

# Energy Conservation Recap

$$W_{12} = \int_1^2 \vec{F} \cdot d\vec{r} = \int_1^2 d\left(\frac{1}{2}mv^2\right) = T_2 - T_1 \quad (T = \frac{1}{2}mv^2)$$

(always true, indep. of type of force  $\vec{F}$ )

\* If  $\vec{F} = -\vec{\nabla}U(\vec{r}, t)$ :

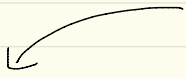
$$W_{12} = T_2 - T_1 = -\int_1^2 dU = U_1 - U_2$$

(\*1) Necessary & Sufficient condition is  $\vec{\nabla} \times \vec{F} = 0$

and

$$\frac{dE}{dt} = \frac{dT}{dt} + \frac{dU}{dt} = (\vec{F} \cdot \vec{\nabla}U) \cdot \dot{\vec{r}} + \frac{\partial U}{\partial t}$$

$$\Rightarrow \frac{dE}{dt} = \frac{\partial U}{\partial t}$$



\* Finally, if  $\vec{F} = U(\vec{r})$ :

$$\frac{d}{dt} E = 0 \Rightarrow E = T + U = \text{constant}$$

"Conservative force  $\vec{F}$ "

Non-conservative forces?

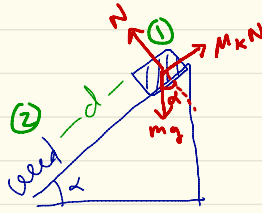
e.g.,  $\vec{F} = \vec{F}_c + \vec{F}_{nc}$  where  $\vec{F}_c = -\vec{\nabla}U^c$   
 $\vec{F}_{nc} \neq -\vec{\nabla}U$

$$\begin{aligned} \Rightarrow W_{12} = T_2 - T_1 &= \int_1^2 \vec{F}_c \cdot d\vec{r} + \int_1^2 \vec{F}_{nc} \cdot d\vec{r} \\ &= U_1^c - U_2^c + W_{12}^{nc} \end{aligned}$$

$$\Rightarrow \Delta T + \Delta U = W_{12}^{nc}$$

Example (HW2):

find  $v_2$ :



$$W_{12} = T_2 - \int_1^2 \vec{F}_1 = \int_1^2 (\vec{F}_g + \vec{F}_{\text{fric}}) \cdot d\vec{r}$$

conservative

Not conservative

$$\therefore \frac{1}{2} m v_2^2 = U_1^g - U_2^g + W_{12}^{\text{fric}}$$

$$= m g d \sin \alpha - \mu_k N d$$

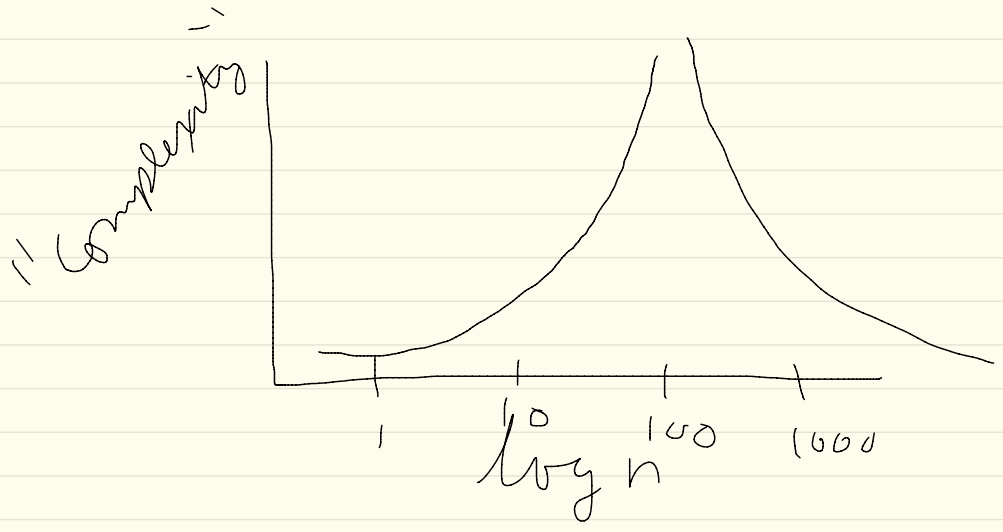
$$\frac{1}{2} m v_2^2 = m g d \sin \alpha - \mu_k m g \cos \alpha d$$

$$\Rightarrow v_2 = \sqrt{2 g d [\sin \alpha - \mu_k \cos \alpha]}$$

# Dynamics of a System of Particles (Reading: ch. 9)

\* Exact solutions of the eqns. of motion can almost always be done (sometimes analytically, other times numerically) for  $n=1$  or  $2$  bodies

