

Poisson Brackets

$$\{f, g\} = \sum_k \left(\frac{\partial f}{\partial q_k} \frac{\partial g}{\partial p_k} - \frac{\partial g}{\partial q_k} \frac{\partial f}{\partial p_k} \right)$$

Properties of the Brackets will be proven in HW.

- $\{f, g\} = -\{g, f\}$
- $\{\alpha f + \beta g, h\} = \alpha \{f, h\} + \beta \{g, h\}$ for $\alpha, \beta \in \mathbb{R}$
- $\{f, g, h\} = f \{g, h\} + \{f, h\} g = -\{h, fg\}$
- Jacobi: $\{f, \{g, h\}\} + \{g, \{h, f\}\} + \{h, \{f, g\}\} = 0$

We also have

- $\{q_i, q_j\} = 0$, $\{p_i, p_j\} = 0$, $\{q_i, p_j\} = \delta_{ij}$ ← Lock families?
- $\frac{df}{dt} = \{f, H\} + \frac{\partial f}{\partial t}$ - Hamiltonian generates time evolution.
- $\frac{dH}{dt} = \underbrace{\{H, H\}}_{=0, \text{trivially}} + \frac{\partial H}{\partial t} \Rightarrow$ ~~f(H)~~ $\{f, H\} = 0 \Rightarrow \frac{df}{dt} = \frac{\partial f}{\partial t}$
~~f(p,q)~~ with $\{f, H\} = 0 \Rightarrow \frac{df}{dt} = 0$.
~~H~~ $H = \text{const.}$ if $\frac{\partial H}{\partial t} = 0$

Example: Angular Momentum

$$L_i = \sum_{jk} \epsilon_{ijk} x_j p_k$$

$$\begin{aligned} \Rightarrow \{L_i, L_j\} &= \sum_{jk} \sum_{st} \epsilon_{ijk} \epsilon_{rst} \{x_j p_k, x_s p_t\} \\ &= \sum_{jk} \sum_{st} \epsilon_{ijk} \epsilon_{rst} \left(x_j \{p_k, x_s p_t\} + \{x_j, x_s p_t\} p_k \right) \\ &= \sum_{jk} \sum_{st} \epsilon_{ijk} \epsilon_{rst} \left(x_j \left(-x_s \overbrace{\{p_k, p_t\}}^{=0} - \{x_s, p_k\} p_t \right) \right. \\ &\quad \left. + \left(-x_s \{p_t, x_j\} - \underbrace{\{x_s, x_j\}}_{=0} p_t \right) p_k \right) \\ &= \sum_{jk} \sum_{st} \epsilon_{ijk} \epsilon_{rst} \left(-x_j p_t \delta_{ks} + x_s \delta_{jt} p_k \right) \\ &= \sum_{jk} \sum_{st} \epsilon_{ijk} \epsilon_{rst} \left(x_s p_k \delta_{jt} - x_j p_t \delta_{ks} \right) \\ &= \sum_{jks} \epsilon_{ijk} \epsilon_{rsj} x_s p_k - \sum_{jkt} \epsilon_{ijk} \epsilon_{rkt} x_j p_t \\ &= \sum_{jks} \epsilon_{kij} \epsilon_{rsj} x_s p_k - \sum_{jkt} \epsilon_{ijk} \epsilon_{brk} x_j p_t \end{aligned}$$

$$= \sum_{ks} (\delta_{kr} \delta_{is} - \delta_{ks} \delta_{ir}) x_s p_k - \sum_{jt} (\delta_{it} \delta_{jr} - \delta_{ir} \delta_{jt}) x_j p_t$$

$$= \cancel{x_i p_r} - \sum_s (x_s p_s) \delta_{ir} - x_r p_i + \sum_t (x_t p_t) \delta_{ir} = x_i p_r - x_r p_i$$

Make the example explicit

$$L_1 = x_2 p_3 - x_3 p_2 \quad \text{etc.}$$

$$\{L_1, L_2\} = \{x_2 p_3 - x_3 p_2, x_3 p_1 - x_1 p_3\} = \{x_2 p_3, x_3 p_1\} - \{x_2 p_3, x_1 p_3\} - \{x_3 p_2, x_3 p_1\} + \{x_3 p_2, x_1 p_3\}$$

~~$$\{AB, CD\} = A\{B, CD\} + \{A, CD\}B = -A\{CD, B\} - \{CD, A\}B$$

$$= -AC\{D, B\} - A\{C, B\}D - C\{D, A\}B - \{D, A\}CB$$

$$\Rightarrow \{x_2 p_3, x_3 p_1\} = -x_2 x_3 \{p_3, p_1\} - p_1 x_2$$~~

$$\{AB, CD\} = A\{B, CD\} + \{A, CD\}B = -A\{CD, B\} - \{CD, A\}B$$

$$= -AC\{D, B\} - A\{C, B\}D - C\{D, A\}B - \{C, A\}DB$$

$$= AC\{B, D\} + A\{B, C\}D + C\{A, D\}B + \{A, C\}DB$$

(if we have to worry about non-commutativity... of course, we don't have to here).

