

What happened here physically?



Which moments should change?

I_{xx} , I_{zz} , but I_{yy} is physically the same!

→ "parallel" axis theorem / Steiner's theorem.

The MOI around a given axis \vec{a} is equal to the parallel axis through the COM plus the MOI of the body w.r.t. \vec{a} , treated as a point mass (sounds awkward)

$$\hat{I}_{ab} = \left[\int dm \left(r^2 \delta_{ab} - \vec{x}_a \cdot \vec{x}_b \right) \right]$$

$$= \sum_i \Delta m_i \left(r_i^2 \delta_{ab} - \vec{x}_{ia} \cdot \vec{x}_{ib} \right) = \sum_i \Delta m_i \left((\vec{r}_i + \vec{R})^2 \delta_{ab} - (\vec{x}_{ia} + \vec{R}_a) \cdot (\vec{x}_{ib} + \vec{R}_b) \right)$$

$$= \sum_i \Delta m_i \left(\vec{r}_i^2 \delta_{ab} + \vec{R}^2 \delta_{ab} + 2\vec{r}_i \cdot \vec{R} \delta_{ab} - \vec{x}_{ia} \cdot \vec{x}_{ib} - \vec{x}_{ia} \cdot \vec{R}_b - \vec{x}_{ib} \cdot \vec{R}_a \right)$$

$$= \sum_i \Delta m_i \left(\vec{r}_i^2 \delta_{ab} - \vec{x}_{ia} \cdot \vec{x}_{ib} + \vec{R}^2 \delta_{ab} - \vec{R}_a \cdot \vec{R}_b + 2\vec{r}_i \cdot \vec{R} \delta_{ab} - \vec{x}_{ia} \cdot \vec{R}_b - \vec{x}_{ib} \cdot \vec{R}_a \right)$$

$$= \hat{I}_{ab}^1 + M \left(\vec{R}^2 \delta_{ab} - \vec{R}_a \cdot \vec{R}_b \right) + M \left(\vec{R}^2 \delta_{ab} - 2\vec{R}_a \cdot \vec{R}_b - M \vec{x}_a^T \vec{x}_b - M \vec{x}_a \cdot \vec{x}_b^T \right) \\ = \hat{I}_{ab}^1 + M \left(\vec{R}^2 \delta_{ab} - \vec{x}_a \cdot \vec{x}_b \right)$$

Now that we have

$$L = \vec{T} = \frac{1}{2} \vec{\omega}^T \vec{I} \vec{\omega}$$

let's look at E_{CM}/s .

From Newton

(will do Lagrange in ~~lecture notes~~ - it's a bit
numerical b/c Euler angles)

lab/inertial frame.

$$\frac{d\vec{P}}{dt} = \vec{F}$$

\vec{P}, \vec{E} total mom. / ang. mom.

$$\frac{d\vec{\Sigma}}{dt} = \vec{N}$$

\vec{P}, \vec{N} total force / torque

$$\vec{L} = \sum_{i=1}^A \vec{r}_i \times \vec{p}_i = \sum_{i=1}^A m_i \vec{r}_i \times \vec{v}_i$$

body-fixed (same origin)

~~Body Center~~ \Rightarrow $\vec{r}_i = 0$ (\vec{r}_i rotates with body)

$$\dot{\vec{r}}_i = \left(\frac{d\vec{r}_i}{dt} \right)_{BF} + \vec{\omega} \times \vec{r}_i =$$

$$\Rightarrow \vec{L} = \sum_i m_i (\vec{r}_i \times (\vec{\omega} \times \vec{r}_i)) = \sum_i m_i (r_i^2 \vec{\omega} - (\vec{r}_i \cdot \vec{\omega}) \vec{r}_i)$$

$$\Rightarrow L_a = \sum_{i=1}^A \sum_{b=1}^3 m_i (r_i^2 \delta_{ab} - x_{ib} \omega \times x_{ia}) w_b = \sum_b \hat{I}_{ab} w_b$$

$$\Rightarrow \vec{L} = \hat{I} \vec{\omega}$$

• \vec{L} is in general not parallel to $\vec{\omega}$!!

• It will be for rotation around a principal axis
(or with appropriate symmetries; e.g. for a
cylindrical rigid body.)

Now

$$\left(\frac{d\vec{L}}{dt} \right)_{LF} = \left(\frac{d\vec{L}}{dt} \right)_{BF} + \vec{\omega} \times \vec{L}$$

$$= \hat{I} \dot{\vec{\omega}} + \vec{\omega} \times \vec{L} = \hat{I} \dot{\vec{\omega}} + \vec{\omega} \times (\hat{I} \vec{\omega}) = \vec{N}$$

\hat{I} does not depend on time in the BF
fixed frame, but it will in a lab frame
(since rot. matrix to BF frame depends on
time - dep. angles!)

