

## \* Recap from last time

Problems w/ Constraints:

$L(q_i, \dot{q}_i, t)$   $i=1, \dots, S$  generalized coordinates  
plus  $m$  eqns of constraint  $f_k(q_1, \dots, q_s) = 0$   $k=1, \dots, m$   
(e.g., motion confined to some surface)

2 ways to deal w/ these situations

① Use  $f_k(q_1, \dots, q_s, t) = 0$  to express  $m$  of the  $q_i$  in terms of the remaining  $(S-m)$   $q_j$ .

$$\Rightarrow L(q_1, \dot{q}_1, \dots, q_s, \dot{q}_s) \rightarrow L(q_1, \dot{q}_1, \dots, q_{s-m}, \dot{q}_{s-m})$$

$$\Rightarrow \frac{\partial L}{\partial q_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} = 0 \quad j=1, \dots, S-m$$

• EASY 😊

• Can't say anything about  $\vec{F}_{\text{constraint}}$  😞

② Lagrange Multipliers: Keep  $L = L(q_i, \dot{q}_i)$   $i=1, \dots, S$ , but solve

$$\frac{\partial L}{\partial q_i} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} + \sum_{k=1}^m \lambda_k(t) \frac{\partial f_k}{\partial q_i} = 0 \quad i=1, \dots, S$$

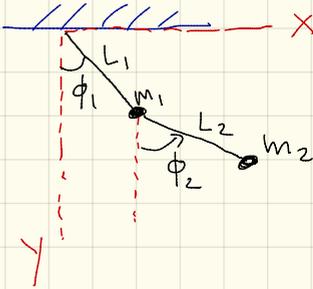
$$f_k(q_1, \dots, q_s, t) = 0 \quad k=1, \dots, m$$

( $S+m$  eqns. for  $S+m$  unknowns  
 $q_1, \dots, q_s + \lambda_1, \dots, \lambda_m$ )

• More eqns ( $S+m$ ) 😞

• Allows you to calculate  $\vec{F}_{\text{constraint}}$  😊

Example: Co-planar double pendulum. Find  $L$



2 masses  $\Rightarrow$  Naively, 6 DOF



Constrained motion

- planar  $z_1, z_2$
- $L_1 + L_2$  fixed  $r_1, r_2$

$\therefore$  Really only 2 independent generalised  
Coords. Needed.

$$\Rightarrow \underline{T(\dot{\phi}_1 + \dot{\phi}_2)}$$

\* Construct  $L(\phi_1, \phi_2, \dot{\phi}_1, \dot{\phi}_2)$ :  $L = (T_1 + T_2) - (U_1 + U_2)$

$$T_1 = \frac{1}{2} m_1 L_1^2 \dot{\phi}_1^2$$

$$T_2 = \frac{m_2}{2} (\dot{x}_2^2 + \dot{y}_2^2)$$

$$U_1 = -m_1 g L_1 \cos \phi_1$$

$$U_2 = -m_2 g y_2$$

$$\text{but } x_2 = L_1 \sin \phi_1 + L_2 \sin \phi_2$$

$$y_2 = L_1 \cos \phi_1 + L_2 \cos \phi_2$$

$$\Rightarrow T_2 = \frac{1}{2} m_2 \left[ L_1^2 \dot{\phi}_1^2 + L_2^2 \dot{\phi}_2^2 + 2L_1 L_2 \cos(\phi_1 - \phi_2) \dot{\phi}_1 \dot{\phi}_2 \right]$$

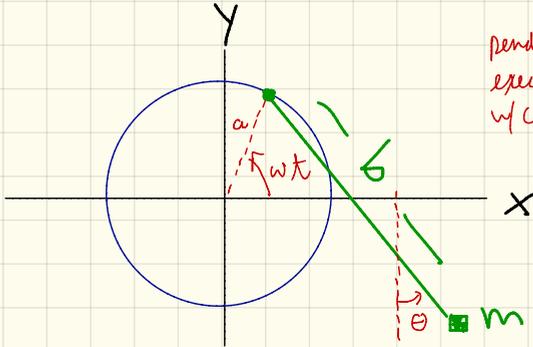
↓

$$L(\phi_1, \dot{\phi}_1, \phi_2, \dot{\phi}_2) = \frac{1}{2} (m_1 + m_2) L_1^2 \dot{\phi}_1^2 + \frac{1}{2} m_2 L_2^2 \dot{\phi}_2^2 + m_2 L_1 L_2 \dot{\phi}_1 \dot{\phi}_2 \cos(\phi_1 - \phi_2) \\ + (m_1 + m_2) g L_1 \cos \phi_1 + m_2 g L_2 \cos \phi_2$$

