!



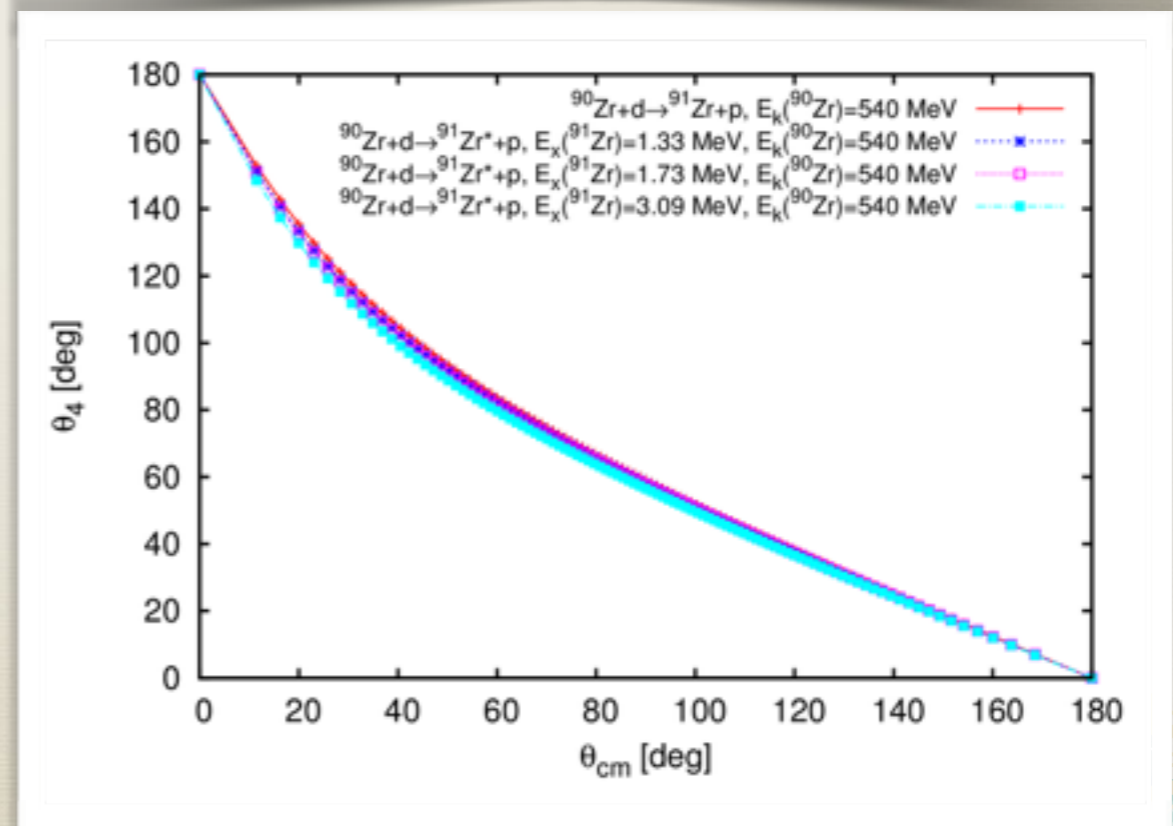
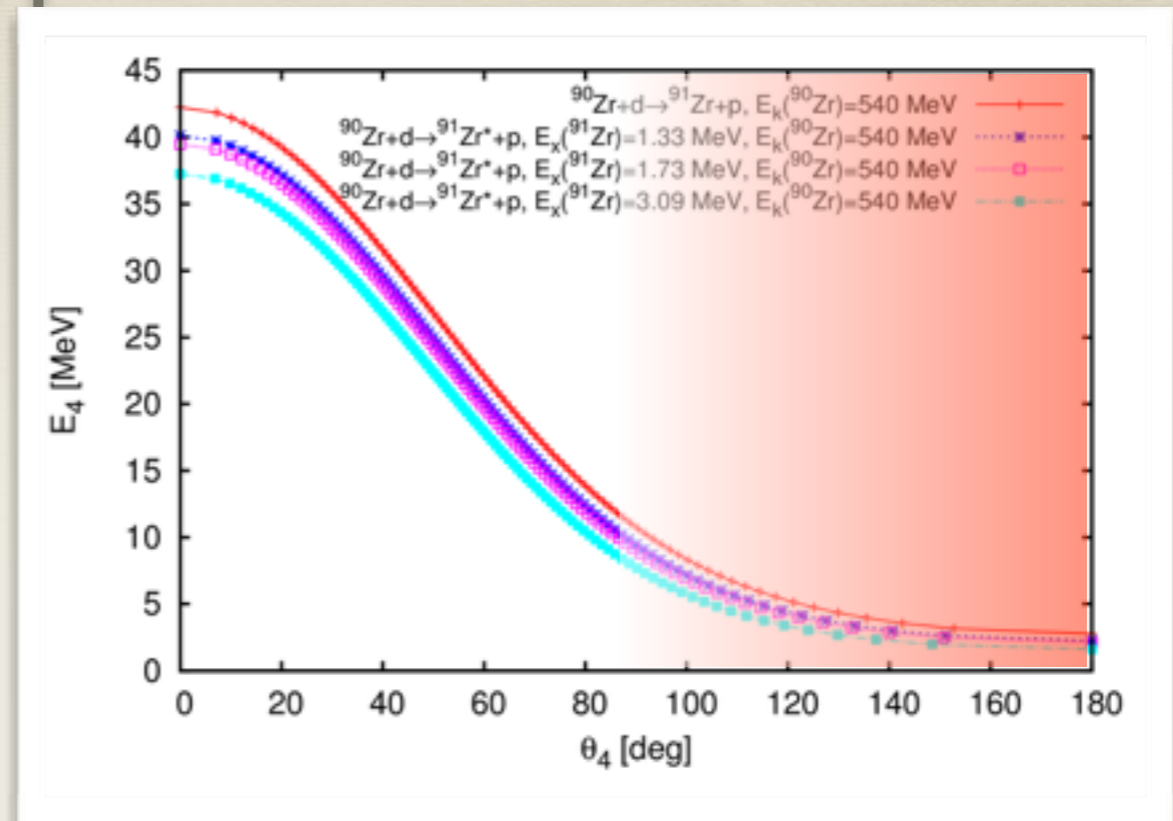
$d(^{90}\text{Zr}, ^{91}\text{Zr})p$ reaction

- ▶ What happened!
- ▶ Center-of-mass motion now with projectile, not target



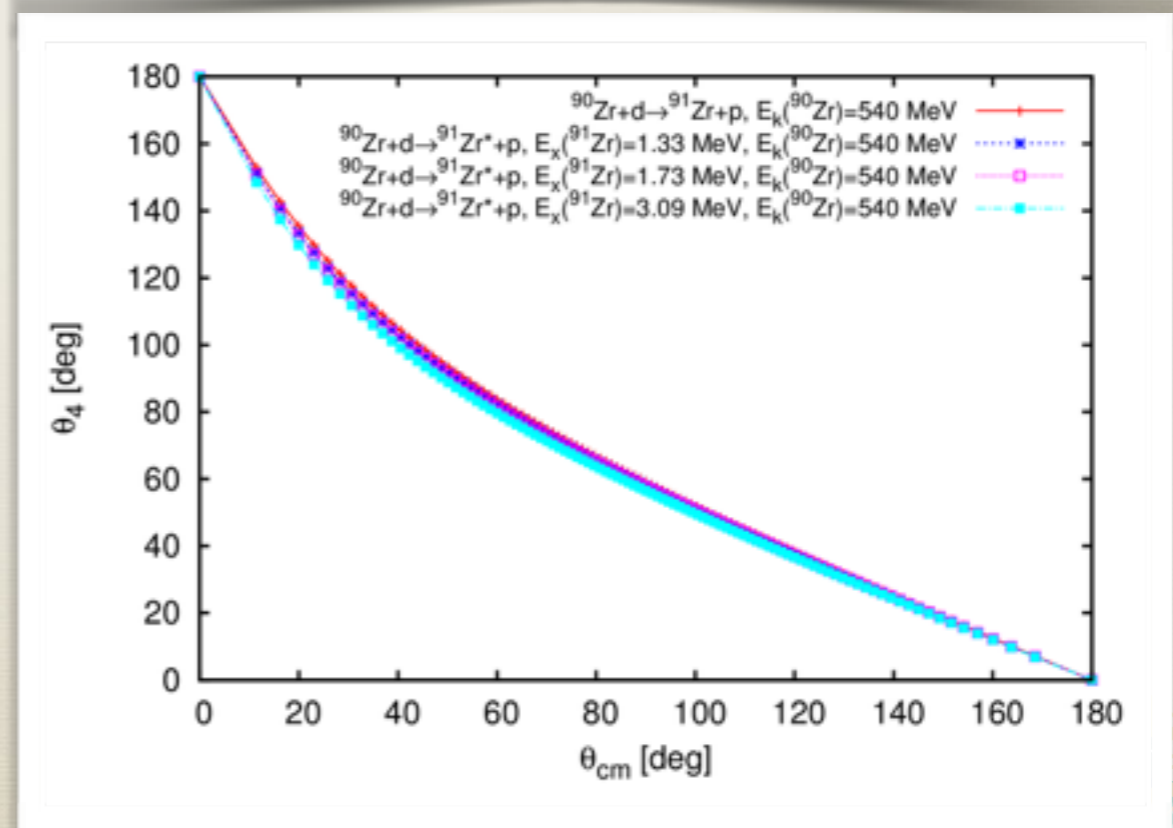
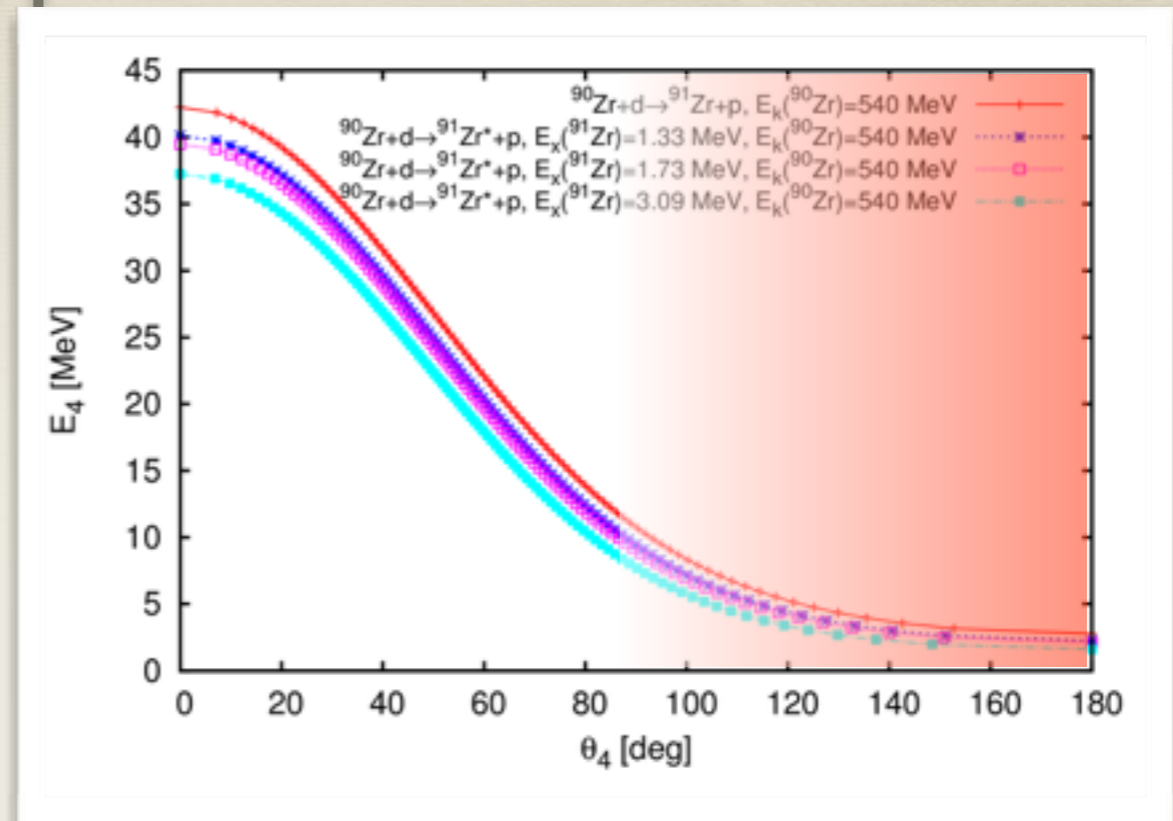
$d(^{90}\text{Zr}, ^{91}\text{Zr})p$ reaction

- ▶ **What happened!**
 - ▶ Center-of-mass motion now with projectile, not target
 - ▶ Scattered protons now have very wide range of energies



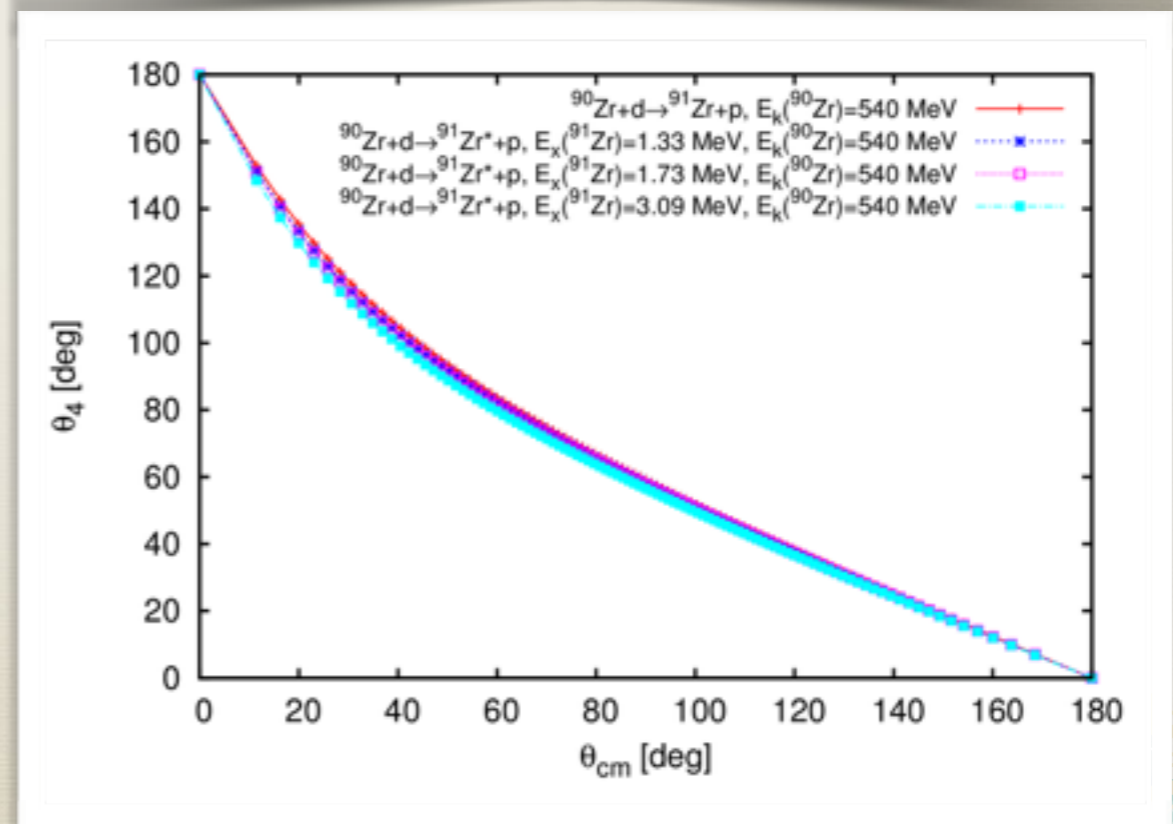
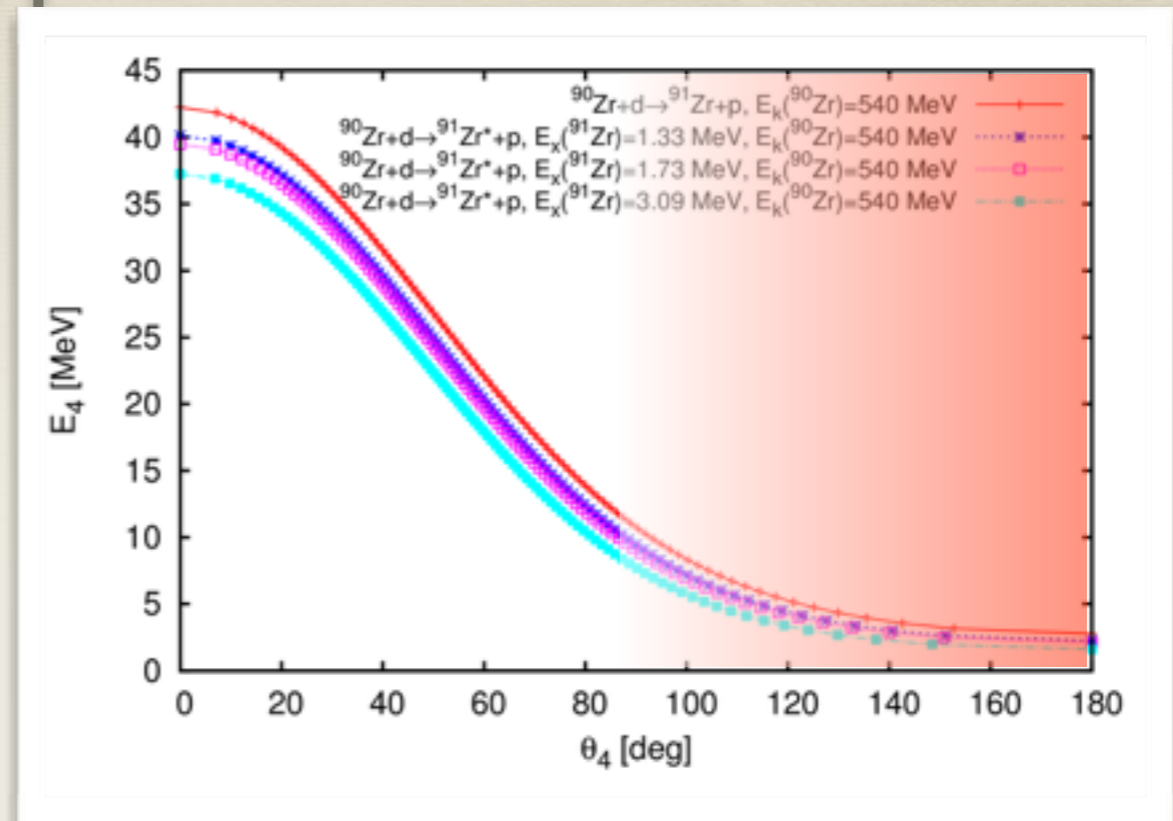
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 - ▶ Worse: largest cross sections at the lowest end of energy range