~~Derive Euler angles from Euler eqns~~

Then we have

$$(\overset{x}{\vec{\omega}})_{BF} + \vec{\omega} \times \vec{L} = \vec{N}$$

$$\Rightarrow \vec{\omega} = \omega_x \hat{e}_x + \omega_y \hat{e}_y + \omega_z \hat{e}_z$$

$$\vec{L} = A \omega_x \hat{e}_x + B \omega_y \hat{e}_y + C \omega_z \hat{e}_z$$

(with $\vec{\omega}$ components
as shown before/
earlier)

$$\Rightarrow N_x = A \dot{\omega}_x + (C-B) \omega_y \omega_z$$

$$N_y = B \dot{\omega}_y + (A-C) \omega_x \omega_z$$

$$N_z = C \dot{\omega}_z + (B-A) \omega_y \omega_x$$

Euler's eqns for rigid body (in the body-fixed frame)
principal axis

General procedure:

- first determine external forces that give torque \vec{N} in lab frame
- transform to BF frame (ideally ~~not~~ principal axis system) using rotational matrices with Euler angles.

Free Top

$\vec{N} = 0$ \rightarrow Note that this need not mean that there are no forces! e.g. if Col fixed in lab frame

$$\vec{M} \vec{R} \times \vec{g} = 0$$

$$\rightarrow A \dot{\omega}_x + (C-B) \omega_y \omega_z = 0$$

$$B \dot{\omega}_y + (A-C) \omega_x \omega_z = 0$$

$$C \dot{\omega}_z + (B-A) \omega_y \omega_x = 0$$

Conserved quantities:

multiply each equation by the w_i and add:

$$\begin{aligned} A\ddot{\omega}_x w_x + (\underline{C-B})\dot{\omega}_y \dot{\omega}_z w_x \\ + B\ddot{\omega}_y w_y + (\underline{A-C})\dot{\omega}_x \dot{\omega}_z w_y \\ + C\ddot{\omega}_z w_z + (\underline{B-A})\dot{\omega}_x \dot{\omega}_y w_z = 0 \\ = A\ddot{\omega}_x w_x + B\ddot{\omega}_y w_y + C\ddot{\omega}_z w_z = \frac{d}{dt} \frac{1}{2} (A\omega_x^2 + B\omega_y^2 + C\omega_z^2) \\ = \frac{d}{dt} \vec{r}_{\text{tot}} = 0 \quad \left(\begin{array}{l} \text{= total energy if we have no} \\ \text{translation or forces} \end{array} \right) \end{aligned}$$

Multiply by $A\dot{\omega}_x$, $B\dot{\omega}_y$, $C\dot{\omega}_z$:

$$\begin{aligned} \cancel{\text{if}} \quad A^2 \ddot{\omega}_x \dot{\omega}_x + A(C-B)\dot{\omega}_y \dot{\omega}_z \dot{\omega}_x \\ + B^2 \ddot{\omega}_y \dot{\omega}_y + B(A-C)\dot{\omega}_x \dot{\omega}_z \dot{\omega}_y \\ + C^2 \ddot{\omega}_z \dot{\omega}_z + C(B-A)\dot{\omega}_x \dot{\omega}_y \dot{\omega}_z = 0 \\ = A^2 \ddot{\omega}_x \dot{\omega}_x + B^2 \ddot{\omega}_y \dot{\omega}_y + C^2 \ddot{\omega}_z \dot{\omega}_z = 0 \\ = \frac{d}{dt} \cancel{\text{if}} \frac{1}{2} (A^2 \omega_x^2 + B^2 \omega_y^2 + C^2 \omega_z^2) = \frac{d}{dt} (\frac{1}{2} \vec{L}^2) \\ \Rightarrow \frac{d(\vec{L}^2)}{dt} = 0 \end{aligned}$$

$$\begin{aligned} \vec{L} &= \hat{I} \vec{\omega} \\ \vec{L}^2 &= (\hat{I} \vec{\omega})^2 \\ \vec{\omega}^T \hat{I}^{-1} \hat{I} \vec{\omega} &= \underbrace{\hat{I}^2}_{A^2 B^2 C^2} \end{aligned}$$

→ absolute value of \vec{L} is conserved.

To preserve \vec{L} direction:

$$A\ddot{\omega}_x = B\ddot{\omega}_y = C\ddot{\omega}_z = 0$$

$$\hookrightarrow (C-B)\dot{\omega}_y \dot{\omega}_z = 0 \quad \circ \text{ If } A=B=C \text{ that's trivially true}$$

$$(A-C)\dot{\omega}_x \dot{\omega}_z = 0 \quad \circ \text{ If } A \neq B \neq C: \omega_x = 0, \omega_y = 0$$

$$(B-A)\dot{\omega}_x \dot{\omega}_y = 0 \quad \rightarrow \vec{L} \text{ is on principal axis}$$

Okay

Stability:

Let's assume rotation around z' axis, plus perturbation

$$\omega_{z1} = \omega_0 + \epsilon_{z1}$$

$$\omega_{x1} = \epsilon_{x1}$$

$$\omega_{y1} = \epsilon_{y1}$$

$$\Rightarrow A \dot{\epsilon}_{x1} + (C-B) \epsilon_{y1} (\omega_0 + \epsilon_{x1}) = 0$$

$$B \ddot{\epsilon}_{y1} + (A-C) \epsilon_{x1} (\omega_0 + \epsilon_{x1}) = 0$$

$$C \dot{\epsilon}_{z1} + \underbrace{(B-A) \epsilon_{y1} \epsilon_{x1}}_{=0(\epsilon^2)} = 0 \Rightarrow C \dot{\epsilon}_{z1} = 0 \Rightarrow \dot{\epsilon}_z = \text{const.}$$

$$\Rightarrow B \ddot{\epsilon}_{y1} + (A-C) \dot{\epsilon}_{x1} (\omega_0 + \epsilon_{x1}) = 0 \Rightarrow \dot{\epsilon}_{x1} = - \frac{B}{(A-C) \omega_0} \ddot{\epsilon}_{y1}$$

in first eqn.

$$-\frac{AB}{(A-C)\omega_0} \ddot{\epsilon}_{y1} + (C-B) \epsilon_{y1} = 0$$

$$\ddot{\epsilon}_{y1} - \frac{(A-C)(C-B)}{AB} \omega_0^2 \epsilon_{y1} = 0$$

Analogously,

$$A \ddot{\epsilon}_{x1} + (C-B) \dot{\epsilon}_{y1} \omega_0 = 0 \Rightarrow \dot{\epsilon}_{y1} = - \frac{A}{(C-B)\omega_0} \ddot{\epsilon}_{x1}$$

in second eq.

$$\Rightarrow -\frac{AB}{(C-B)\omega_0} \ddot{\epsilon}_{x1} + (A-C) \omega_0 \epsilon_{x1} = 0$$

$$\ddot{\epsilon}_{x1} - \frac{(A-C)(C-B)}{AB} \omega_0^2 \epsilon_{x1} = 0$$

Define

$$\Sigma^2 = \frac{(A-C)(B-C)}{AB}$$

$$\ddot{\epsilon}_{x1} + \Sigma^2 \epsilon_{x1} = 0$$

$$\ddot{\epsilon}_{y1} + \Sigma^2 \epsilon_{y1} = 0$$

stable solutions for $\Sigma^2 > 0$,
 so $A > C, B > C$ or
 $A < C \& B < C$ (principal)
 i.e., C is smallest or largest MoI

for $B < C < A$ or $A < C < B$ rotation is unstable.

Intermediate Axis Theorem

Free symmetric top:

$$A=B \neq C$$

$$\Rightarrow A\ddot{\omega}_{x1} + (C-A)\omega_{y1}\omega_{z1} = 0 \quad (1)$$

$$A\ddot{\omega}_{y1} + (A-C)\omega_{x1}\omega_{z1} = 0 \quad (2)$$

$\underline{\omega_{01}}$?

$$C\ddot{\omega}_{z1} = 0 \Rightarrow \omega_{z1} = \text{const} \equiv \omega_{0z1}$$

$$\Rightarrow A\ddot{\omega}_{x1} + (C-A)\omega_{0z1}\omega_{y1} = 0 \quad (1)$$

$$A\ddot{\omega}_{y1} + (A-C)\omega_{0z1}\omega_{x1} = 0 \quad (2)$$

$$\Rightarrow A\ddot{\omega}_{x1} + (C-A)\omega_{0z1}\omega_{y1} = 0 \Rightarrow \ddot{\omega}_{y1} = -\frac{A}{(C-A)\omega_{0z1}}\ddot{\omega}_{x1}$$

in (1)

$$\cancel{A\ddot{\omega}_{x1}} + \cancel{(C-A)\omega_{0z1}\omega_{y1}} = 0$$

$$\frac{A^2}{(A-C)}\omega_{0z1}\ddot{\omega}_{x1} + (A-C)\omega_{0z1}\omega_{x1} = 0$$

$$\ddot{\omega}_{x1} + \frac{(A-C)^2}{A^2}\omega_{0z1}^2\omega_{x1} = 0$$

Analogously

$$\ddot{\omega}_{y1} + \frac{(C-A)^2}{A^2}\omega_{0z1}^2\omega_{y1} = 0$$

$\equiv \Sigma^2$

$$\Rightarrow \omega_{x1}(t) = \omega_{0z1} \sin(\Omega t + \beta_1)$$

$$\omega_{y1}(t) = \omega_{0z1} \cos(\Omega t + \beta_1)$$

since ω_x and ω_y are related.

ω_{0z1}, β must be the same

$\Rightarrow \omega_{0z1}$ is const.

$$\omega = \sqrt{\omega_{0z1}^2 + \omega_{0x1}^2} = \text{const.}$$

$\vec{\omega}_L(t) = (\omega_{x1}(t), \omega_{y1}(t))$ is a circle in the xy plane

BF