To find EOM, simply take  $\frac{\partial L}{\partial \phi_i} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\phi}_i} \right) = 0 \quad i=1,2$

Example

↓ g



pendulum support  
executes circular motion  
w/ constant  $\omega$  + radius  $a$

\* Particles in 3D  $\Rightarrow S=3$

\* But planar motion + fixed  $b \Rightarrow$  only  $S-2=1$  independent coordinate  $\Theta$

$$x = a \cos \omega t + b \sin \Theta$$

$$y = a \sin \omega t - b \cos \Theta$$

$$\dot{x} = -a \omega \sin \omega t + b \dot{\Theta} \cos \Theta$$

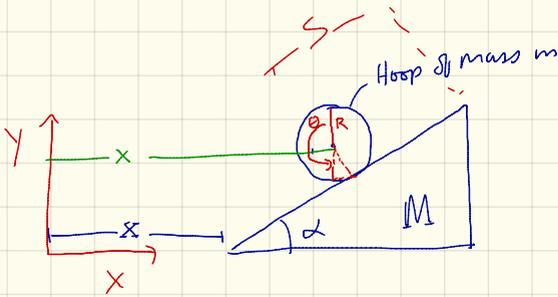
$$\dot{y} = a \omega \cos \omega t + b \dot{\Theta} \sin \Theta$$

$$\Rightarrow T = \frac{m}{2}(\dot{x}^2 + \dot{y}^2) = \frac{m}{2}[a^2\dot{\omega}^2 + b^2\dot{\theta}^2 + 2b\dot{\theta}a\omega \sin(\theta - \omega t)]$$

$$U = mgy = mg(a\sin\omega t - b\cos\theta)$$

$$\Rightarrow L = T - U = \frac{m}{2}[a^2\dot{\omega}^2 + b^2\dot{\theta}^2 + 2b\dot{\theta}a\omega \sin(\theta - \omega t)] - mg(a\sin\omega t - b\cos\theta)$$

### Example (#7.6)



\* Hoop rolls without slipping

$$\Rightarrow S = R\theta$$

\* Wedge can slide w/out friction

### \* Find Lagrange eqns. of motion

- Based on example from last class, expect  $S, \theta, X$  suitable set of generalized coords

$$L = T_{\text{hoop}} + T_{\text{wedge}} - U_{\text{hoop}} - U_{\text{wedge}}$$

$$T_{\text{hoop}} = \frac{m}{2}(\dot{x}^2 + \dot{y}^2) + \frac{I}{2}\dot{\theta}^2$$

$$I = mR^2$$

but

$$x = X + (l - S)\cos\alpha - R\sin\alpha$$

$$y = (l - S)\sin\alpha + R\cos\alpha$$

Quick aside on moment of inertia for  $T_{rot}$  for the hoop



$$dT = \frac{dm}{2} v_{rot}^2 = \frac{dm}{2} \left( \frac{R d\theta}{dt} \right)^2 = \frac{R^2}{2} \dot{\theta}^2 dm$$

$$T_{rot} = \frac{mR^2}{2} \dot{\theta}^2 \equiv \frac{I}{2} \dot{\theta}^2 \Rightarrow I = mR^2$$

Likewise, for the rolling disk example last time



$$dT = \frac{dm}{2} r^2 \dot{\theta}^2$$

$$dm = \sigma 2\pi r dr = \frac{M}{\pi R^2} \cdot 2\pi r dr$$

$$dm = \frac{2M}{R^2} r dr$$

$$\Rightarrow T_{rot} = \int dT_{rot} = \frac{1}{2} \cdot \frac{2M}{R^2} \dot{\theta}^2 \int_0^R r^3 dr$$

$$= \frac{M \dot{\theta}^2 R^2}{4} \equiv \frac{I}{2} \dot{\theta}^2$$

$$\Rightarrow I_{disk} = \frac{1}{2} MR^2$$

$$L = T_{\text{hoop}} + T_{\text{wedge}} - U_{\text{hoop}} - \cancel{U_{\text{wedge}}}$$

$$T_{\text{hoop}} = \frac{m}{2} (\dot{x}^2 + \dot{y}^2) + \frac{mR^2}{2} \dot{\theta}^2$$

but

$$x = X + (l-s) \cos \alpha - R \sin \alpha$$

$$y = (l-s) \sin \alpha + R \cos \alpha$$

↓

$$\dot{x} = \dot{X} - \dot{s} \cos \alpha$$

$$\dot{y} = -\dot{s} \sin \alpha$$

$$\therefore T_{\text{hoop}} = \frac{m}{2} \left( (\dot{X} - \dot{s} \cos \alpha)^2 + \dot{s}^2 \sin^2 \alpha \right) + \frac{mR^2}{2} \dot{\theta}^2$$