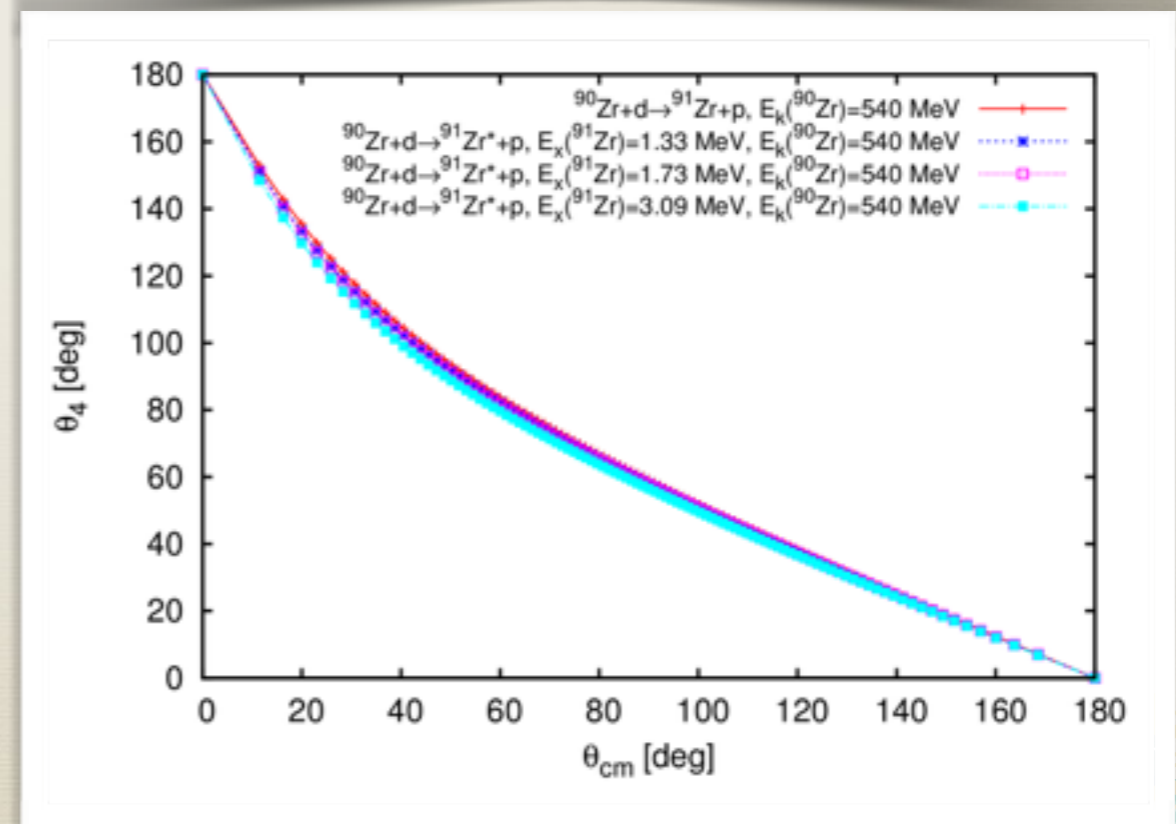
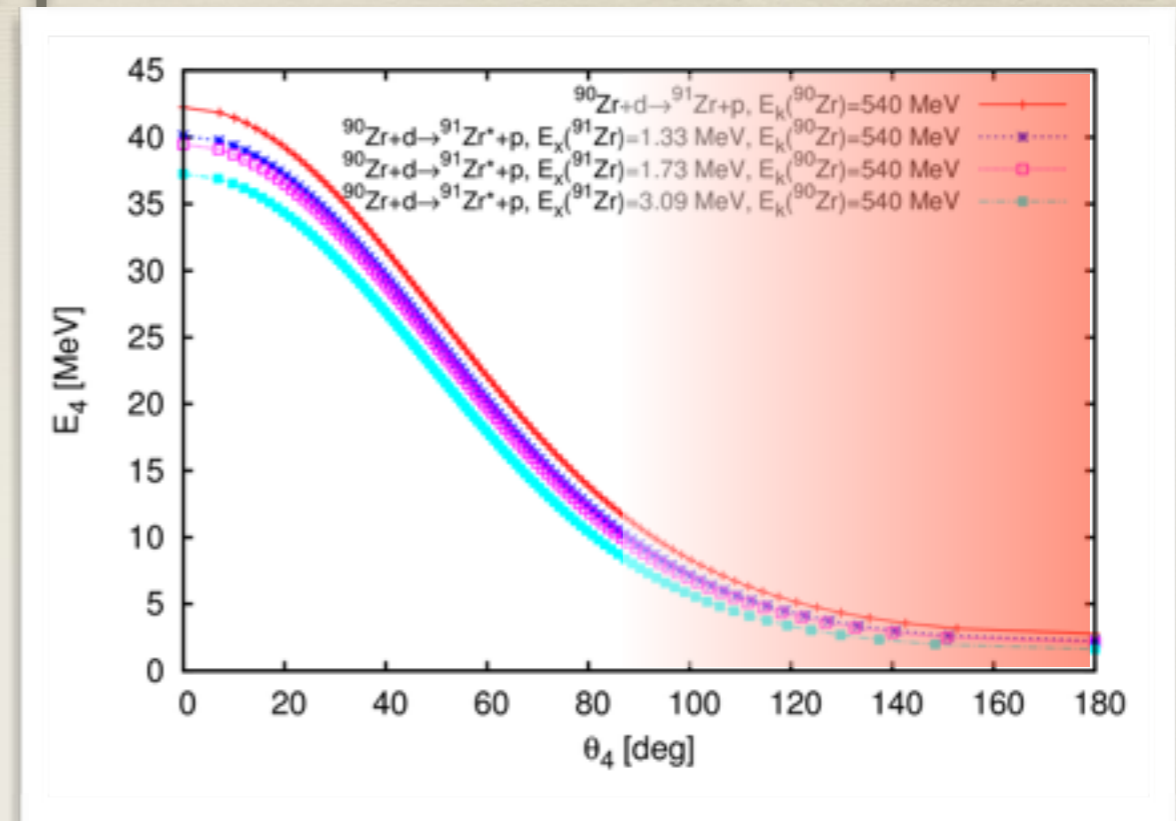
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 - ▶ More difficult for proton to emerge from target



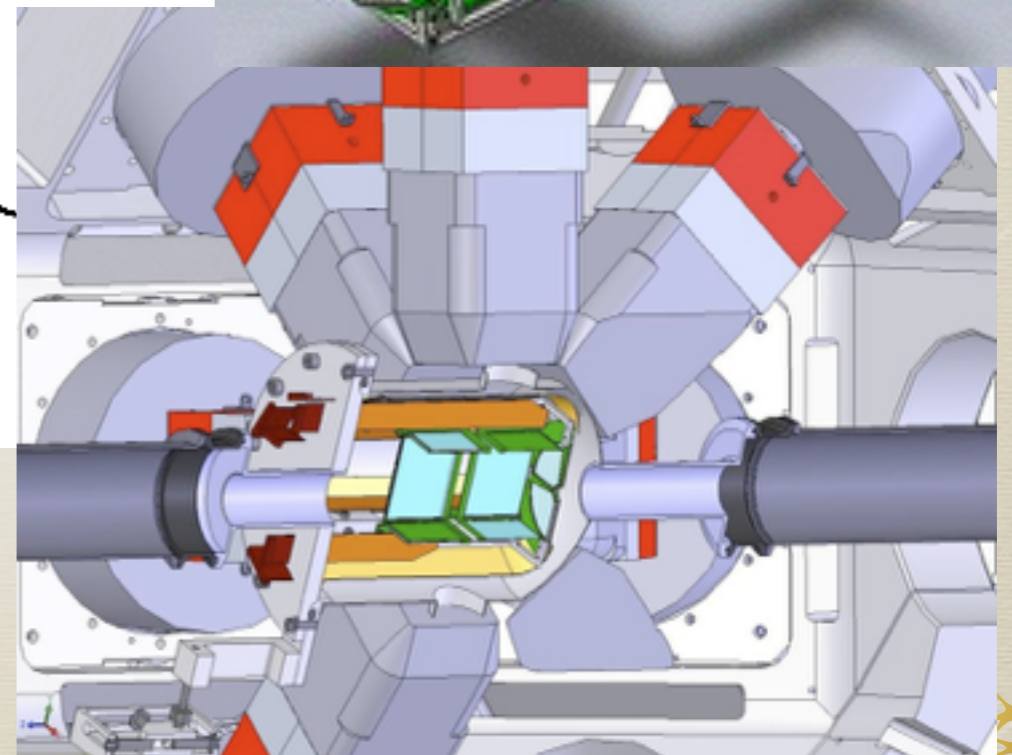
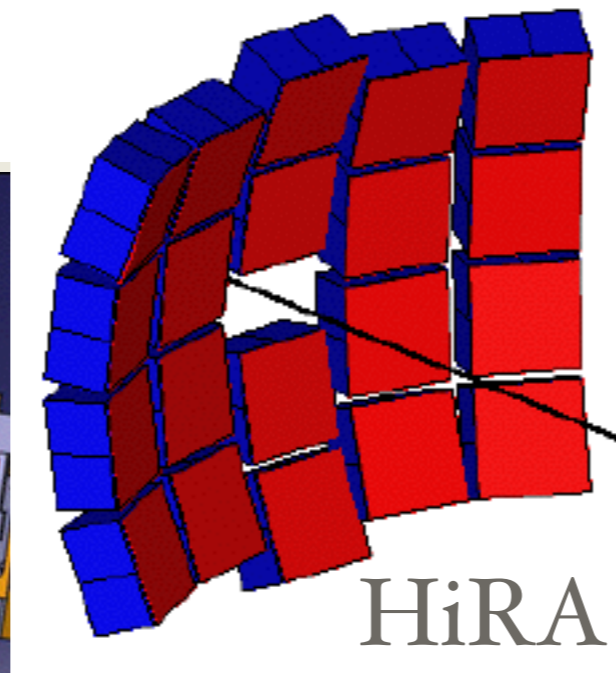
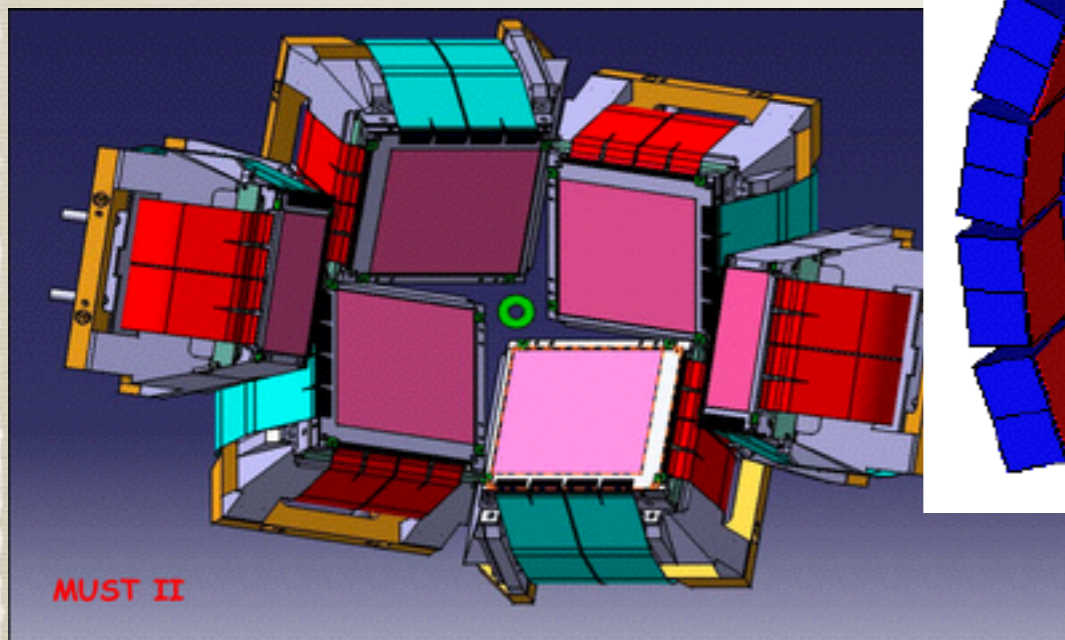
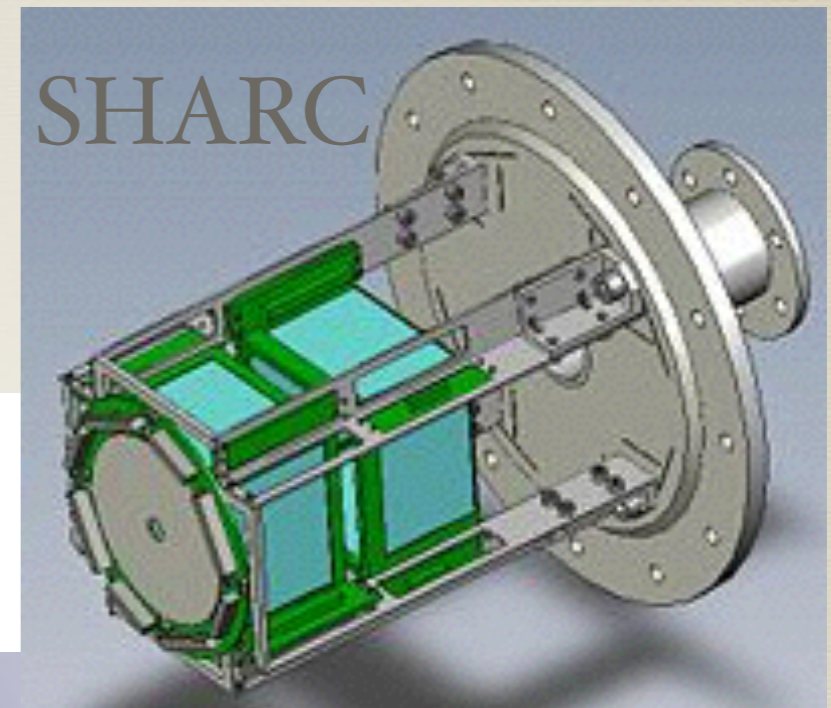
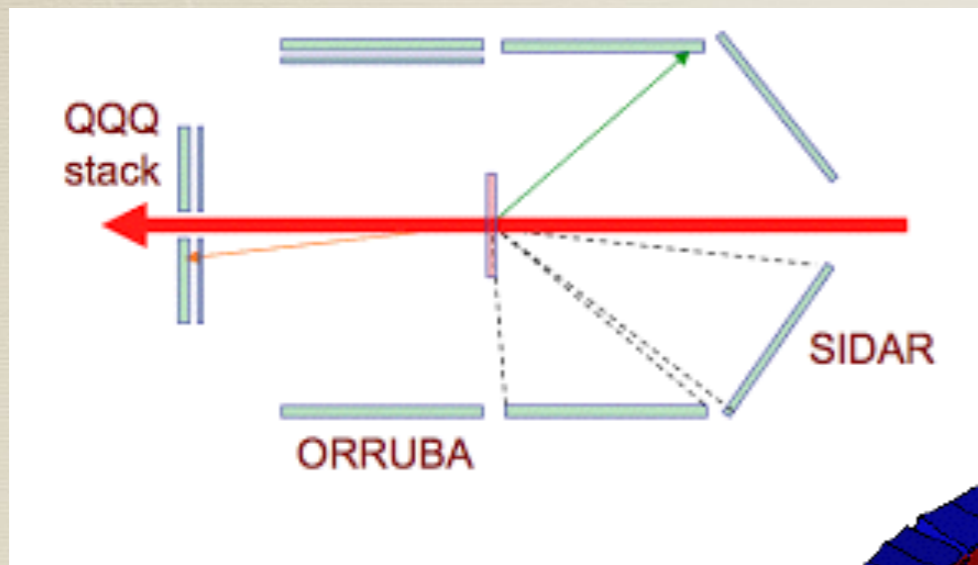
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 - ▶ More difficult for proton to emerge from target
 - ▶ Proton energies much more dependent on scattering angle



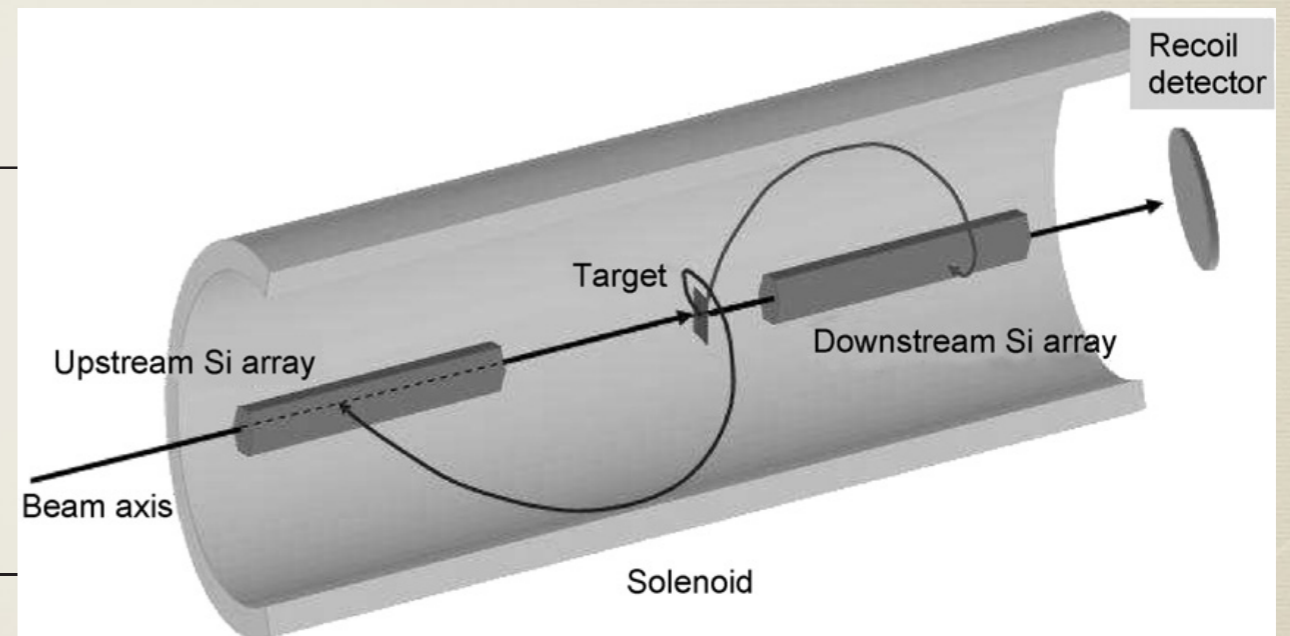
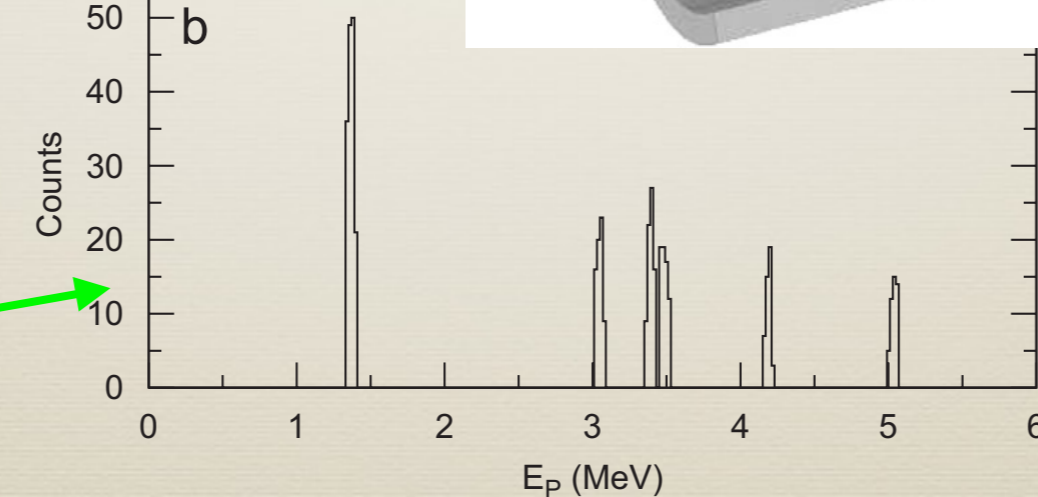
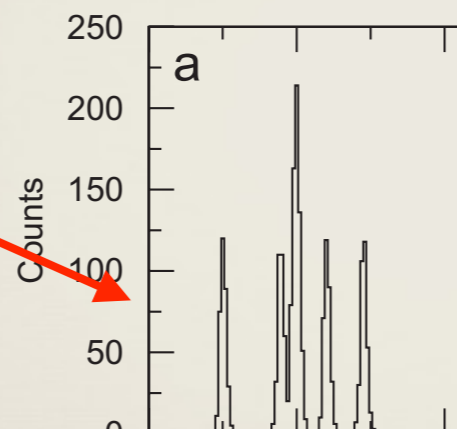
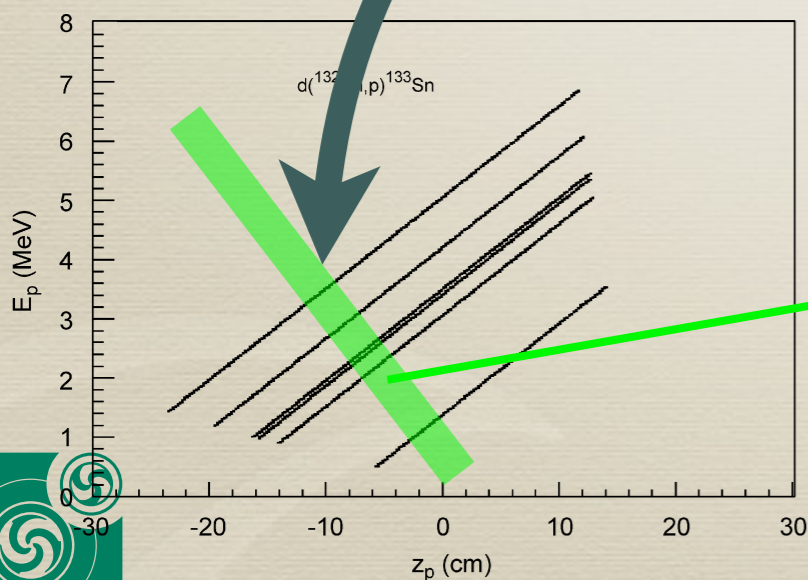
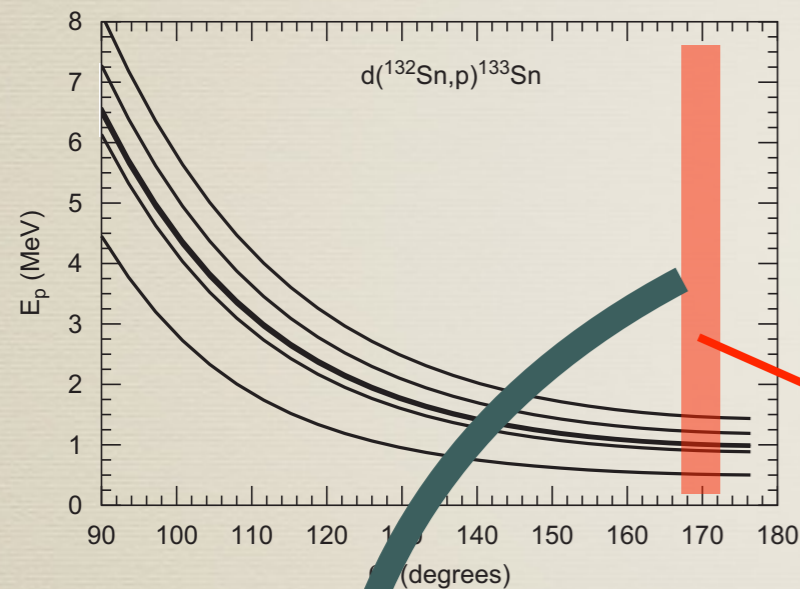
Solutions...