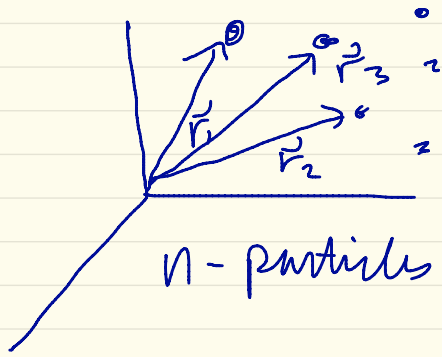
\* Beyond this, problems can only (typically) be solved on a computer. The complexity/computational cost (in CM as well as QM) looks like



- \* large  $n \Rightarrow$  statistical methods simplify the problem
- \* small  $n \Rightarrow$  analytic sol'n or "cheap" numerical sol'n
- \* medium  $n \Rightarrow$  HARD

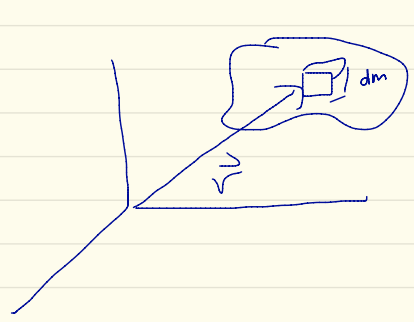
Nevertheless, we can often times make general statements about a system of  $n$ -bodies by appealing to conservation laws, etc., even if the direct solution of the eqns. of motion is impossible/impractical.

# Center of Mass & total $\vec{P}$ of a System



\* let  $M = \sum_{\alpha=1}^n m_{\alpha}$  (total mass) and  $\vec{R} = \frac{\sum_{\alpha} m_{\alpha} \vec{r}_{\alpha}}{M}$  "Center of Mass"

For a continuous distribution of mass



$$\vec{R} = \frac{\int \vec{r} dm}{\int dm}$$

density  $\rho(\vec{r})$ :  $dm = \rho(\vec{r}) d^3r$

$$\Rightarrow \vec{R} = \frac{\int \vec{r} \rho(\vec{r}) d^3r}{\int \rho(\vec{r}) d^3r}$$

\* Obvious generalization for 2d (i.e. surface) or 1d (linear) continuous mass distributions

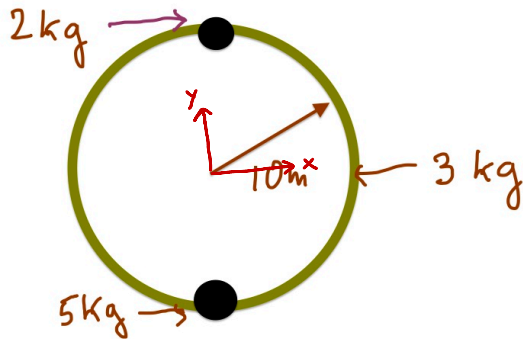
2d:  $\sigma(\vec{r}) dA = dm \Rightarrow \vec{R} = \frac{\int \vec{r} \sigma(\vec{r}) dA}{\int \sigma(\vec{r}) dA}$

1d:  $\lambda(\vec{r}) dl = dm \Rightarrow \vec{R} = \frac{\int \vec{r} \lambda(\vec{r}) dl}{\int \lambda(\vec{r}) dl}$

## Clicker Question

L7-5

A THIN UNIFORM RING OF RADIUS  $10\text{m}$  AND MASS  $3\text{kg}$  HAS TWO POINT MASSES OF  $2\text{kg}$  AND  $5\text{kg}$  ATTACHED ON THE OPPOSITE SIDES OF ITS CIRCUMFERENCE. HOW FAR IS THE CM OF THE SYSTEM AWAY FROM THE CENTER OF THE RING?