$$\ast \text{ but } \underline{s = R\theta} \Rightarrow \dot{\theta} = \frac{\dot{s}}{R}$$

$$\Rightarrow T_{\text{hoop}} = \frac{m}{2} \left[ \dot{X}^2 - 2\dot{X}\dot{s} \cos \alpha + \dot{s}^2 \right] + \frac{m}{2} \dot{s}^2$$

$$\boxed{T_{\text{hoop}} = \frac{m}{2} \left[ \dot{X}^2 - 2\dot{X}\dot{s} \cos \alpha + 2\dot{s}^2 \right]} \quad (2)$$

$$\boxed{U_{\text{hoop}} = mgy = mg[(l-s) \sin \alpha + R \cos \alpha]} \quad (3)$$

$$\therefore \underline{L = (1) + (2) + (3) = m\dot{s}^2 + \frac{1}{2}(m+M)\dot{X}^2 - m\dot{X}\dot{s} \cos \alpha - mgs \sin \alpha + \text{const}}$$

$$L = \textcircled{1} + \textcircled{2} + \textcircled{3} = m\dot{s}^2 + \frac{1}{2}(m+M)\dot{X}^2 - m\dot{X}\dot{s}\cos\alpha - mgs\sin\alpha + \text{const}$$

---

\* Now apply EL eqns.

$$\frac{\partial L}{\partial s} = -mg\sin\alpha$$
$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{s}}\right) = \frac{d}{dt}(2m\dot{s} - m\dot{X}\cos\alpha) = 2m\ddot{s} - m\ddot{X}\cos\alpha$$

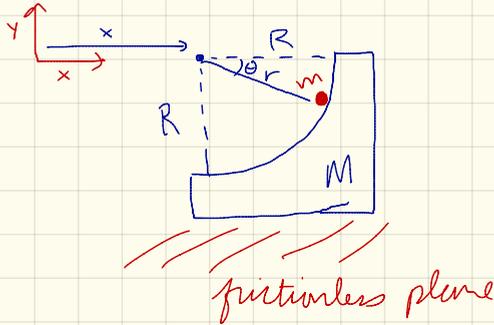
$$\Rightarrow 2m\ddot{s} - m\ddot{X}\cos\alpha + mg\sin\alpha = 0$$

$$\frac{\partial L}{\partial X} = 0$$

$$\frac{\partial L}{\partial \dot{X}} = (m+M)\dot{X} - m\dot{s}\cos\alpha$$

$$\therefore \frac{d}{dt}\left(\frac{\partial L}{\partial \dot{X}}\right) = 0 = (m+M)\ddot{X} - m\ddot{s}\cos\alpha$$

EX: 7.34



Find EOM for  $m + M$  and find reaction force of Wedge on particle  
 $\Rightarrow$  Use Lagrange multiplier

$$\begin{aligned} X_m &= X & x_m &= x + r \cos \theta & \textcircled{X} \\ Y_m &= 0 & y_m &= -r \sin \theta \end{aligned}$$

$$L = \frac{M}{2} \dot{X}_M^2 + \frac{m}{2} (\dot{x}_m^2 + \dot{y}_m^2) - mg Y_m \quad \text{plug in } \textcircled{X}$$

$$\dot{X}_M^2 = \dot{X}^2$$

$$\dot{x}_m^2 + \dot{y}_m^2 = (\dot{x} + \dot{r} \cos \theta - r \dot{\theta} \sin \theta)^2 + (-\dot{r} \sin \theta - r \dot{\theta} \cos \theta)^2$$

etc.

} tedious algebra

$$L = \frac{(M+m)}{2} \dot{X}^2 + \frac{m}{2} (\dot{r}^2 + r^2 \dot{\theta}^2 + 2\dot{x}\dot{r} \cos \theta - 2\dot{x}r\dot{\theta} \sin \theta) + mgr \sin \theta$$

---

Constraint ;  $f(r, \theta, x) = r - R$

$$\Rightarrow \frac{\partial f}{\partial r} = 1$$

$$\frac{\partial f}{\partial \theta} = \frac{\partial f}{\partial x} = 0$$

$\Rightarrow$  Need to solve them

$$\frac{\partial L}{\partial r} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{r}} \right) + \lambda = 0$$

$$\frac{\partial L}{\partial \theta} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) = 0$$

$$\frac{\partial L}{\partial x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = 0$$

plus  $r = R$