- ▶ Cover as much solid angle as possible

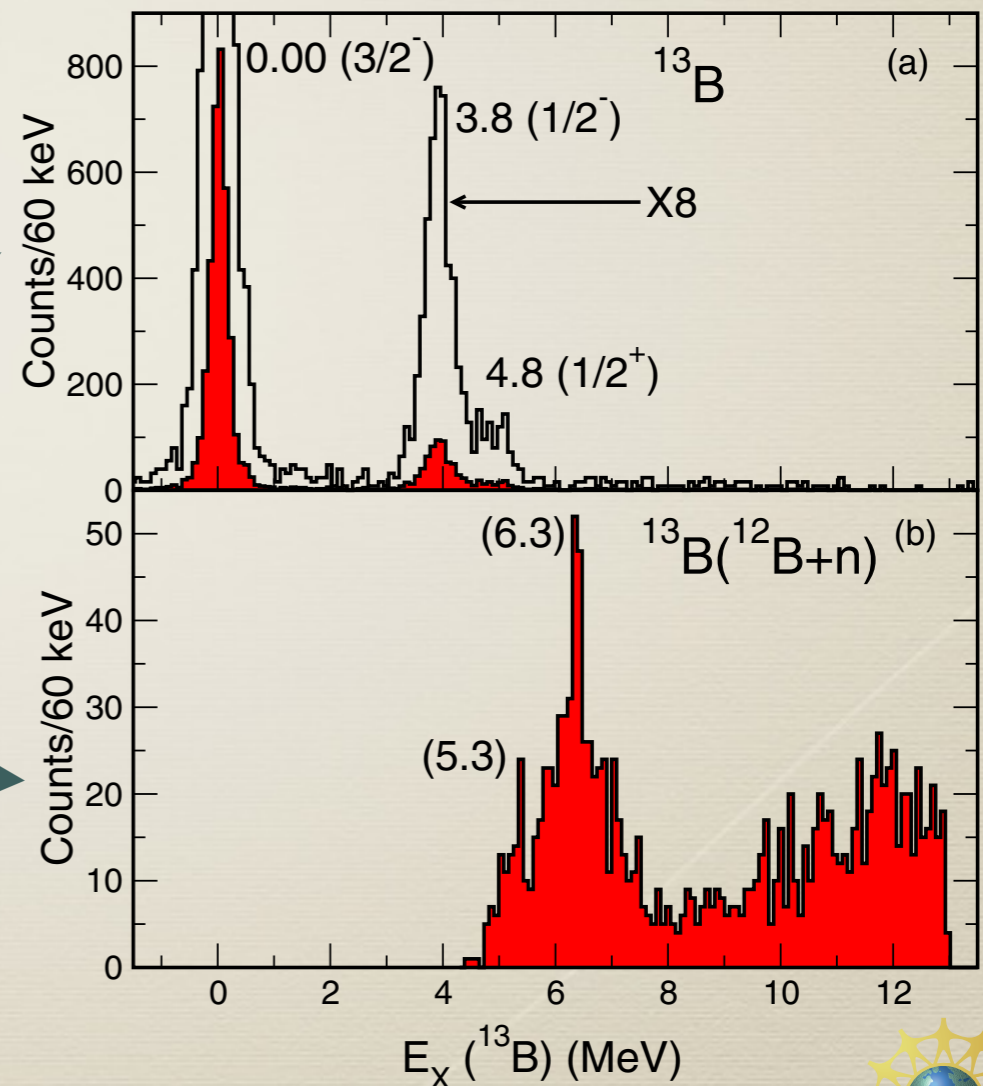
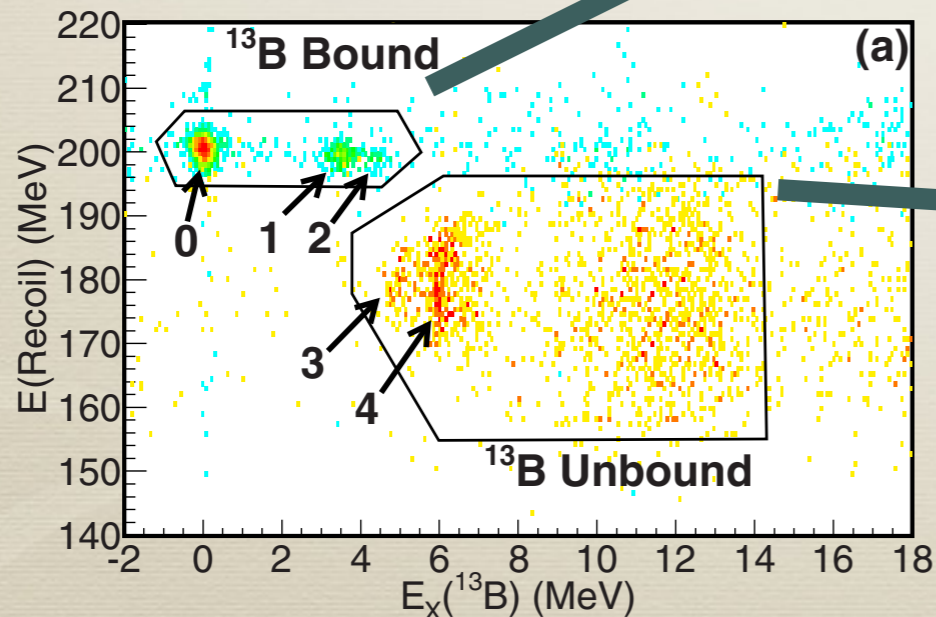
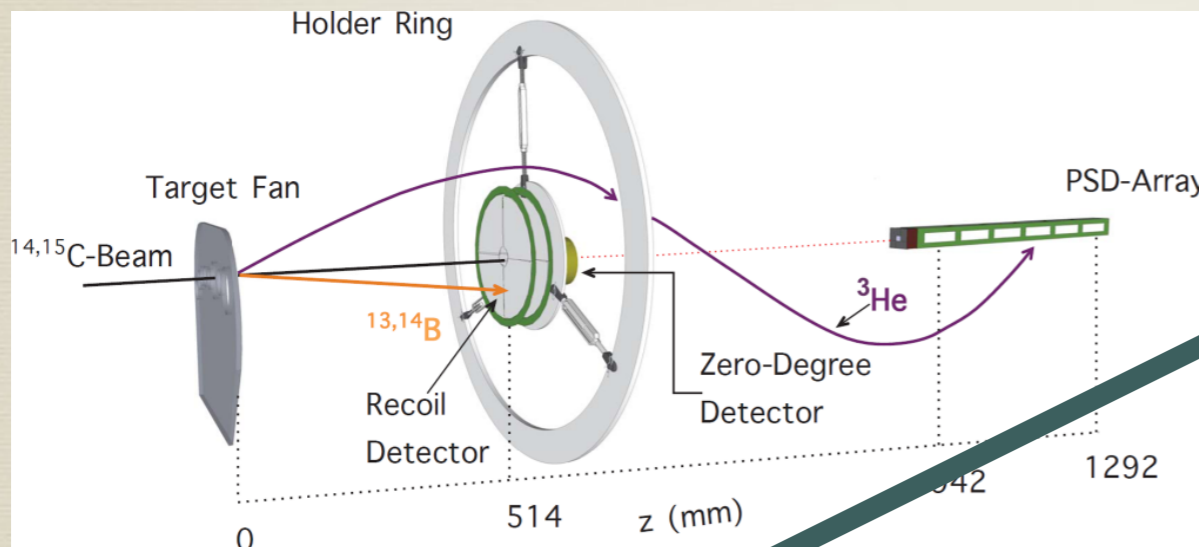
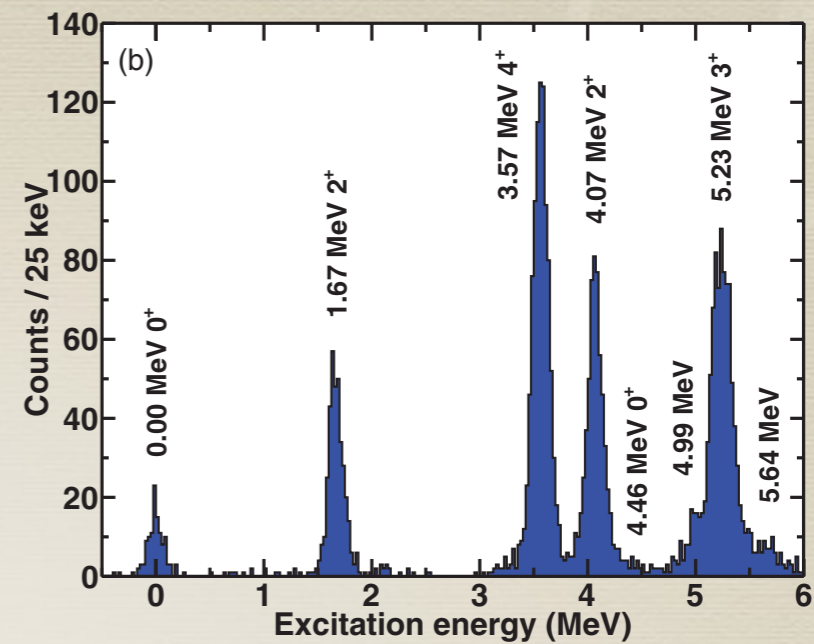
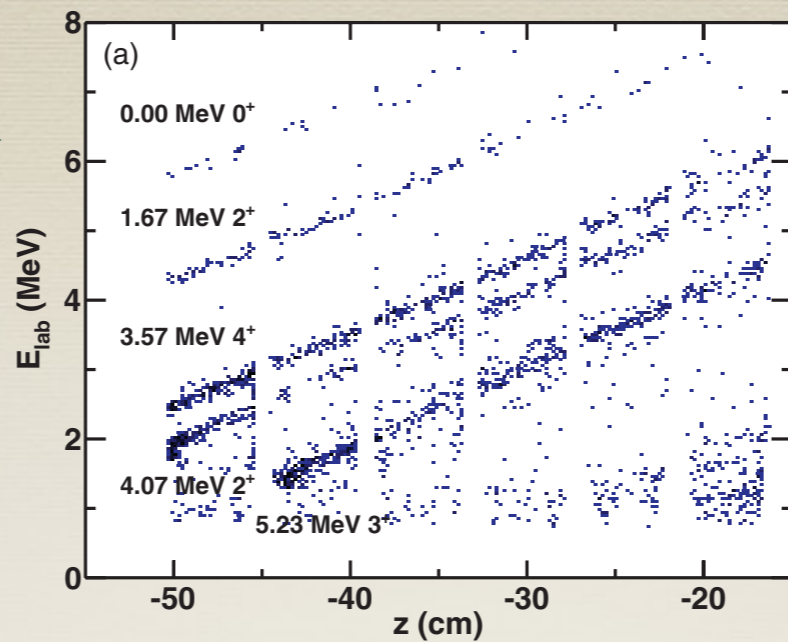
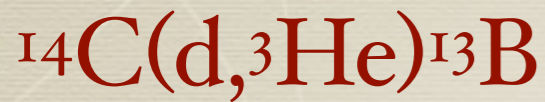
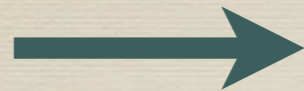


Helios: solving inverse kinematics

- ▶ Kinematic compression of energies
 - ▶ Solenoid spectrometer directly measures center-of-mass energies
 - ▶ Large angular acceptance within solenoid boundaries



A. Wuosmaa *et al.*,
NIMA **507**, 1290 (2007)

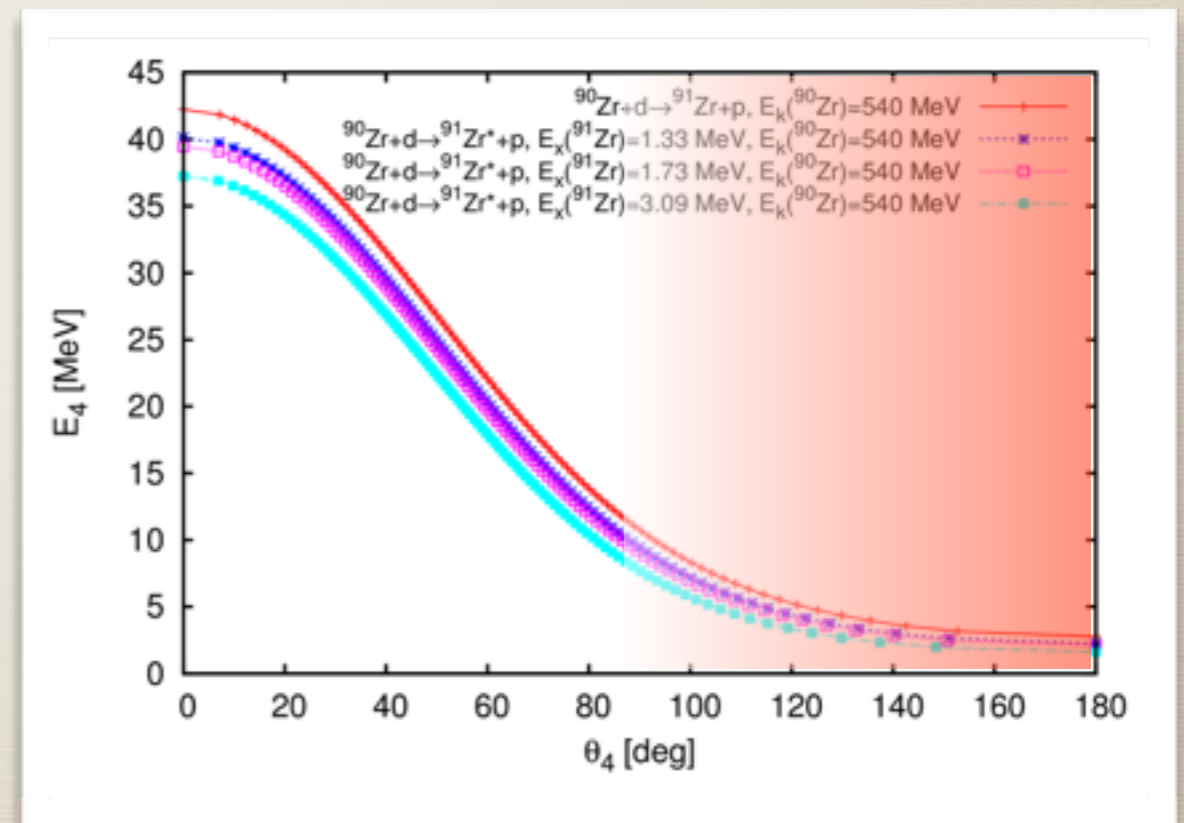


Issue n°2: luminosity!

Event rate in detector $\frac{dR}{dt} = L\sigma_R$ Cross section

$$L = \phi \rho l \varepsilon$$

beam flux
target density
target thickness
detector efficiency



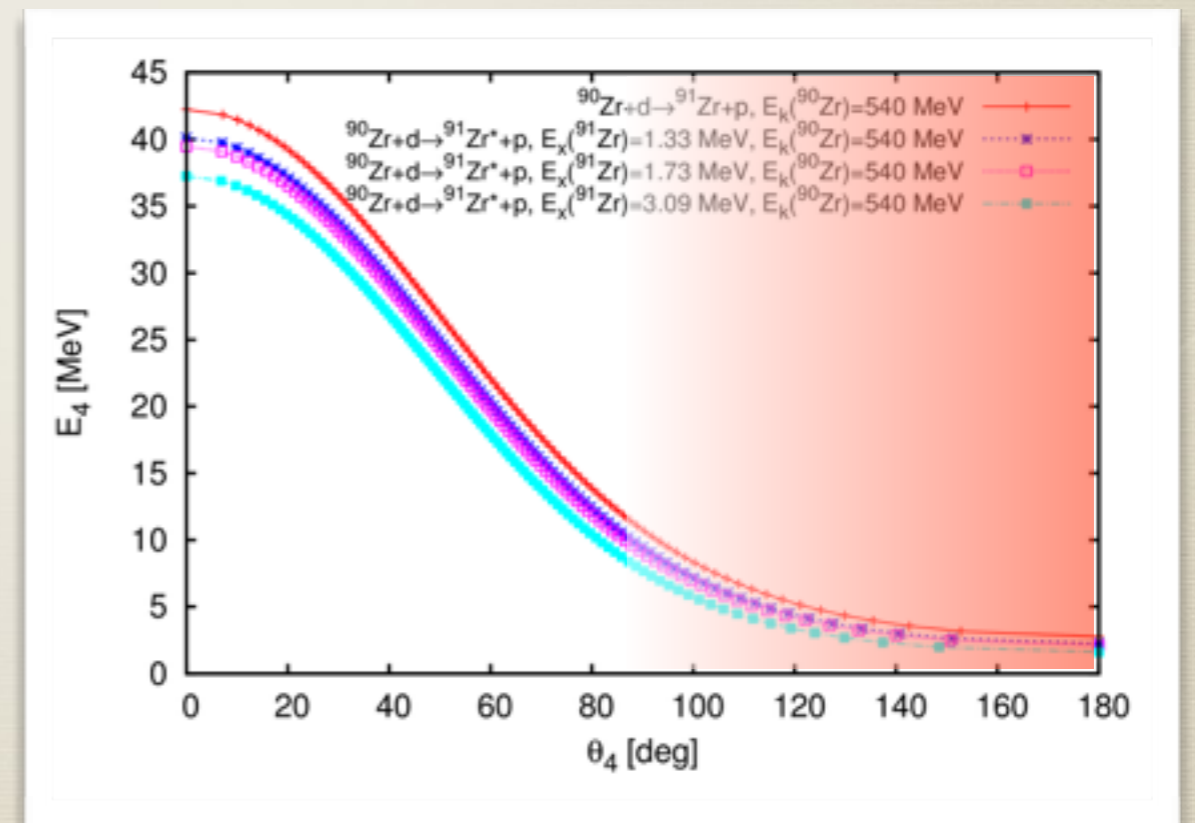
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- ▶ Target cannot be too thick
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 - ▶ Energy lost in target is not recorded: loss of energy resolution
 - ▶ Difficult compromise between resolution and luminosity



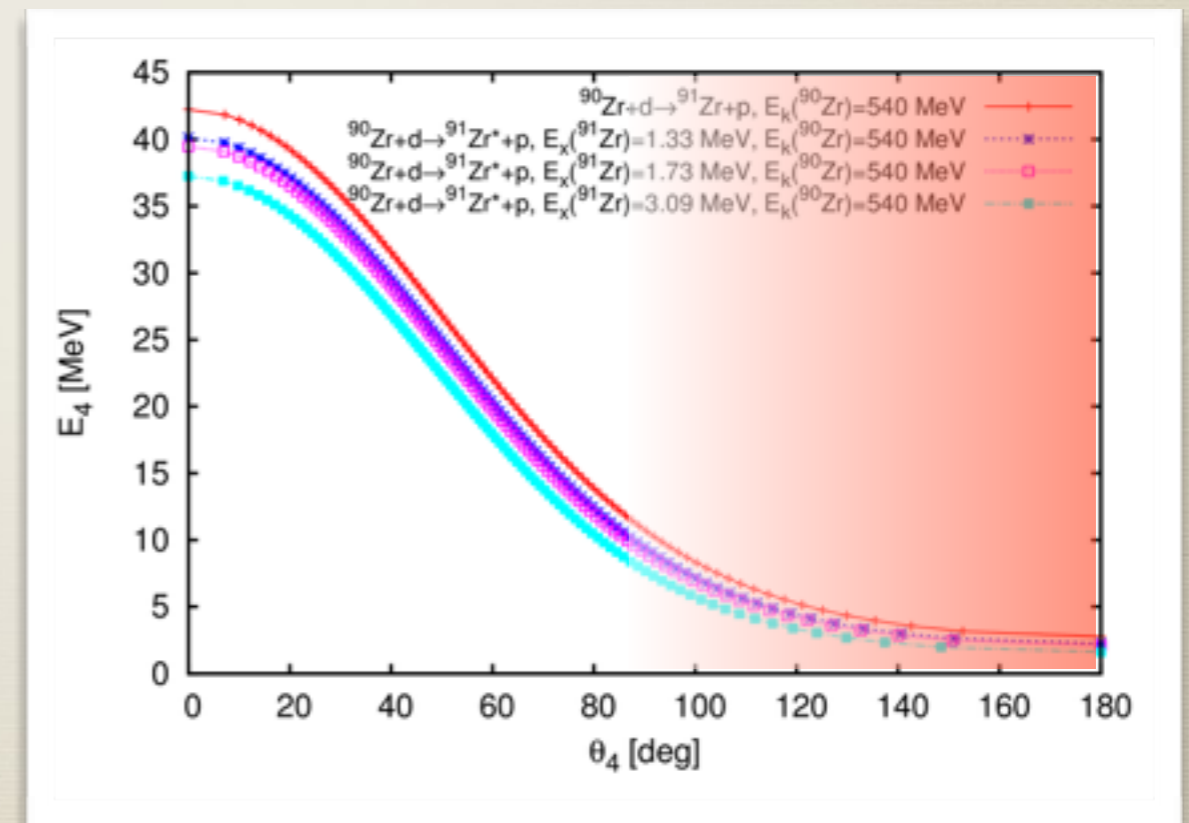
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 - ▶ Difficult compromise between resolution and luminosity
- ▶ Radioactive beams are weak
 - ▶ Intensities several orders of magnitude smaller than stable beams
 - ▶ New experimental techniques needed!