A)  $0.0\text{m}$

B)  $2.0\text{m}$

C)  $3.0\text{m}$

D)  $4.3\text{m}$

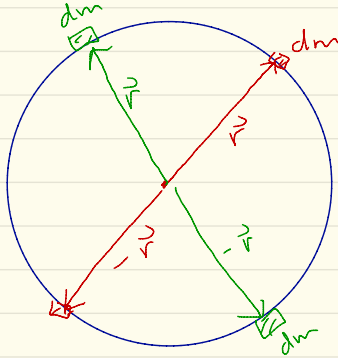
E)  $7.0\text{m}$

Solution:

$$\vec{R} = \frac{\sum_{\alpha=1}^2 \vec{r}_{\alpha} m_{\alpha}}{\sum m_{\alpha}} + \int_{\text{ring}} \vec{r} dm / M_{\text{ring}}$$

L7-6

For the ring



$$\vec{r} dm - \vec{r} dm = 0$$

(cancel pairwise)

$$\Rightarrow \int_{\text{ring}} \vec{r} dm = 0 \text{ by circular symmetry}$$

$$\therefore \vec{R} = \frac{\vec{r}_1 m_1 + \vec{r}_2 m_2}{m_1 + m_2} = \frac{(2\text{kg})(10\text{m})\hat{j} + (5\text{kg})(-10\text{m})\hat{j}}{7\text{kg}} = \frac{-30}{7} \hat{j}$$

$$\therefore \underline{R \approx 4.3 \text{ m}}$$

\* Linear Momentum of the system (n-bodies)

\* let  $\vec{F}_{\alpha\beta}$  = force on particle  $\alpha$  by particle  $\beta$

$$\vec{F}_{\alpha\beta} = -\vec{F}_{\beta\alpha} \quad (\text{Newton's 3rd Law})$$

\* let  $\vec{F}_\alpha$  = total force on particle  $\alpha$

$$= \vec{F}_\alpha^{\text{ext}} + \vec{F}_\alpha$$

external force  
originating from  
outside the n-body  
system

sum of forces from  
all the other (n-1)  
bodies of the system

$$\vec{F}_\alpha = \sum_{\beta} \vec{F}_{\alpha\beta}$$

\* Newton's 2nd =>

$$\vec{F}_\alpha = \frac{d^2}{dt^2} (m_\alpha \vec{r}_\alpha) = \vec{F}_\alpha^{\text{ext}} + \sum_{\beta} \vec{F}_{\alpha\beta}$$

$$\therefore \sum_{\alpha} \vec{F}_\alpha = \sum_{\alpha} \vec{F}_\alpha^{\text{ext}} + \sum_{\substack{\alpha, \beta \\ (\alpha \neq \beta)}} \vec{F}_{\alpha\beta}$$

\* but  $\sum_{\substack{\alpha, \beta \\ (\alpha \neq \beta)}} \vec{F}_{\alpha\beta} = \sum_{\alpha < \beta} (\vec{F}_{\alpha\beta} + \vec{F}_{\beta\alpha}) = 0$  (Newton's 3rd =>  $\vec{F}_{\alpha\beta} = -\vec{F}_{\beta\alpha}$ )

$$\Rightarrow \frac{d^2}{dt^2} \sum_{\alpha} m_\alpha \vec{r}_\alpha = \sum_{\alpha} \vec{F}_\alpha^{\text{ext}} \equiv \vec{F}$$

\* but  $\sum_{\alpha} m_\alpha \ddot{\vec{r}}_\alpha = M \ddot{\vec{R}} \Rightarrow \boxed{M \ddot{\vec{R}} = \vec{F}}$  \*\*

So then, provided the internal forces obey Newton's 3<sup>rd</sup> Law ( $\vec{f}_{ij} = -\vec{f}_{ji}$ )

$$M \ddot{\vec{R}} = \sum_i \vec{F}_i^{\text{ext}} \equiv \vec{F}$$

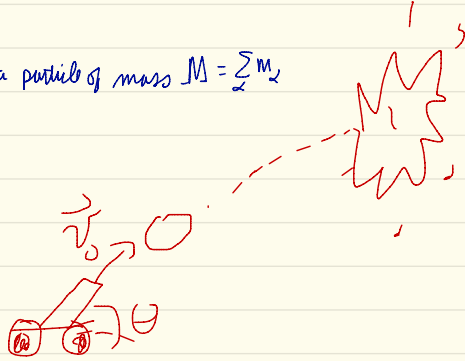
$$\text{def } \dot{\vec{P}} = \sum_i m_i \dot{\vec{v}}_i = \sum_i \vec{p}_i \Rightarrow \vec{F} = \dot{\vec{P}} = M \ddot{\vec{R}}$$



1) COM moves as if it were a particle of mass  $M = \sum_i m_i$

2) Total  $\dot{\vec{P}} = \text{const}$  if  $\vec{F} = 0$

\*Example: Exploding artillery shell



$$M \ddot{\vec{R}} = \sum_i m_i \vec{g} = M \vec{g}$$

$\Rightarrow$  same as a point mass  
(see, e.g., ex 2.6)

$\Rightarrow$  COM  $\vec{R}$  follows parabolic traj. just as in ex 2.6.