Self-Consistent Simulations of Beam and Plasma Systems Homework 1

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Problem 1 - Derivatives on a non-uniform mesh.

In class, we discretized a first derivative on a uniform grid

as:

$$\left. \frac{\partial f}{\partial x} \right|_{i} = \frac{f_{i+1} - f_{i-1}}{2\Delta x} + \mathcal{O}(\Delta x^2)$$

whereas a second derivative is discretized as

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$$\frac{\partial^2 f}{\partial x^2}\Big|_i = \frac{f_{i+1} - 2f_i + f_{i-1}}{(\Delta x)^2} + \mathcal{O}(\Delta x^2)$$

Similarly, discretize the first and second derivative for a non-uniform grid

$$\begin{array}{ccccccccc} f_{i-1} & f_i & & f_{i+1} \\ \hline & & & \\ \hline & & & \\ i-1 & \Delta x_- & i & & \Delta x_+ & & i+1 \end{array}$$

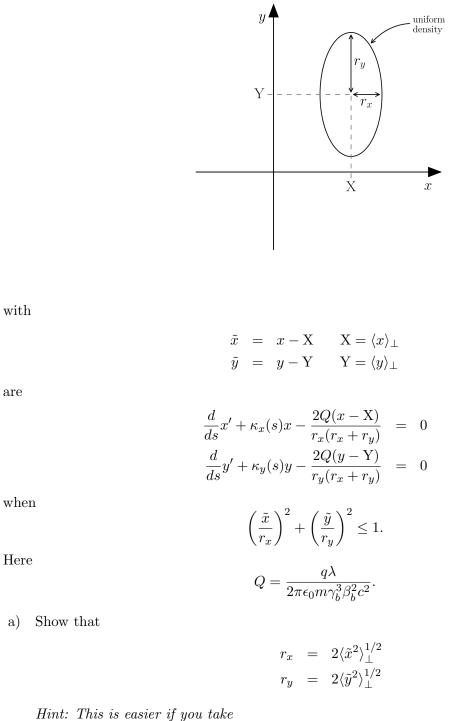
i.e. find:

$$\frac{\partial f}{\partial x}\Big|_{i} = \cdots + \mathcal{O}(\Delta x^{2}_{+,-})$$
$$\frac{\partial^{2} f}{\partial x^{2}}\Big|_{i} = \cdots + \mathcal{O}(\Delta x_{+,-})$$

Notice that we find an $\mathcal{O}(\Delta x_{+,-})$ error term for the second derivative instead of an $\mathcal{O}(\Delta x^2)$ term as we did for the uniform grid. Establish the order of error with your answer!

Problem 2 - Moment formulation of the K-V distribution

The equations of motion for a particle evolving within a K-V beam



$$\int_{ellipse} d\tilde{x} \int d\tilde{y} \cdots = r_x r_y \int_0^1 d\rho \rho \int_{-\pi}^{\pi} d\theta \cdots$$

$$\begin{aligned} \tilde{x} &= r_x \rho \cos \theta \\ \tilde{y} &= r_y \rho \sin \theta \end{aligned}$$

with

$$\rho \in [0,1], \quad \theta \in [-\pi,\pi], \quad and \quad d\tilde{x}d\tilde{y} = r_x r_y \rho d\rho d\theta.$$

b) Derive the equations of motion summarized in class as:

$$\frac{d}{ds} \begin{bmatrix} \langle x \rangle_{\perp} \\ \langle x' \rangle_{\perp} \\ \langle y \rangle_{\perp} \\ \langle y' \rangle_{\perp} \end{bmatrix} = \begin{bmatrix} \langle x' \rangle_{\perp} \\ -\kappa_x(s) \langle x \rangle_{\perp} \\ \langle y' \rangle_{\perp} \\ \langle y' \rangle_{\perp} \\ -\kappa_y(s) \langle y \rangle_{\perp} \end{bmatrix}$$
$$\frac{d}{ds} \begin{bmatrix} \langle \tilde{x}^2 \rangle_{\perp} \\ \langle \tilde{x}^2 \rangle_{\perp} \\ \langle \tilde{x}^2 \rangle_{\perp} \\ \langle \tilde{y}^2 \rangle_{\perp} \\ \langle \tilde{y}^2 \rangle_{\perp} \\ \langle \tilde{y}^2 \rangle_{\perp} \end{bmatrix} = \begin{bmatrix} 2 \langle \tilde{x} \tilde{x}' \rangle_{\perp} \\ \langle \tilde{x}'^2 \rangle_{\perp} - \kappa_x(s) \langle \tilde{x}^2 \rangle_{\perp} + \frac{Q \langle \tilde{x}^2 \rangle_{\perp}^{1/2}}{2[\langle \tilde{x}^2 \rangle_{\perp}^{1/2} + \langle \tilde{y}^2 \rangle_{\perp}^{1/2}]} \\ -2\kappa_x(s) \langle \tilde{x} \tilde{x}' \rangle_{\perp} + \frac{Q \langle \tilde{x}^2 \rangle_{\perp}}{\langle \tilde{x}^2 \rangle_{\perp}^{1/2} + \langle \tilde{y}^2 \rangle_{\perp}^{1/2}} \\ 2 \langle \tilde{y} \tilde{y}' \rangle_{\perp} \\ \langle \tilde{y}'^2 \rangle_{\perp} - \kappa_y(s) \langle \tilde{y}^2 \rangle_{\perp} + \frac{Q \langle \tilde{y}^2 \rangle_{\perp}^{1/2}}{2[\langle \tilde{x}^2 \rangle_{\perp}^{1/2} + \langle \tilde{y}^2 \rangle_{\perp}^{1/2}]} \\ -2\kappa_y(s) \langle \tilde{y} \tilde{y}' \rangle_{\perp} + \frac{Q \langle \tilde{y} \tilde{y}' \rangle_{\perp}}{\langle \tilde{y}^2 \rangle_{\perp}^{1/2} + \langle \tilde{y}^2 \rangle_{\perp}^{1/2}} \end{bmatrix}$$

Moments not shown are zero. $\langle \tilde{x}\tilde{y}\rangle = 0$ is obvious and $\langle \tilde{x}\tilde{y}'\rangle = \langle \tilde{x}'\tilde{y}\rangle = 0$ follows from the kinetic distribution.

c) Show from the results in part b) that

$$\epsilon_x = \text{const.}$$
 with $\epsilon_x = 4[\langle \tilde{x}^2 \rangle_{\perp} \langle \tilde{x}'^2 \rangle_{\perp} - \langle \tilde{x}\tilde{x}' \rangle_{\perp}^2]^{1/2}$

d) Show from the results in part b) and c) that

$$\frac{d^2}{ds^2}r_x + \kappa_x(s)r_x - \frac{2Q}{r_x + r_y} - \frac{\epsilon_x^2}{r_x^3} = 0$$

This is the "standard" form of the K-V envelope equation.

Problem 3 - Moment equations and conservation constraints

The non-relativistic Vlasov-equation is:

$$\left\{\frac{\partial}{\partial t} + \vec{v} \cdot \frac{\partial}{\partial \vec{x}} + \frac{q}{m} \left[\vec{E} + \vec{v} \times \vec{B}\right] \cdot \frac{\partial}{\partial \vec{v}}\right\} f(\vec{x}, \vec{v}, t) = 0$$

Define a fluid density n and a fluid flow velocity \vec{V} by

$$\begin{split} n(\vec{x},t) &= \int d^3 v \, f(\vec{x},\vec{v},t) \\ n(\vec{x},t) \vec{\nabla}(\vec{x},t) &= \int d^3 v \, \vec{v} \, f(\vec{x},\vec{v},t) \end{split}$$

a) Operate on the Vlasov equation with

$$\int d^3v \cdots$$

to derive the continuity equation:

$$\frac{\partial}{\partial t}n(\vec{x},t) + \frac{\partial}{\partial \vec{x}} \cdot \left[n(\vec{x},t)\vec{\mathbf{V}}(\vec{x},t)\right] = 0$$