The rise of Active Targets

- ▶ Solving the “too thick target” problem
 - ▶ Active Target: target no longer inert material, but used also to detect particles and/or determine vertex position

Table 1: Active Targets in operation or being constructed.

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S. Beceiro-Novo *et al.*, PPNP **84**, 124 (2015)

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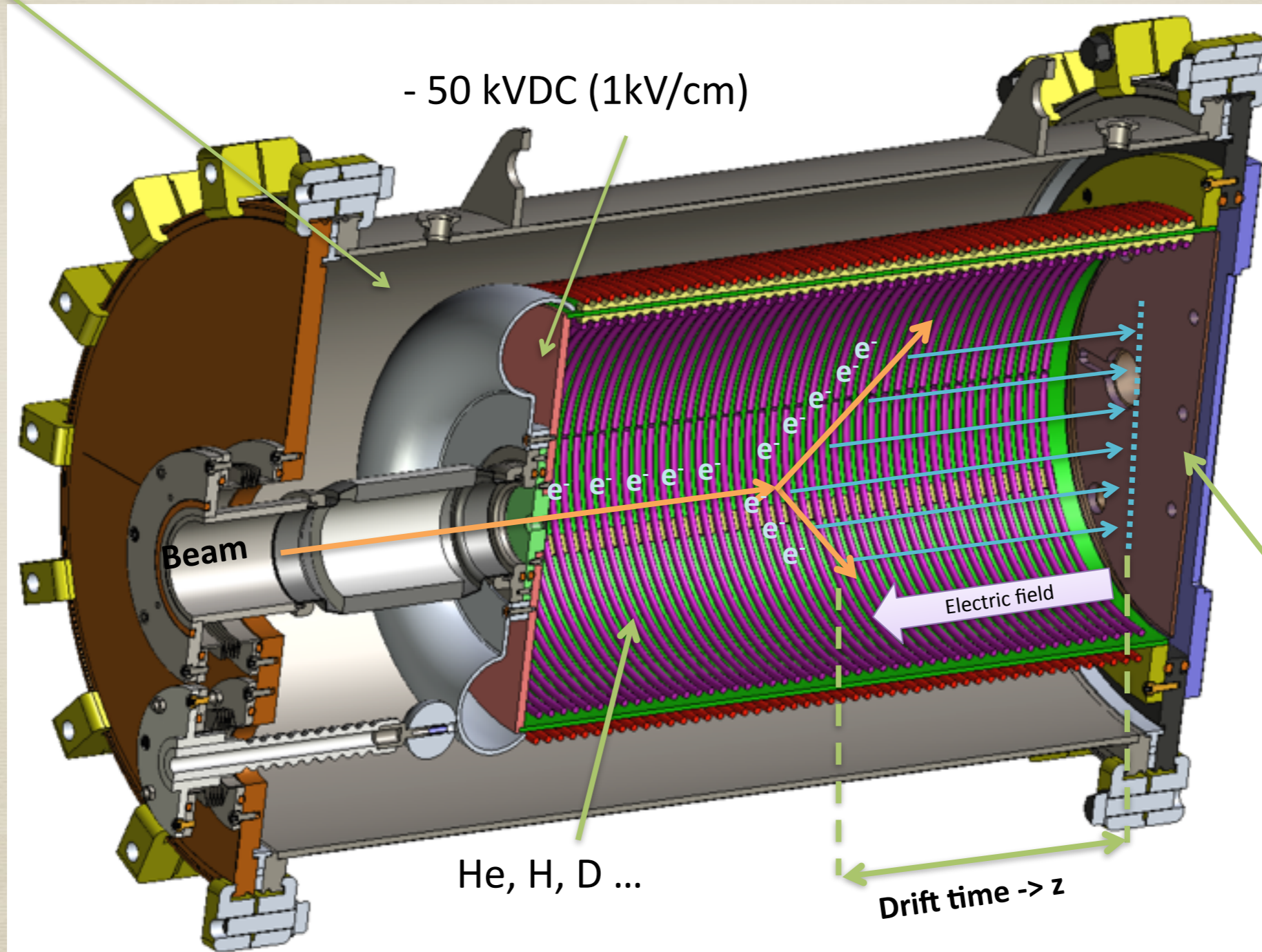
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Example: AT-TPC @ NSCL

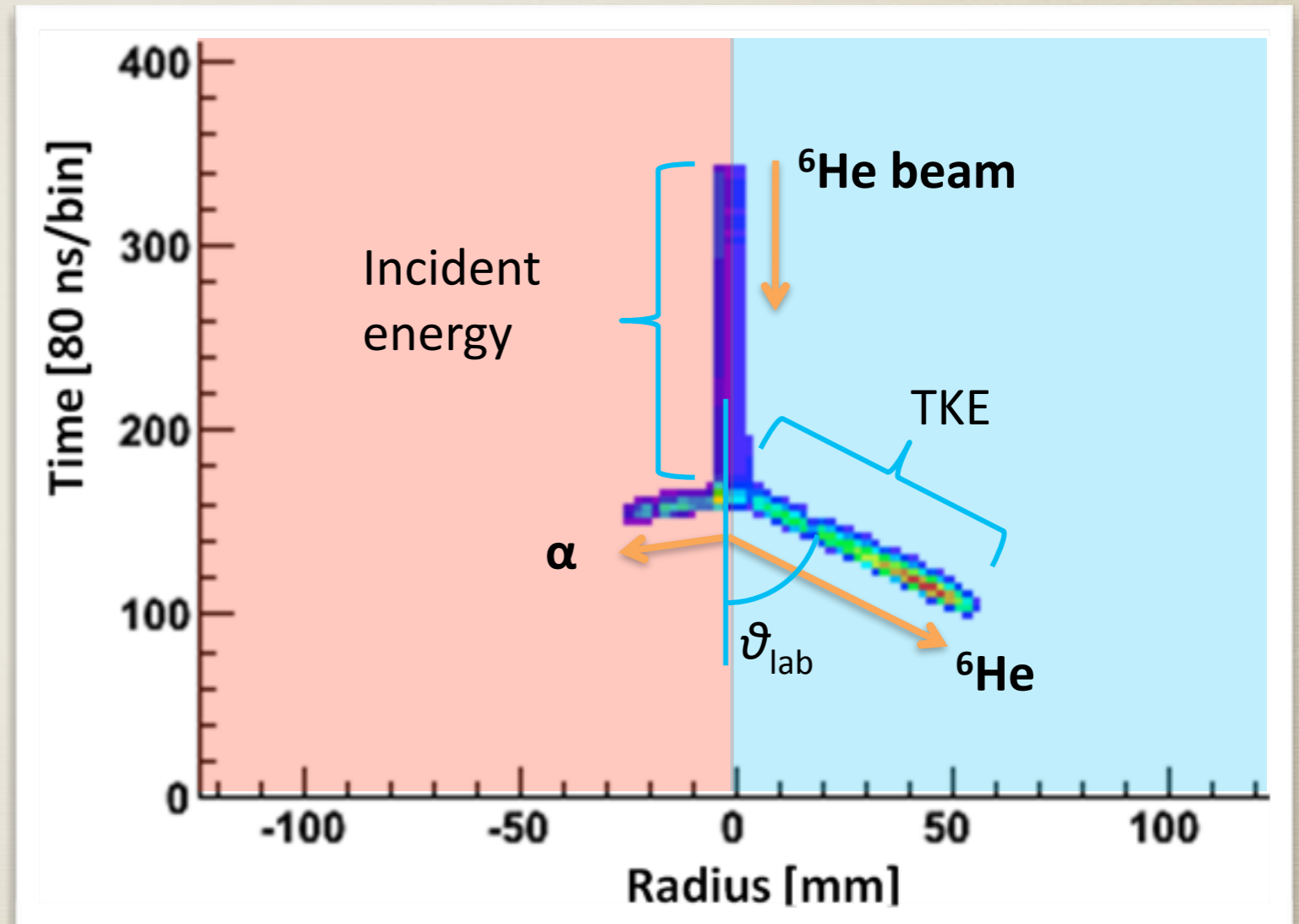
Insulator gas volume

□ N₂ gas 30 kV/cm x 6 cm = 180 kV

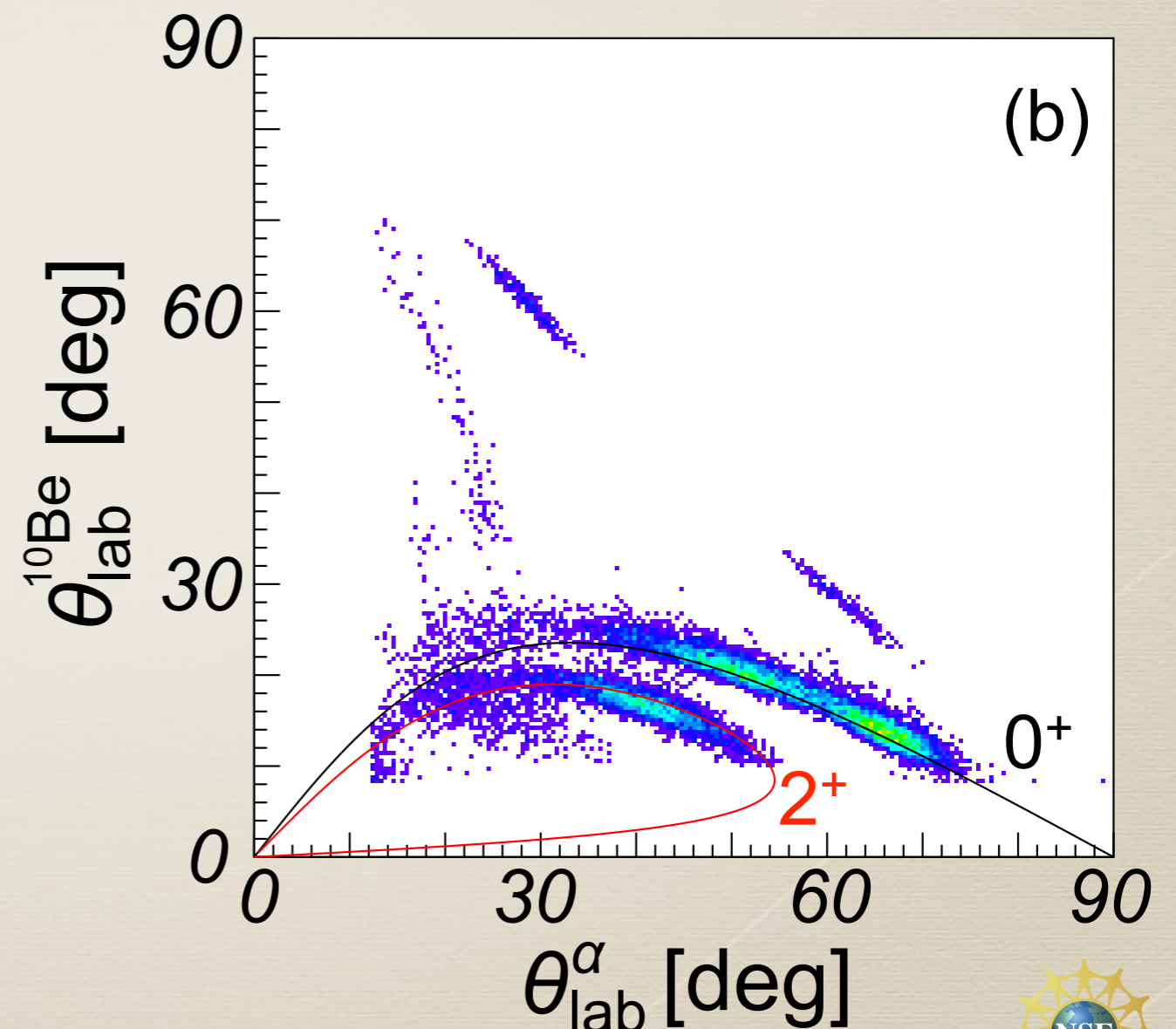
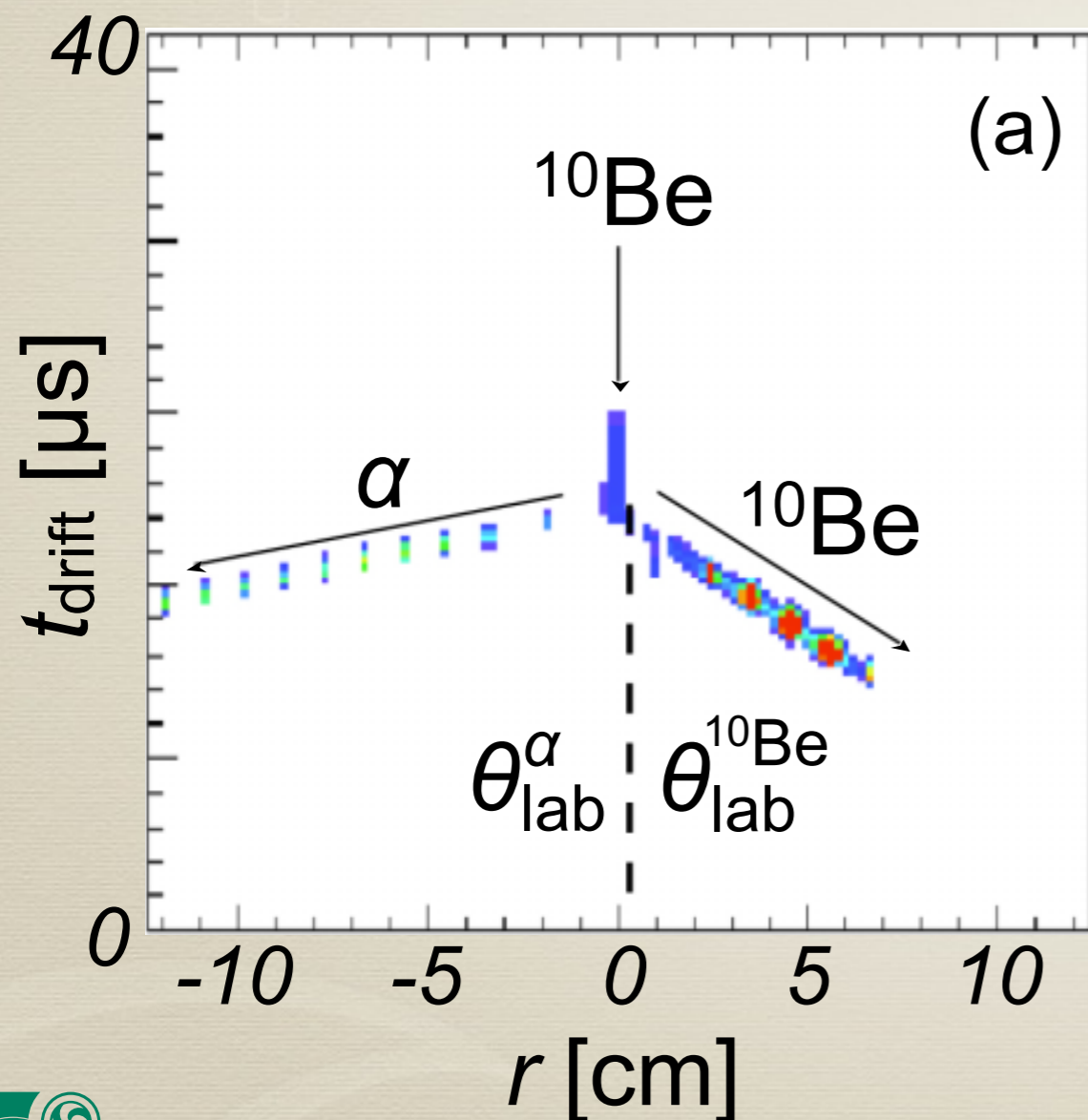


Tracking reactions inside target

- ▶ ${}^6\text{He} + {}^4\text{He}$ elastic scattering
- ▶ Image of charged particle trajectories
- ▶ Beam slows down in gas
- ▶ Vertex location tells reaction energy
- ▶ Energy loss tells which particle is which
- ▶ Length and angle between scattered tracks follow kinematics



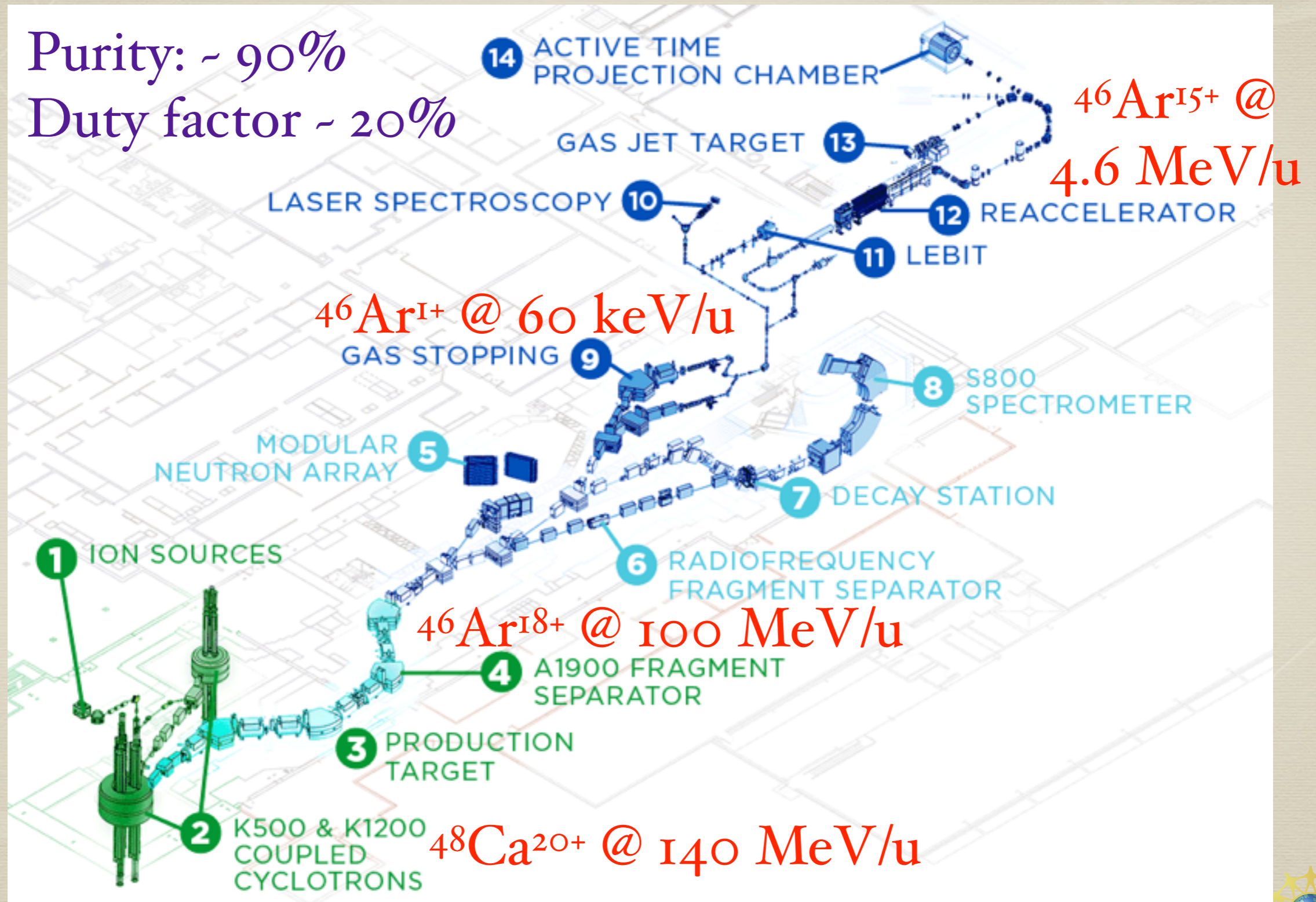
- ▶ $^{10}\text{Be} + ^4\text{He}$ elastic and inelastic scattering
- ▶ Dual gain trick to detect both scattered ^4He and ^{10}Be
- ▶ Clear separation between gs (0^+) and 1st excited state (2^+) at 3.37 MeV



Re-accelerated radioactive beam experiment

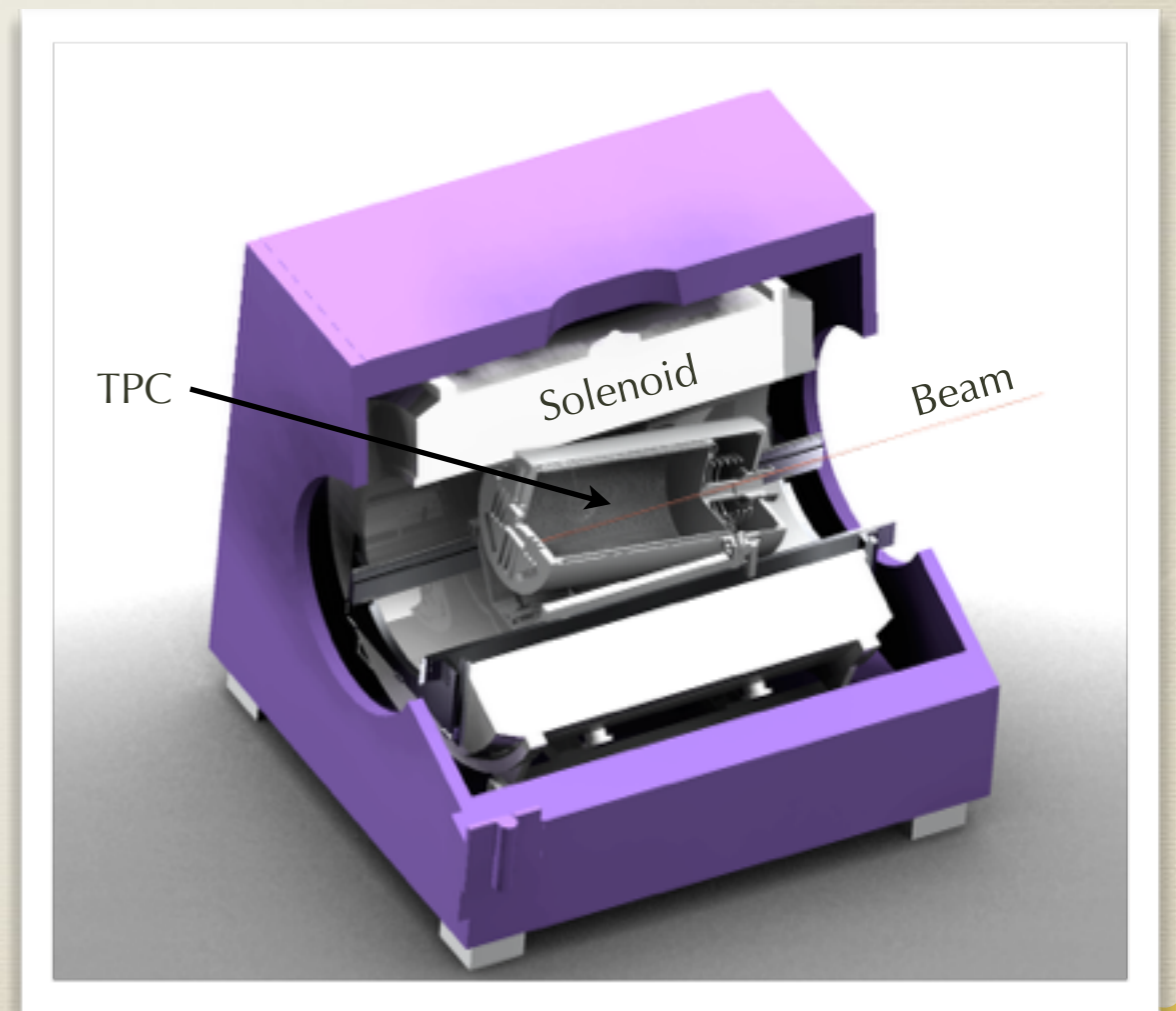
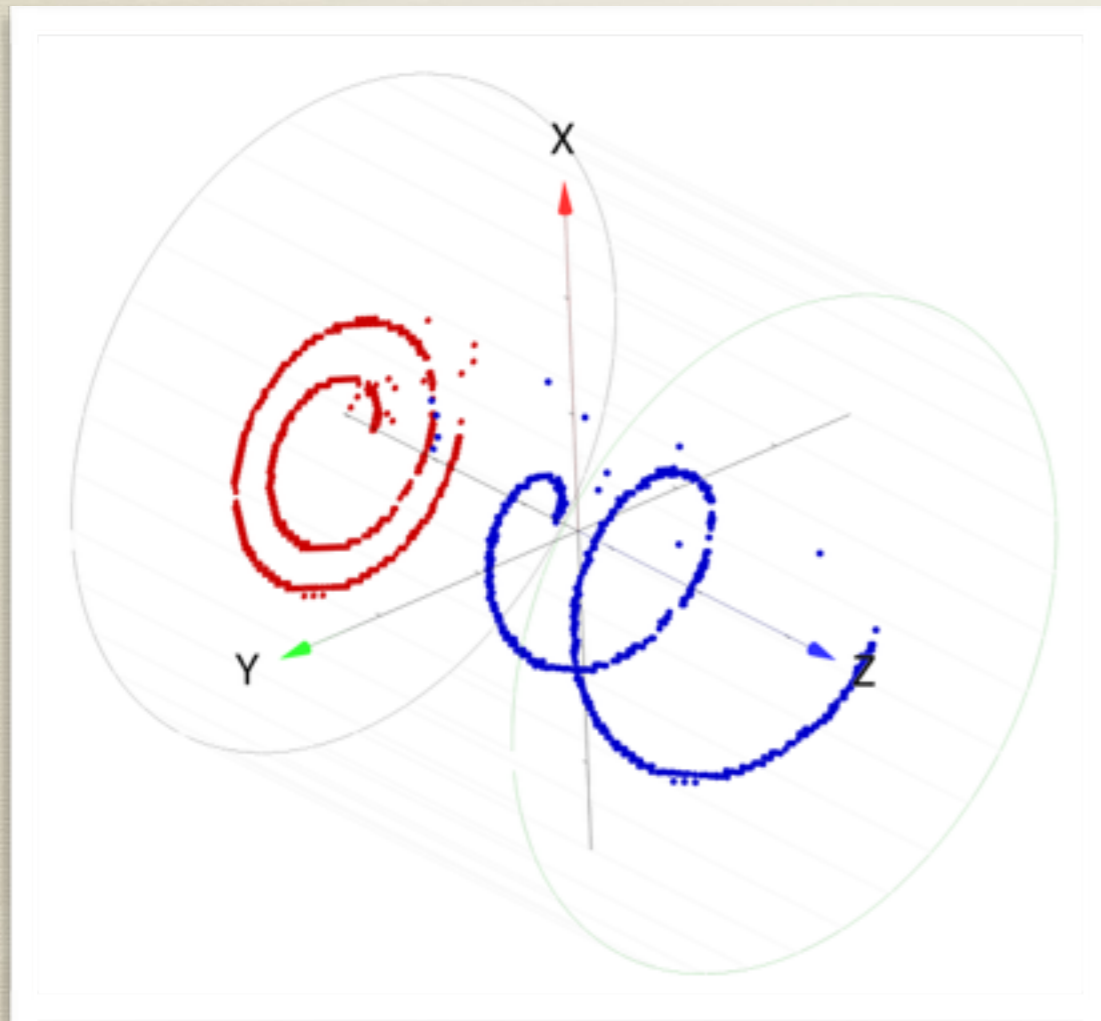
Purity: ~ 90%

Duty factor ~ 20%



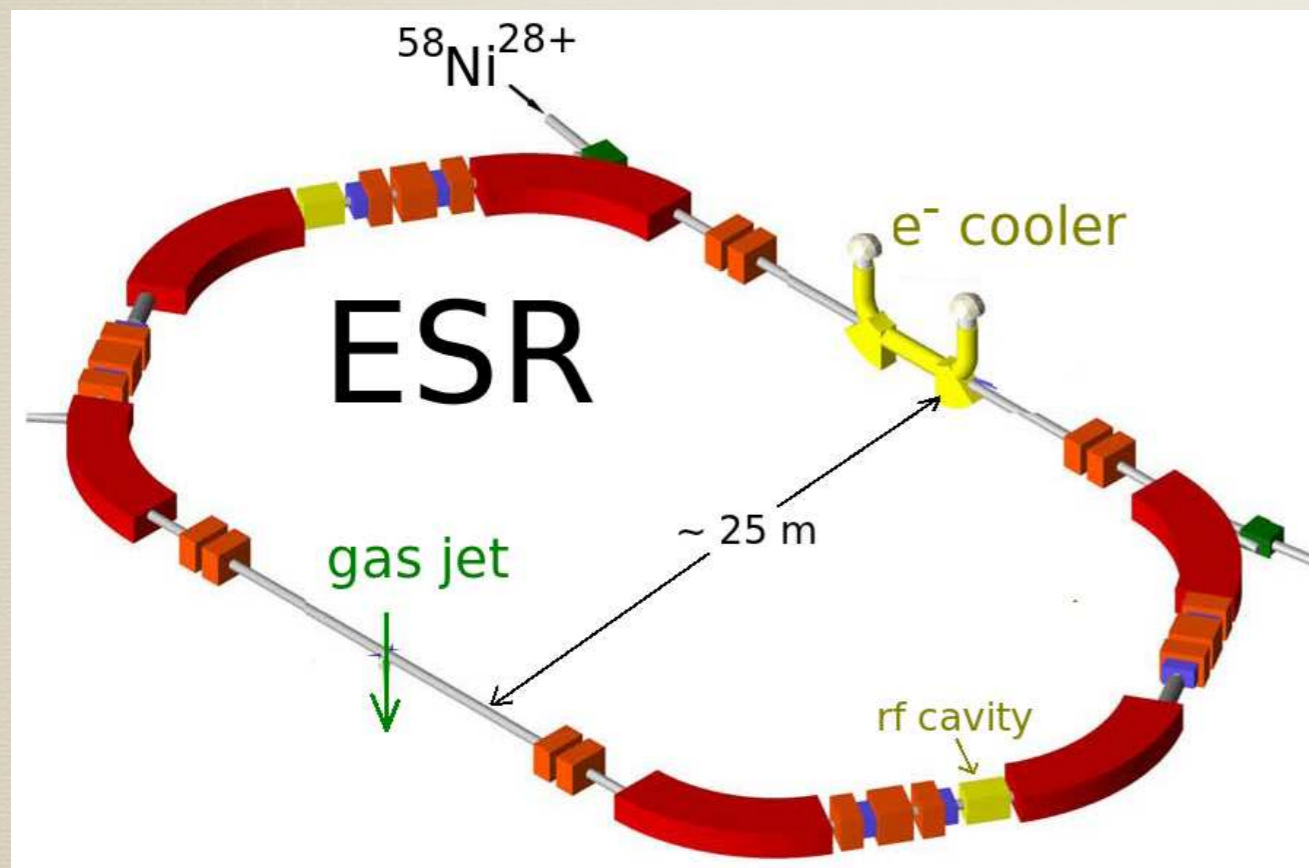
3D camera for nuclear reactions

- ▶ Event from $^{46}\text{Ar}(p,p)$ elastic scattering at 4.6 MeV/u @ NSCL
- ▶ AT-TPC placed inside solenoidal magnetic field to curve trajectories
- ▶ 10,240 pads x 512 time samples = 5.2 Mpixels per event
- ▶ Expected energy resolution 50-100 keV



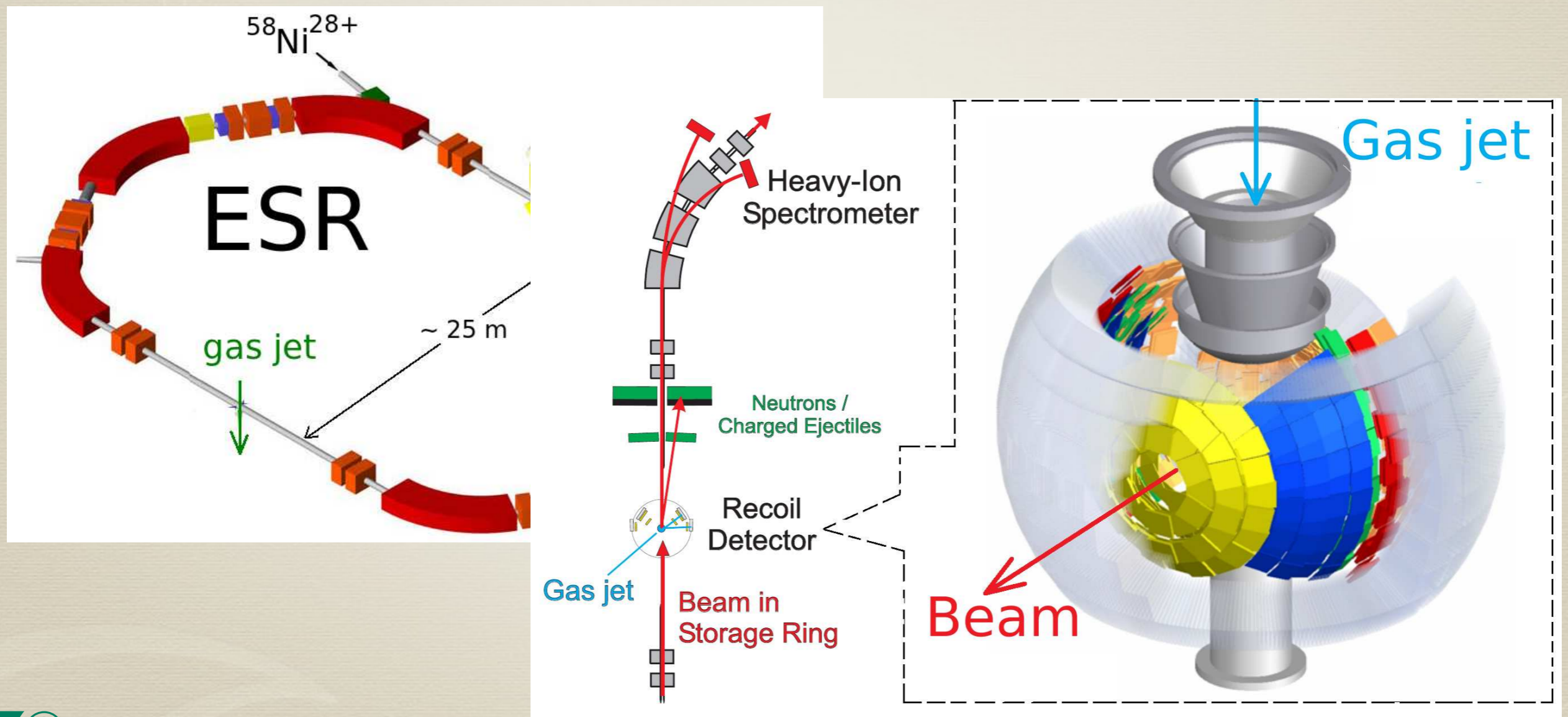
Stored beams at EXL

- ▶ Recirculate radioactive beam onto gas target
 - ▶ Luminosity boost by 10^5 - 10^6 on thin (gas) target



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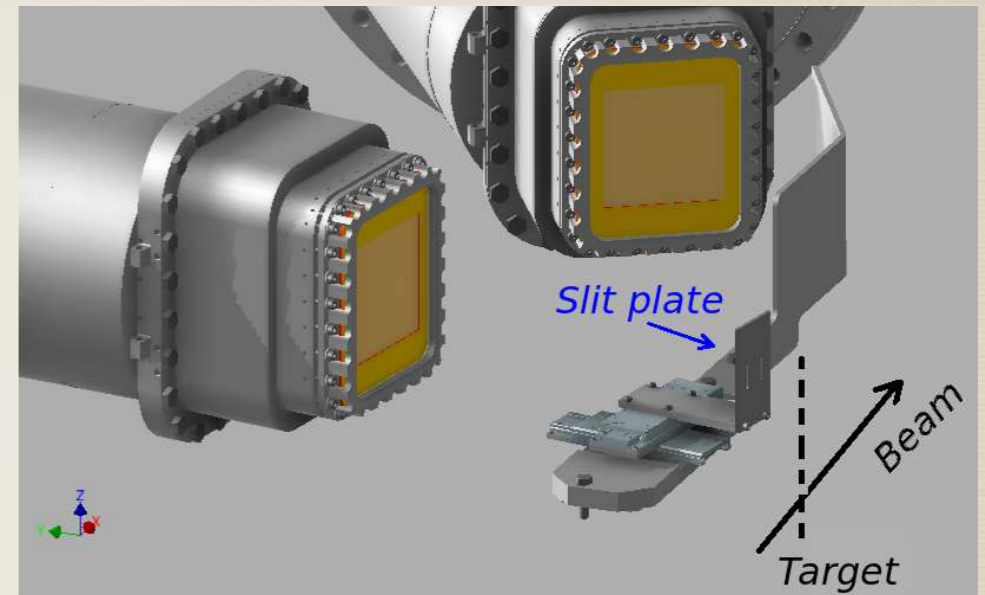


J. C. Zamora, Ph. D. thesis, 2016

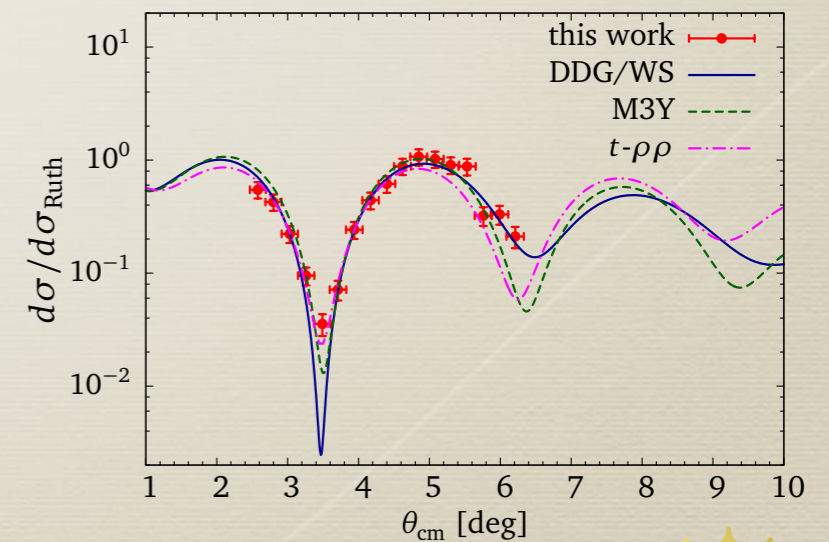
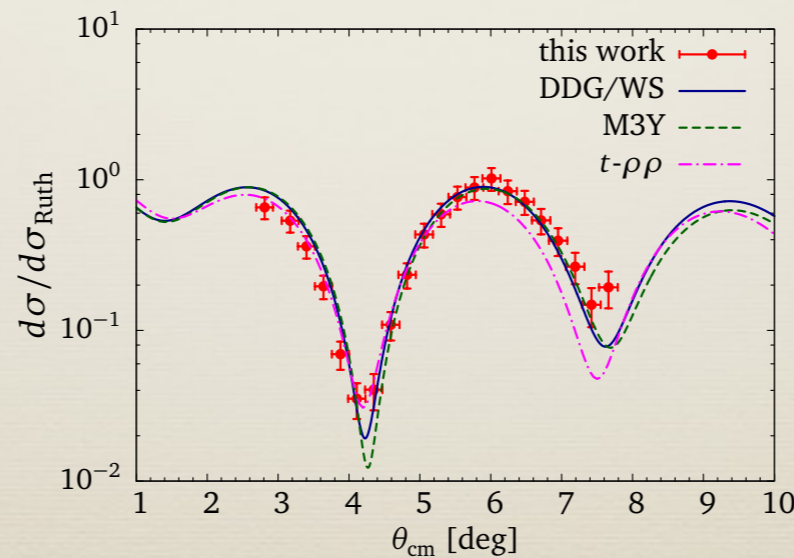
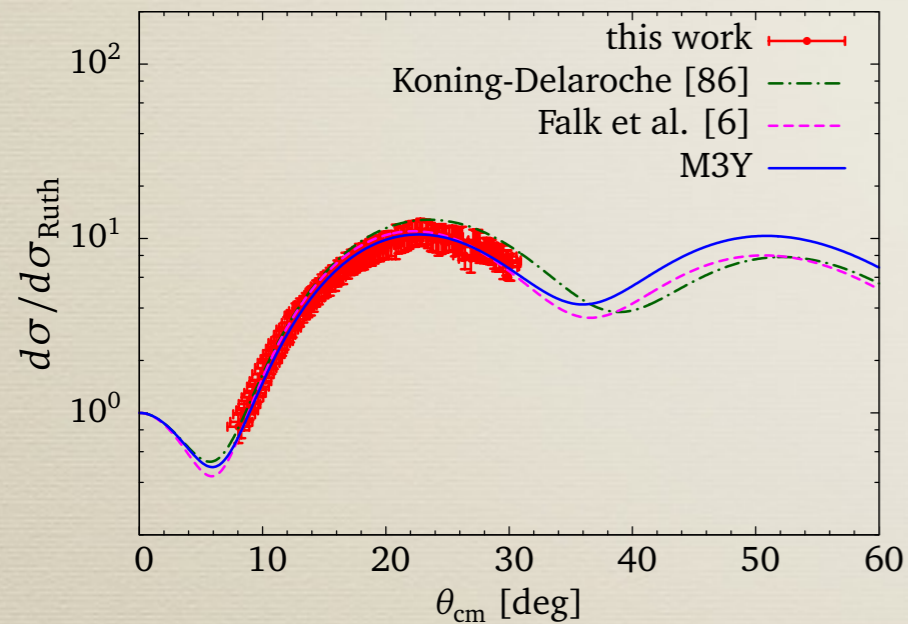
D. Bazin, EBSS 2016, July 22, 2016

First results with stable beams

- ▶ Elastic and inelastic scattering of $^{20}\text{Ne}+p$ and $^{58}\text{Ni}+\alpha$
 - ▶ Luminosities from $6 \cdot 10^{25}$ to $6 \cdot 10^{27} \text{ cm}^{-2}\text{s}^{-1}$ achieved
 - ▶ High vacuum environment
 - ▶ Emittance growth



$^{20}\text{Ne}+p$
50 MeV/u



$^{58}\text{Ni}+\alpha$
100 & 150 MeV/u

(a) 100 MeV/u.