- b) Can the continuity equation be solved by itself if you specify the initial density field $n(\vec{x}, t = 0)$? Why?
- c) Operate on the Vlasov equation with

$$\int d^3v \, \vec{v} \cdots$$

to derive the fluid force equation.

$$\begin{aligned} \frac{\partial}{\partial t} \left(n \vec{\mathbf{V}} \right) + \nabla \cdot \left(n \langle \vec{v} \vec{v} \rangle_v \right) &= \frac{q}{m} n \left(\vec{E} + \vec{\mathbf{V}} \times \vec{B} \right) \\ \text{with } \langle \vec{v} \vec{v} \rangle_v &\equiv \frac{\int d^3 v \, \vec{v} \vec{v} f}{\int d^3 v \, f} \end{aligned}$$

Defining a pressure tensor as

$$\underline{\underline{P}} = m \int d^3 v \, (\vec{v} - \vec{V}) (\vec{v} - \vec{V}) f(\vec{x}, \vec{v}, t)$$
$$= mn \langle \vec{v}\vec{v} \rangle_v - mn \vec{V} \vec{V},$$

the fluid force equation can be expressed as

$$\frac{\partial}{\partial t}\vec{\mathbf{V}} + \vec{\mathbf{V}} \cdot \frac{\partial}{\partial \vec{x}}\vec{\mathbf{V}} = \frac{q}{m}\left(\vec{E} + \vec{\mathbf{V}} \times \vec{B}\right) - \frac{1}{mn}\frac{\partial}{\partial \vec{x}} \cdot \underline{\underline{\mathbf{P}}}.$$

This form is often used in fluid/plasma analysis.

- d) If the continuity and force equation derived in parts a) and c) are analyzed, can they be solved in principle if you specify the initial density field $n(\vec{x}, t = 0)$ and the velocity field $\vec{V}(\vec{x}, t = 0)$? Why? Does the answer change if we assume a cold initial beam with $\underline{P} = 0$? Why?
- e) Let G(f) be some smooth, differentiable function of f, satisfying $G(f \to 0) = 0$. Show that

$$\int d^3x \int d^3v \, G(f) = \text{const.}$$

This so-called "generalized entropy" measure with G specified can be used to check Vlasov simulations. For example:

$$\begin{split} G(f) &= f: \qquad \int d^3x \int d^3v \, f = \text{const.} \to \text{charge conservation} \\ G(f) &= f^2: \qquad \int d^3x \int d^3v \, f^2 = \text{const.} \to \text{``enstropy'' conservation} \\ G(f) &= f \ln f: \qquad \int d^3x \int d^3v \, f \ln f = \text{const.} \to \text{entropy conservation} \end{split}$$

Problem 4 - Python program for Runge-Kutta order 2

We wish to integrate the differential equation

$$\frac{d x(t)}{dt} = x(t)\cos(t) \qquad x(0) = 1$$

with Euler's method and the Runge-Kutta method (order 2), and compare the results.

a) Download the file euler.py from:

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http://raw.githubusercontent.com/RemiLehe/uspas_exercise/master/euler.py.
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Copy it to a new name rungekutta.py and modify it to implement the Runge-Kutta method, order 2. In particular:

- Change the name of the class from EulerSolver to RKSolver
- Change the name of the euler_integration method to rk_integration, and implement the Runge-Kutta method, order 2 (see the presentation Overview of Basic Numerical Methods)
- b) In ipython, import classes EulerSolver from euler.py and RKSolver from runge_kutta.py. Evaluate the results of both methods for N=100 ; what is the RMS error in both cases?