(b) 150 MeV/u.

Luminosity from fast beams

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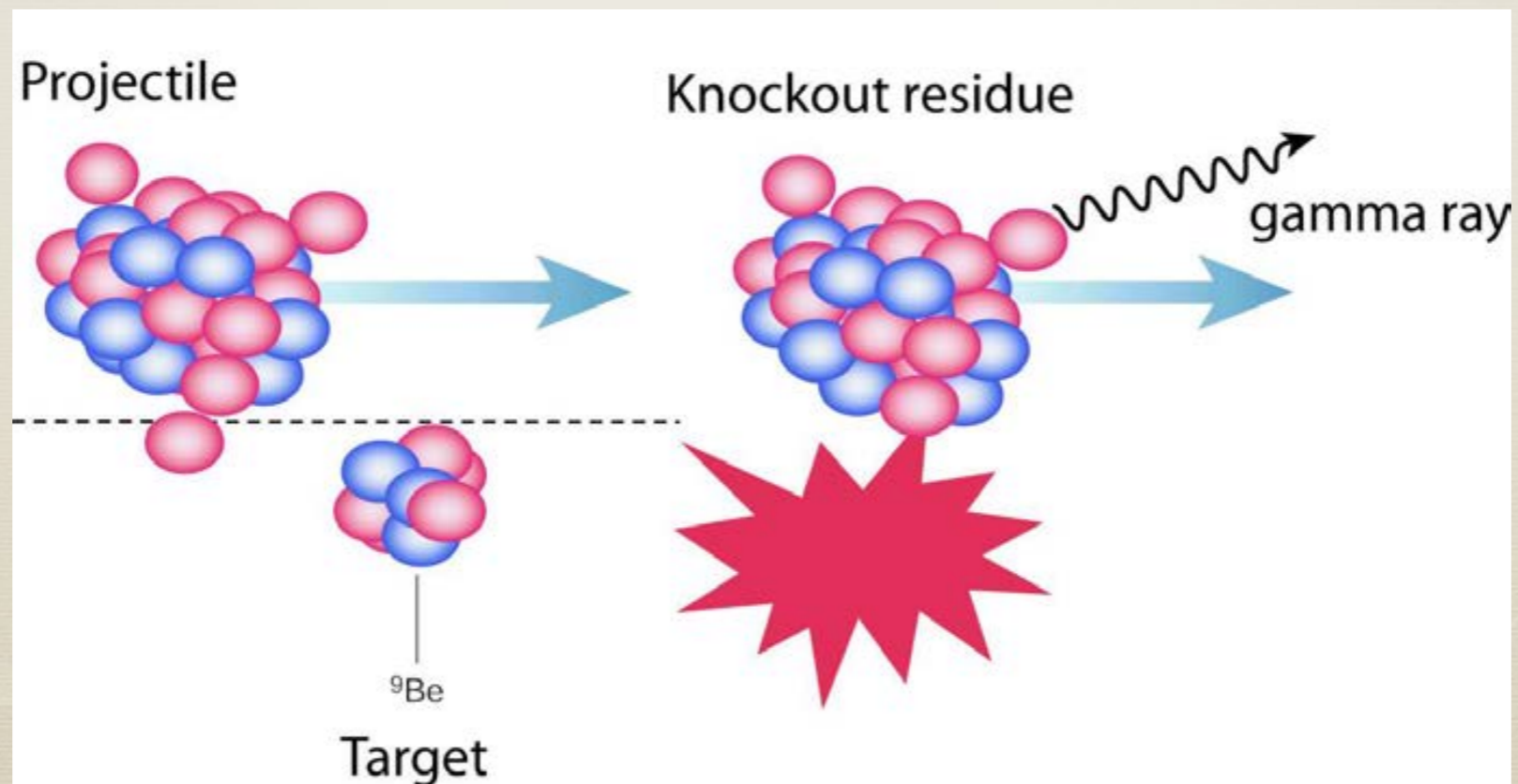
- ▶ In-beam γ -ray spectroscopy and invariant-mass spectroscopy
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Luminosity from fast beams

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 - ▶ For unbound states: long time-of-flight neutron arrays placed around 0° to reconstruct invariant mass
- ▶ Large number of devices are available and developing around the world

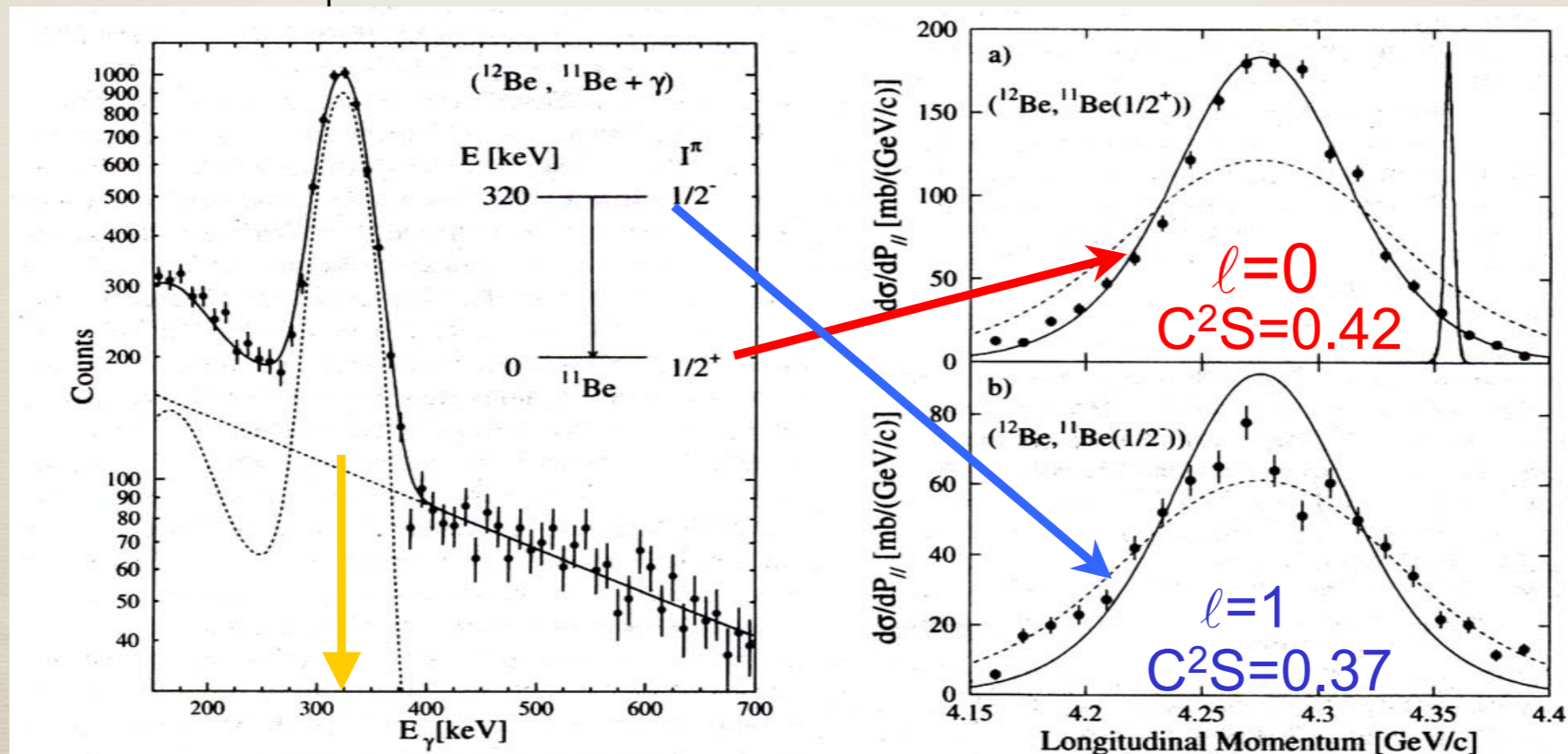
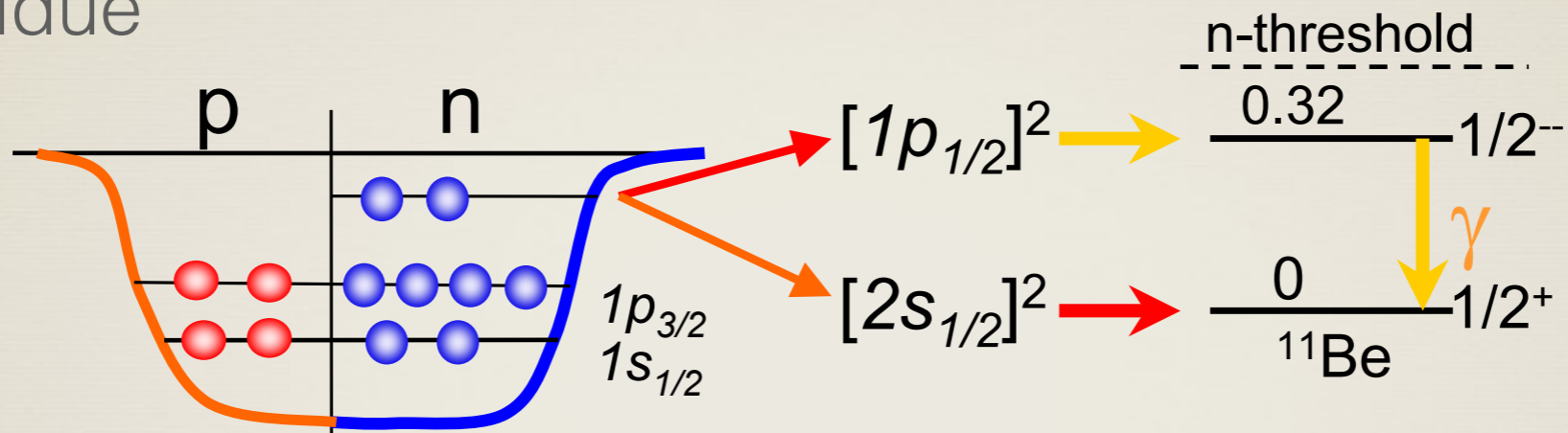
in-beam γ -ray spectroscopy

- ▶ Reactions with fast radioactive beam
 - ▶ Reaction residue loses one or a few nucleons
 - ▶ Only residue is collected and detected at forward angles
 - ▶ Thick targets can be used (luminosity!)
 - ▶ High efficiency array detects Doppler-shifted γ -ray from residue

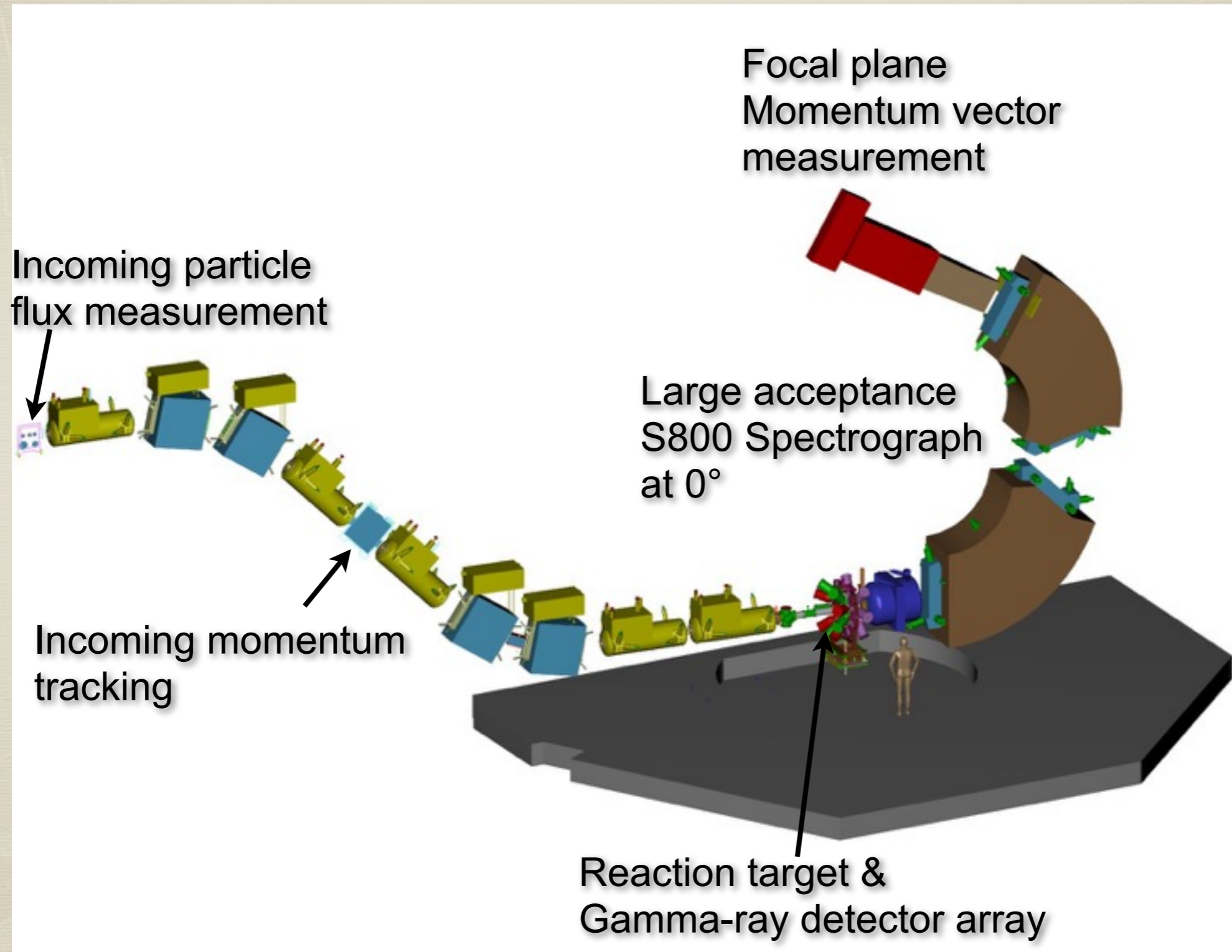


Knockout reactions

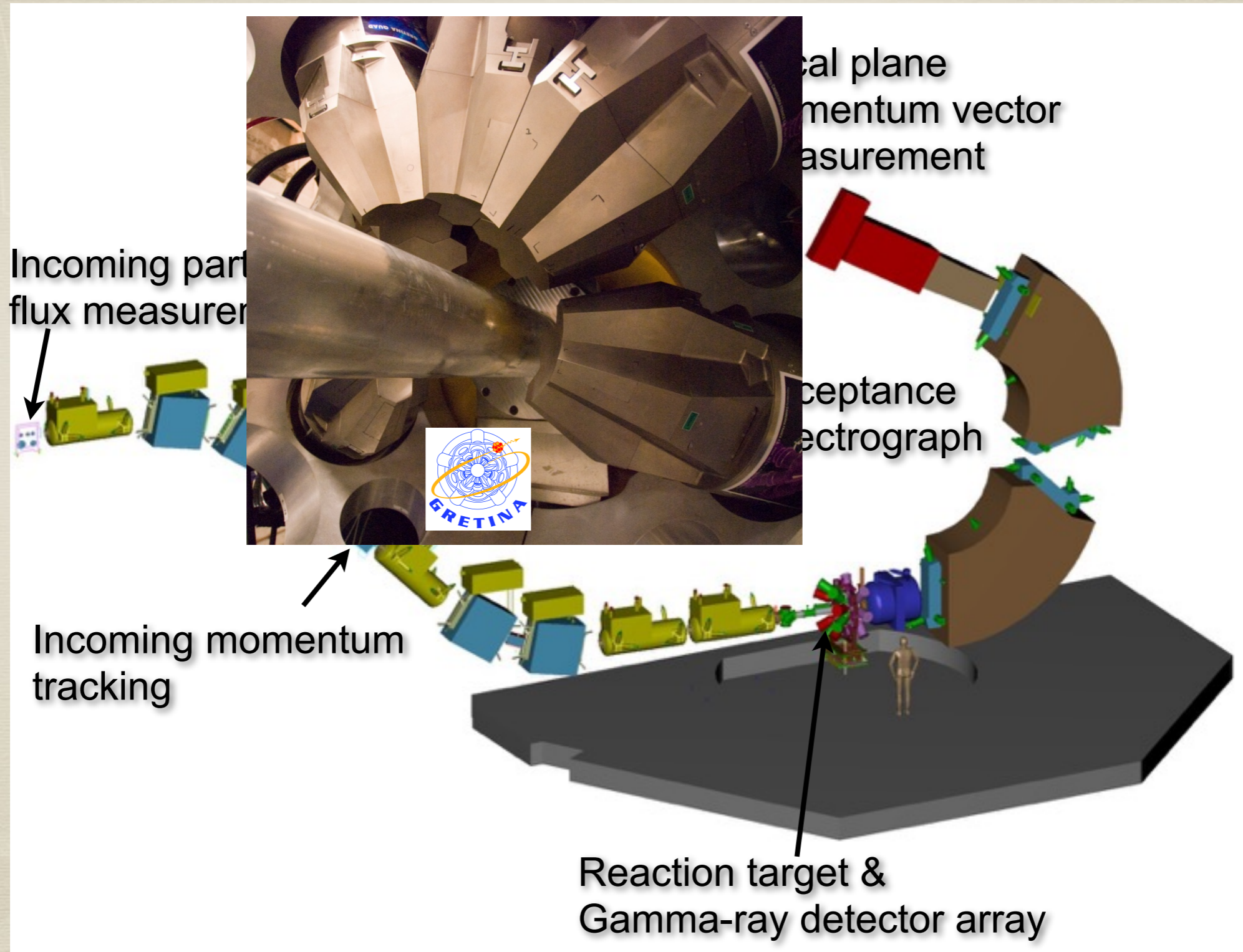
- ▶ Pillar of today's radioactive nuclei spectroscopy
- ▶ Angular momentum of removed nucleon from parallel momentum of residue



S800 spectrograph + Gretina array



S800 spectrograph + Gretina array



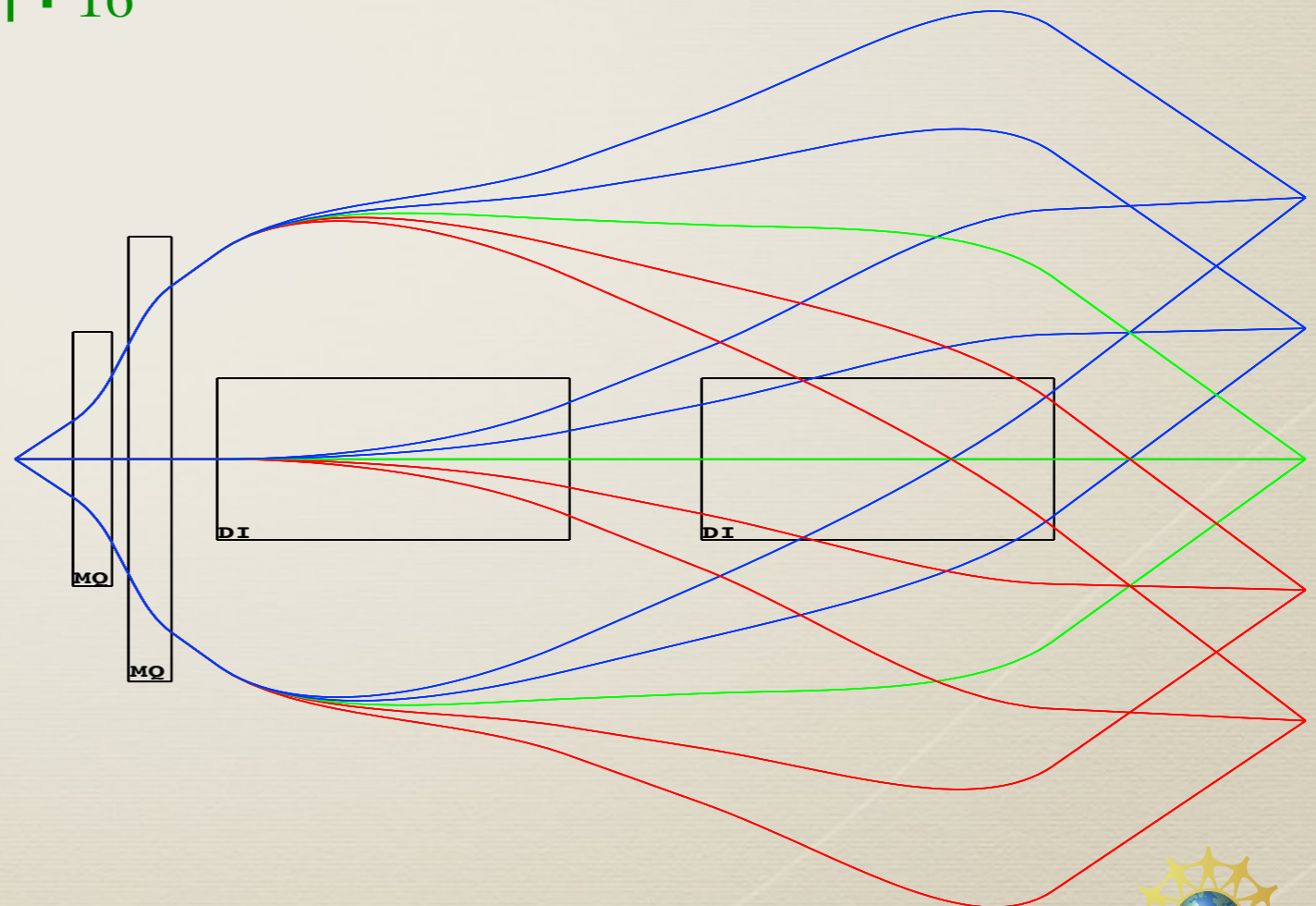
S800: “software” spectrograph

- ▶ Magnetic elements have fringe fields and imperfections
- ▶ Optics are not perfect and deviate from 1st order (linear)
- ▶ Deviations are called “high order aberrations”

$$X_f = X_i T_{11} + a_i T_{12} + d_i T_{16}$$

| | Transfer | Sigma | Inverse | Emittances | | |
|----|----------|-----------|---------|------------|------|--------|
| | x(m) | a(rad) | y(m) | b(rad) | l(m) | d(1) |
| xf | -0.893 | -9.12e-06 | 0 | 0 | 0 | -9.54 |
| af | -0.727 | -1.12 | 0 | 0 | 0 | -0.172 |
| yf | 0 | 0 | -0.59 | -0.704 | 0 | 0 |
| bf | 0 | 0 | 1.33 | -0.102 | 0 | 0 |
| lf | 6.78 | 10.7 | 0 | 0 | 1 | -10.3 |
| df | 0 | 0 | 0 | 0 | 0 | 1 |

Dismiss



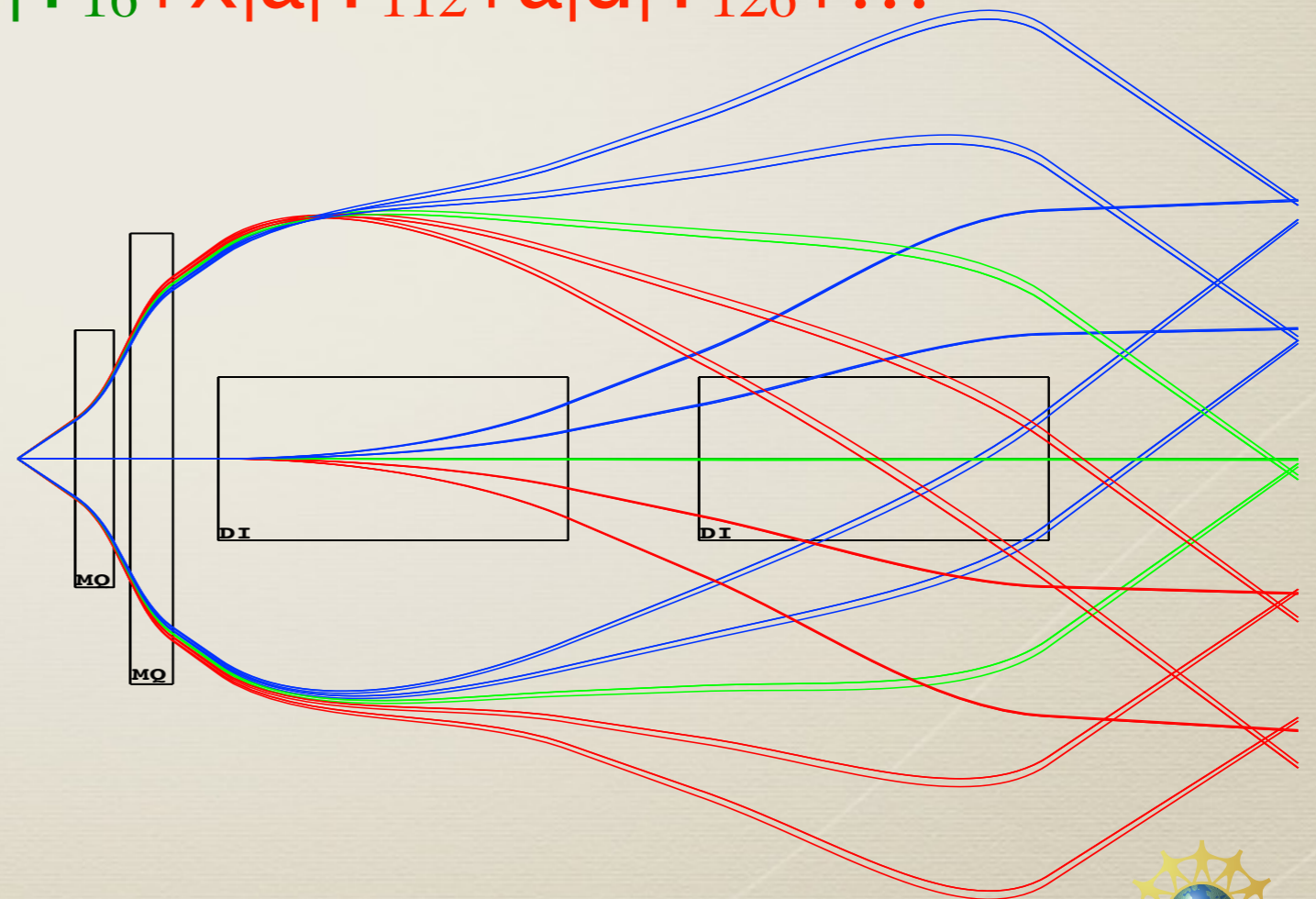
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$$X_f = x_i T_{11} + a_i T_{12} + d_i T_{16} + x_i a_i T_{112} + a_i d_i T_{126} + \dots$$

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Dismiss



Trajectory reconstruction

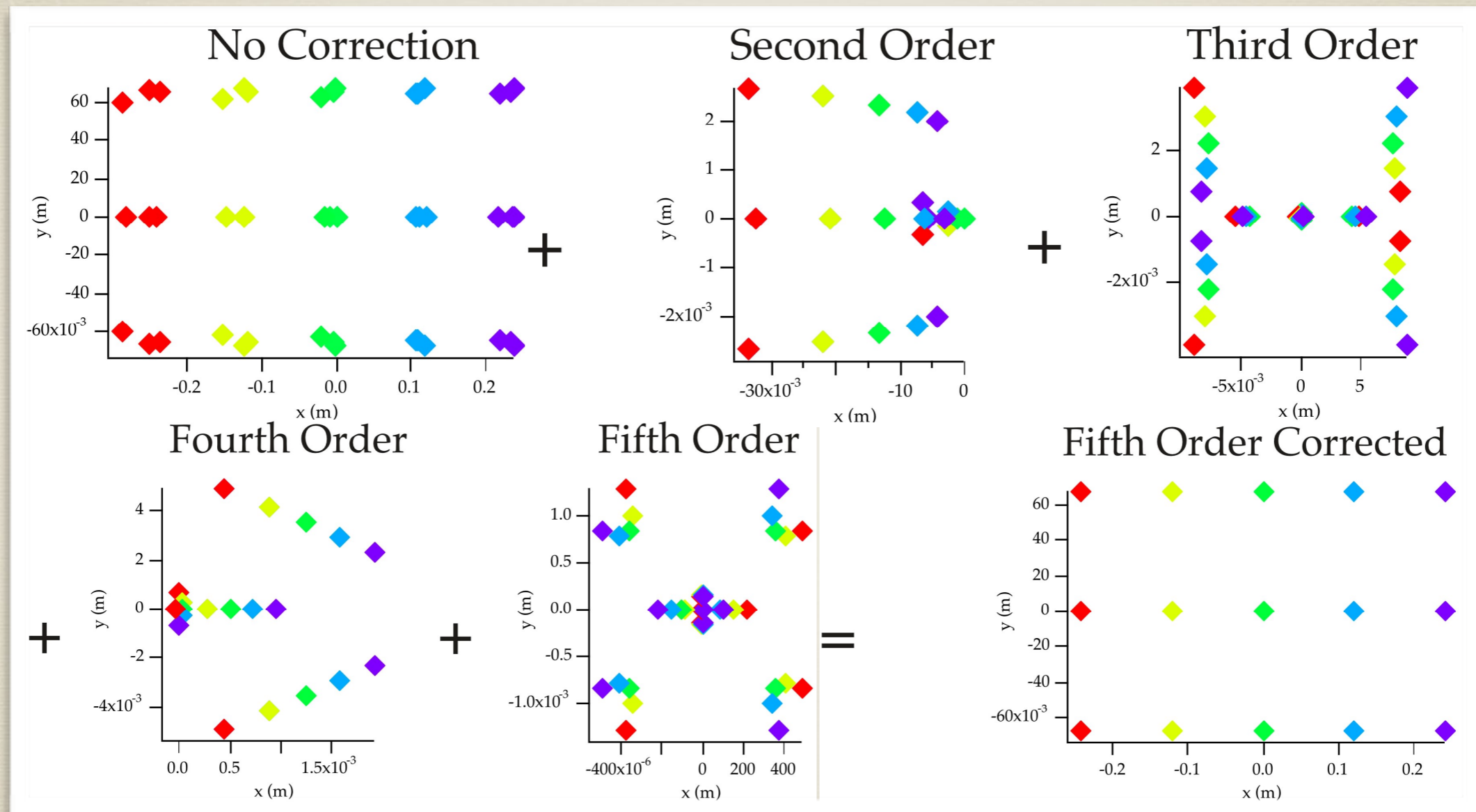
- ▶ Deduce parameters of particle at the target location from measurements at the focal plane
- ▶ Calculate 4x4 forward matrix up to order N (COSY infinity)
- ▶ Inverse forward matrix assuming $x_i=0$ to get d_i
- ▶ Apply inverse matrix (map) to data to extract energy and scattering angle at the target location

$$\begin{pmatrix} x_f \\ \theta_f \\ y_f \\ \phi_f \end{pmatrix} = S \begin{pmatrix} \delta_i \\ \theta_i \\ y_i \\ \phi_i \end{pmatrix} \quad \longrightarrow \quad \begin{pmatrix} \delta_i \\ \theta_i \\ y_i \\ \phi_i \end{pmatrix} = S^{-1} \begin{pmatrix} x_f \\ \theta_f \\ y_f \\ \phi_f \end{pmatrix}$$

$$(x_i = 0)$$

Aberration corrections

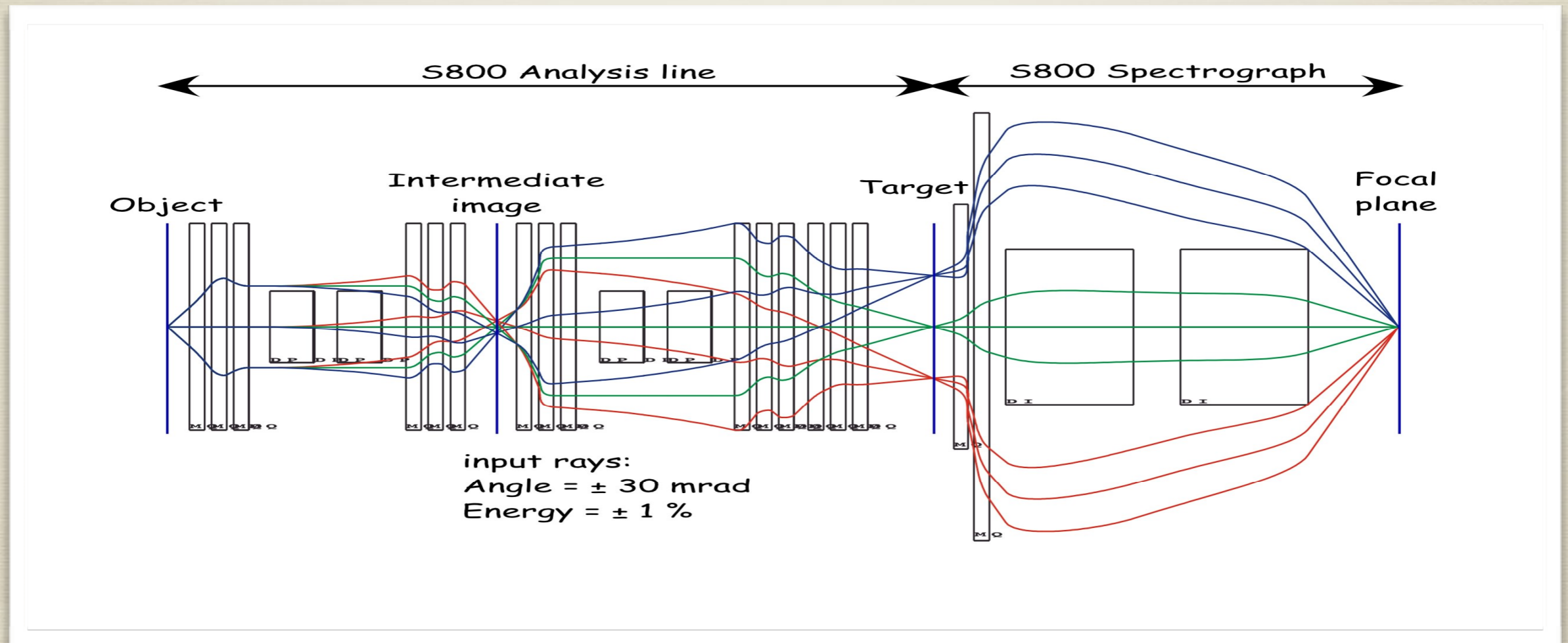
Initial grid: $\Delta d = \pm 5\%$, $\Delta a = \pm 60 \text{ mrad}$, $\Delta b = \pm 90 \text{ mrad}$



Expected energy resolution for a 1mm beam spot size: 1 part in 5,000 (0.02%)

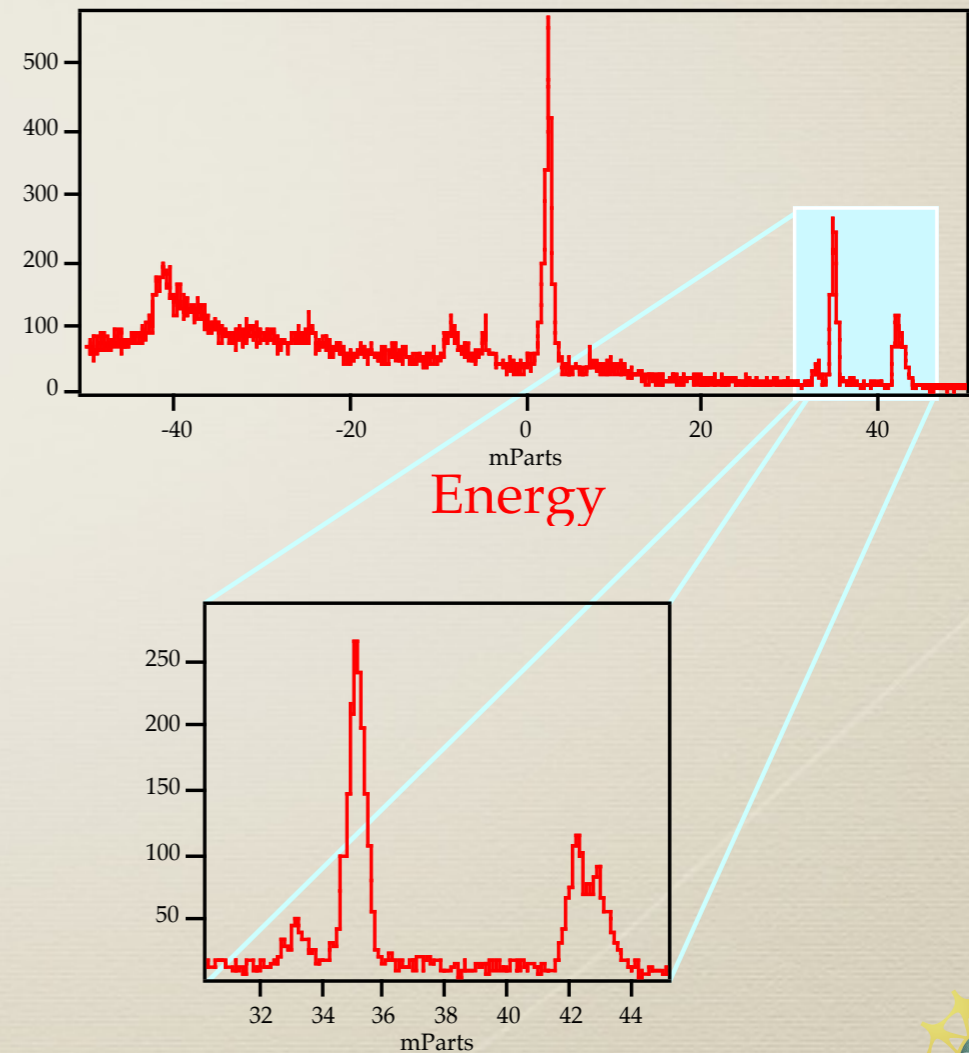
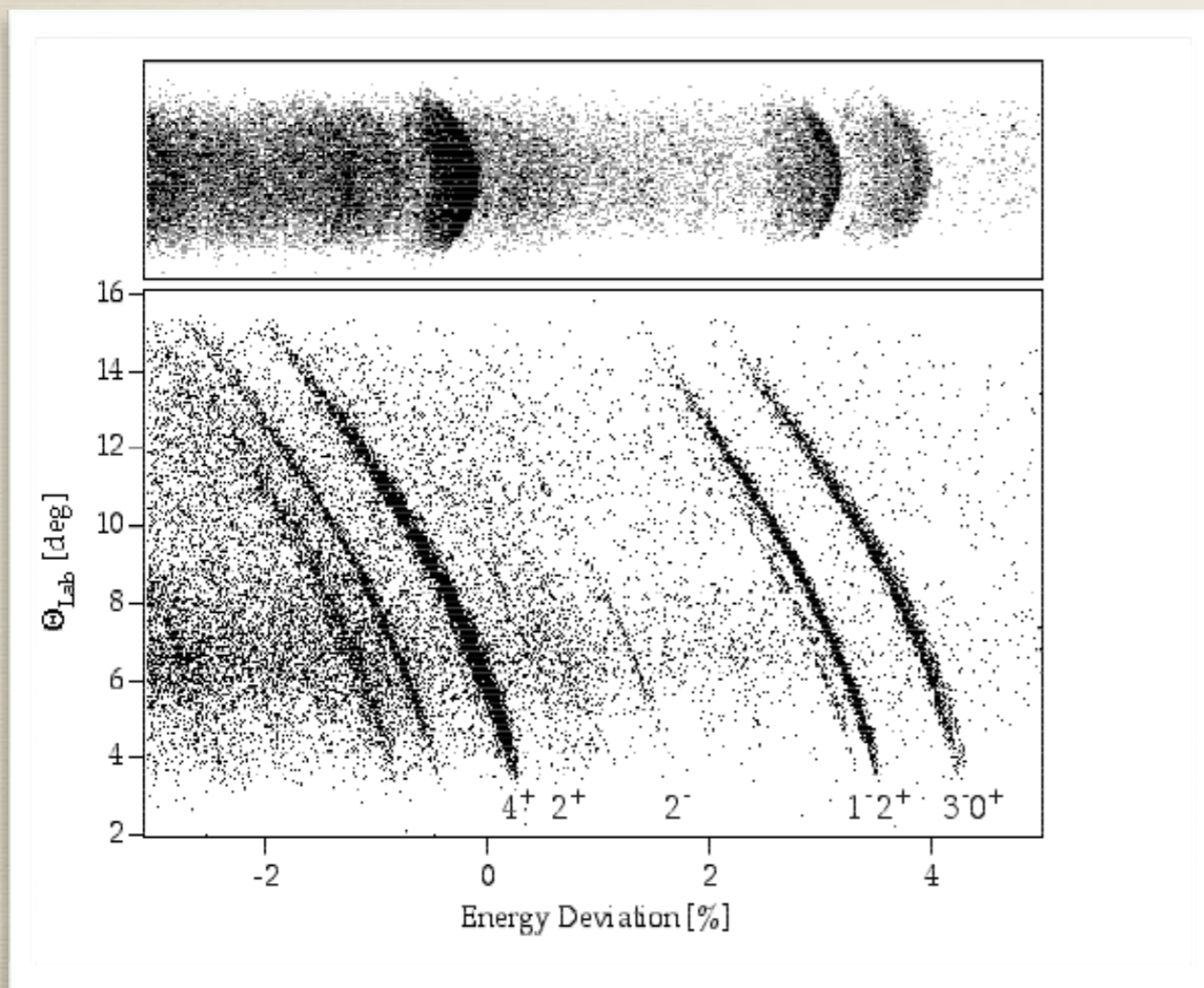
Dispersion matching

- ▶ Analysis line disperses the radioactive beam at target location
- ▶ Without target: all particles refocussed in focal plane
- ▶ Position at focal plane only depends on energy lost in target



Example

- ▶ Reaction: $^{12}\text{C}(^7\text{Li},^3\text{H})^{16}\text{O}$ at 19 MeV/u
- ▶ Spectrograph rotated at 8°
- ▶ Energy resolution of 1/1800 over full acceptance (20 msr)



Particle identification in S800

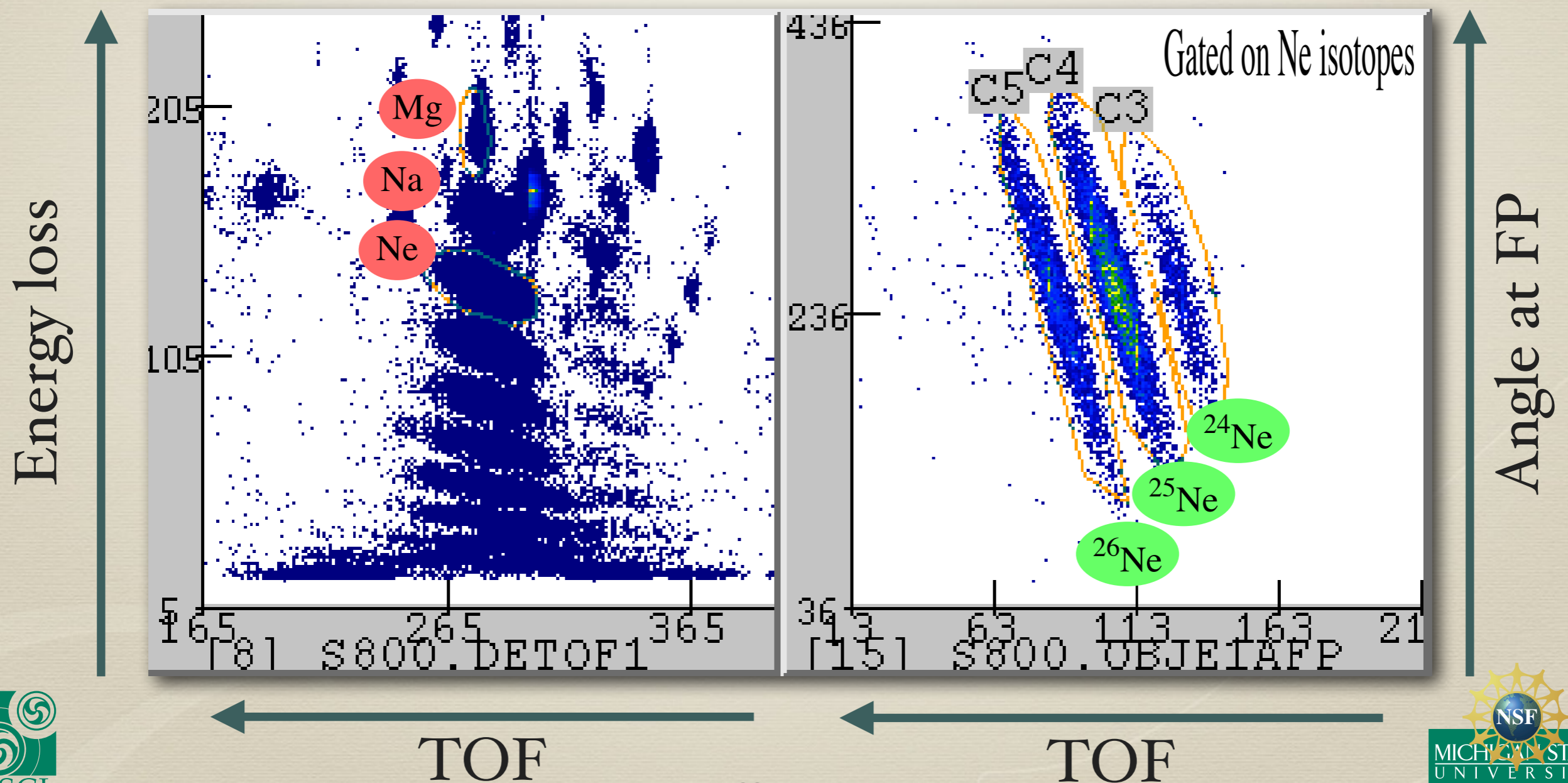
- ▶ How to identify the heavy residues collected in the S800
- ▶ Use time-of-flight and energy loss measurements
- ▶ Time-of-flight depends on velocity and trajectory length
- ▶ Energy loss depends on velocity
- ▶ Corrections are needed to recover resolution and achieve identification

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Dismiss

Example

- ▶ Angle at Focal Plane directly proportional to length of trajectory inside S800
- ▶ Without correction $^{26,25,24}\text{Ne}$ are not resolved



Doppler correction

$$E = E_0 \frac{\sqrt{1 - \beta_0^2}}{1 - \beta_0 \cdot \mathbf{e}}$$

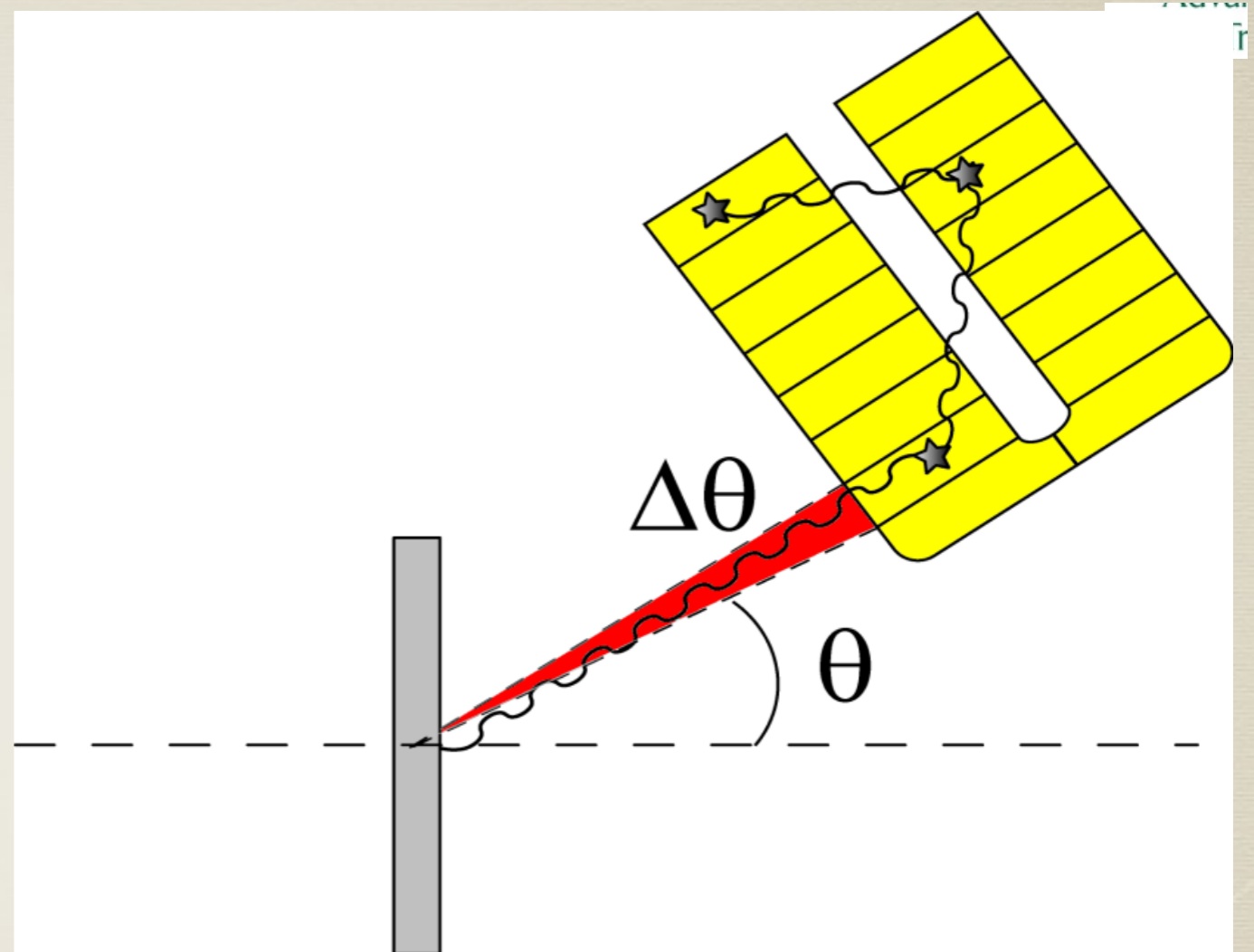
$$\beta_0 \cdot \mathbf{e} = |\beta_0| \cos \theta_0$$

E_0 γ -ray energy in the source frame

E γ -ray energy in the lab frame

β_0 velocity of the source

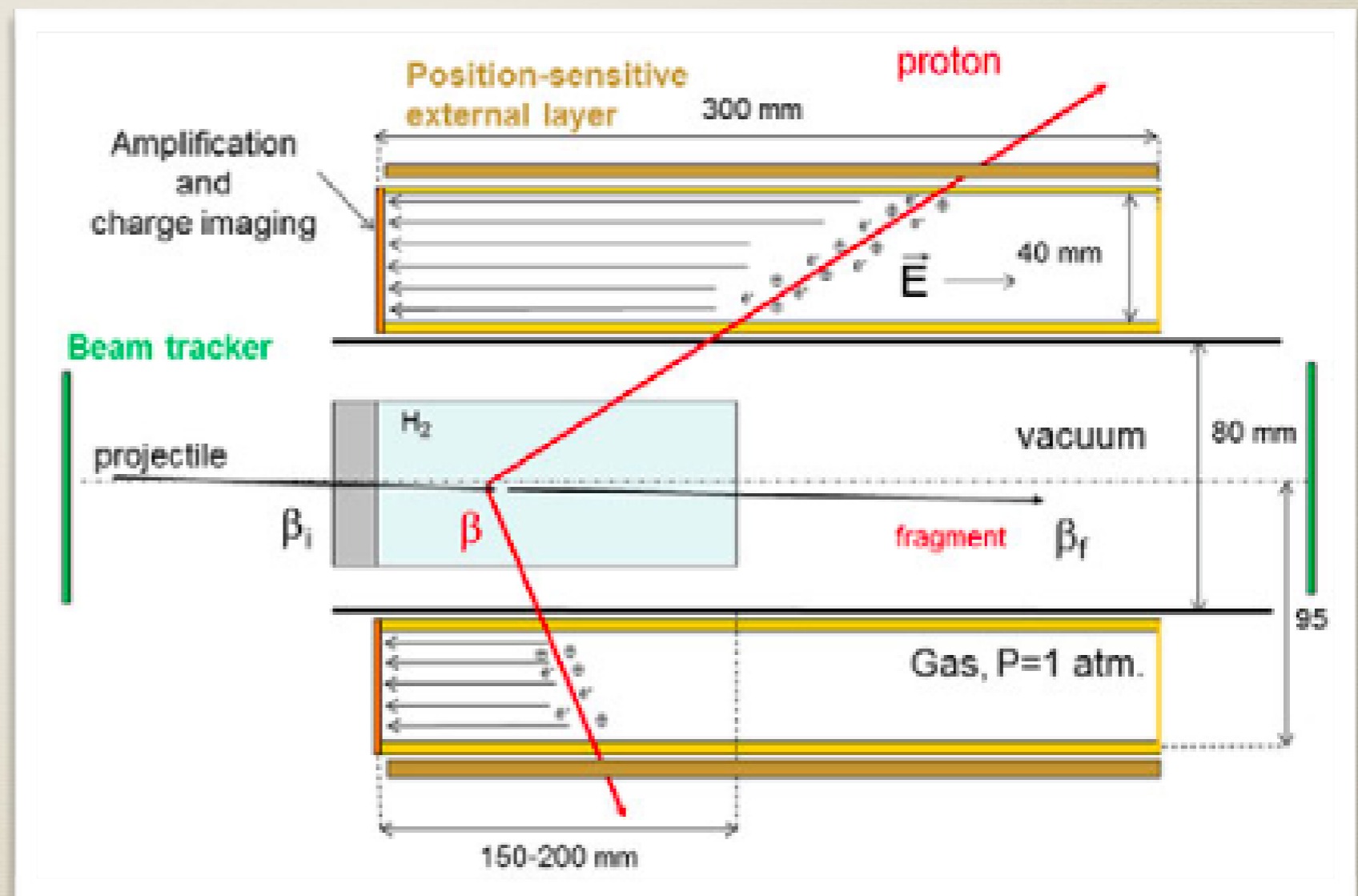
θ_0 γ -ray angle of emission



- ▶ γ -ray energy directly related to velocity at the time of emission
- ▶ Resolution directly related to target thickness

Vertex tracker: MINOS design

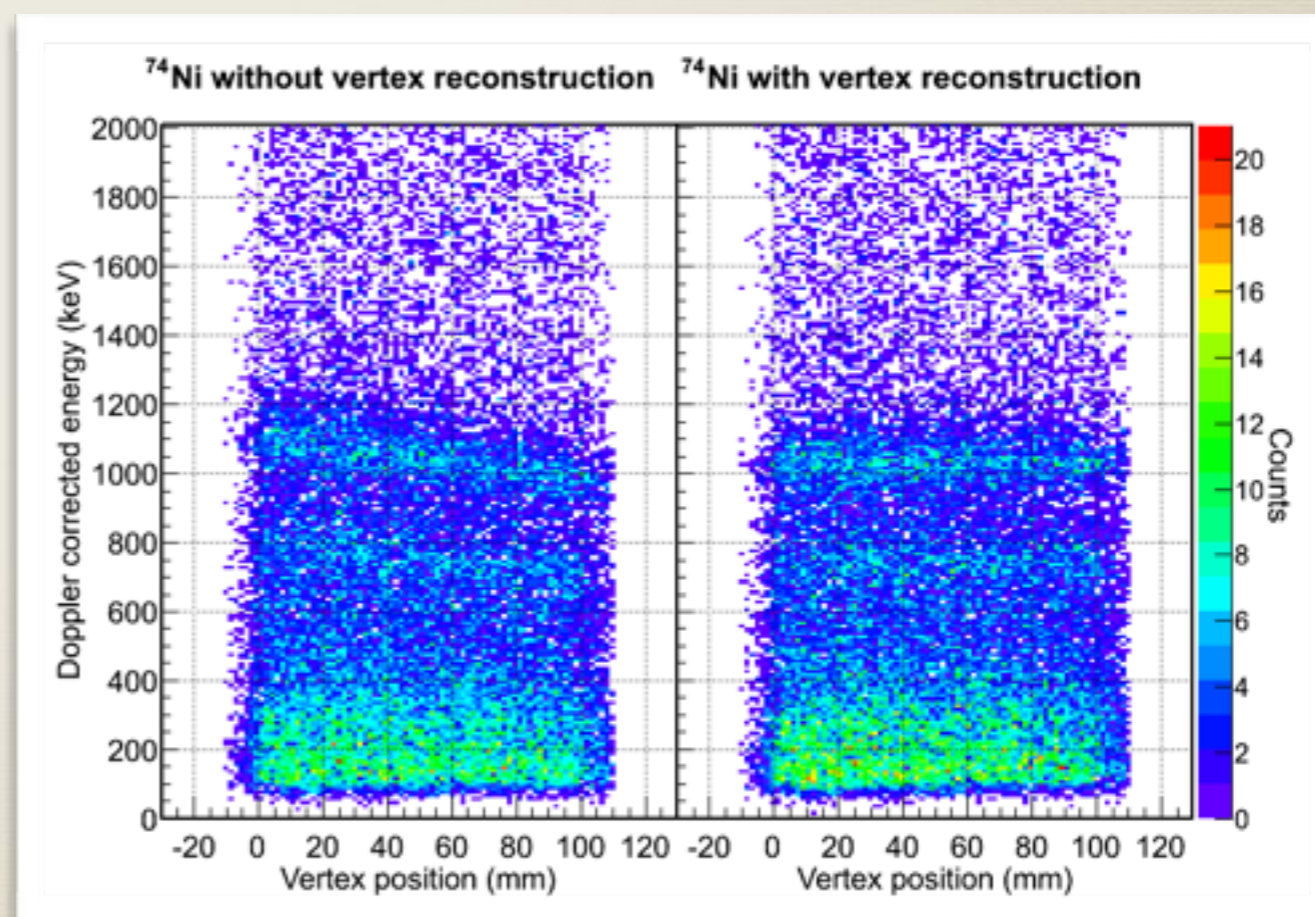
- ▶ In-beam γ -ray spectroscopy relies on Doppler correction
 - ▶ Both γ -ray angle and energy of reaction needed
 - ▶ Vertex tracker measures reaction location inside LH₂ volume
- ▶ (p,2p) or (p,pn) reactions at 200-300 MeV/u
- ▶ Protons can escape target



A. Obertelli et al., EPJA 50, 8 (2014)

MINOS performance

- ▶ Gains in luminosity and resolution
 - ▶ Luminosity increase by a factor 5 to 50
 - ▶ Resolution gain depends greatly on γ -ray array performance
 - ▶ γ -ray tracking Ge array resolution will result in sensitivity gains 100-200
- ▶ Present use @ RIBF/RIKEN
 - ▶ DALI2 NaI array (6-7%)
 - ▶ Example shows clear improvement with vertex determination



Courtesy of A. Obertelli

Take aways...

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 - ▶ Inverse kinematics makes life difficult at low energies, but can be exploited for high luminosity at high energies
 - ▶ Low beam intensities requires to design experiments with high efficiency to recover good luminosity

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- ▶ Experimental developments foster new discoveries (and vice versa)

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- ▶ Don't worry, we will help you (a little...)