→ We will only get contributions if we have x_i and p_i in the same bracket, so

$$\{x_2 p_3, x_3 p_1\} = \cancel{x_2 x_3} \{p_3, p_1\} + x_2 \{p_3, x_3\} p_1 + x_3 \{x_2 p_3\} p_1 + \cancel{x_2 x_3} \{p_1, p_3\}$$

$$= -x_2 p_1$$

$$\{x_2 p_3, x_1 p_3\} = 0$$

$$\{x_3 p_2, x_3 p_1\} = 0$$

$$\{x_3 p_2, x_1 p_3\} = x_1 x_3 \{p_2, p_3\} + x_3 \{p_2, x_1\} p_3 + x_1 \{x_3 p_2\} p_3 + \cancel{x_3 x_1} \{p_3, p_2\}$$

$$= x_1 p_2$$

$$\Rightarrow \{L_1, L_2\} = x_1 p_2 - x_2 p_1 = L_3 \quad \text{etc.}$$

⇒ Lenz - Runge as ungraded HW or practice problem.

Canonical Transformations

- Lagrange equations are invariant under general coordinate transformations in configuration space
 - Hamiltonian formalism treats q_i, p_i as independent, so let's look for form invariance in phase space
- ↳ will make it possible to find canonical variables that simplify H and make it easy to solve Hamilton's eqs.:

$$q_i \rightarrow Q_i(q_k, p_k) \quad p_i \rightarrow P_i(q_k, p_k)$$

We can write Hamilton's eqs. in symplectic form (symplectic, g^0 , "inter-twined").

$$\vec{x} = (q_1, \dots, q_n, p_1, \dots, p_n)^T \quad ; \quad \frac{\partial H}{\partial \vec{x}} = \left(\frac{\partial H}{\partial q_1}, \dots, \frac{\partial H}{\partial q_n}, \frac{\partial H}{\partial p_1}, \dots, \frac{\partial H}{\partial p_n} \right)^T$$
$$J = \begin{pmatrix} q_i & p_i \\ -\mathbb{1} & \mathbb{1} \end{pmatrix}$$

$$\Rightarrow \boxed{\dot{\vec{x}} = J \frac{\partial H}{\partial \vec{x}}}$$

Work in symplectic form and make the transformation

$$x_j \rightarrow y_j(x_k)$$

$$\dot{y}_j = \sum_{k=1}^{2n} \frac{\partial y_j}{\partial x_k} \dot{x}_k = \sum_{k\ell} \frac{\partial y_j}{\partial x_k} J_{k\ell} \frac{\partial H}{\partial x_\ell} = \sum_{k\ell m} \frac{\partial y_j}{\partial x_k} J_{k\ell} \frac{\partial H}{\partial y_m} \frac{\partial y_m}{\partial x_\ell}$$
$$= \sum_{k\ell m} \frac{\partial y_j}{\partial x_k} J_{k\ell} \frac{\partial y_m}{\partial x_\ell} \frac{\partial H}{\partial y_m}$$

Define the Jacobian of the coordinate change

$$M_{jk} = \frac{\partial y_j}{\partial x_k}$$

$$\Rightarrow \dot{y}_j = \sum_{k\ell m} M_{jk} J_{k\ell} M_{\ell m}^T \frac{\partial H}{\partial y_m} \Rightarrow \dot{\vec{y}} = (M J M^T) \frac{\partial H}{\partial \vec{y}}$$

Thus, Hamilton's eqs. are form invariant if

$$M J M^T = J \Leftrightarrow \sum_{k\ell m} \frac{\partial y_j}{\partial x_k} J_{k\ell} \frac{\partial y_m}{\partial x_\ell} = J_{jk}$$

If this relationship holds, the Jacobian is called symplectic and the transformation $x_i \rightarrow y_j(x_k)$ is canonical.

The Poisson brackets are invariant under canonical transformations. Conversely, a trafo which preserves the ^{fundamental} Poisson brackets ~~structure~~

$$\{Q_i, Q_j\} = \{P_i, P_j\} = 0, \quad \{Q_i, P_j\} = \delta_{ij}$$

is canonical.

Proof

$$\{f, g\} = \sum_{k=1}^n \left(\frac{\partial f}{\partial q_k} \frac{\partial g}{\partial p_k} - \frac{\partial g}{\partial p_k} \frac{\partial f}{\partial q_k} \right) = \sum_{rs} \frac{\partial f}{\partial x_r} J_{rs} \frac{\partial g}{\partial x_s}$$

\uparrow
 $1 \dots 2n$

Now transform: $\frac{\partial f}{\partial x_r} = \sum_s \frac{\partial f}{\partial y_s} \frac{\partial y_s}{\partial x_r} = \sum_s \frac{\partial f}{\partial y_s} M_{sr}$

$$\begin{aligned} \hookrightarrow \{f, g\}_x &= \sum_{rstu} \frac{\partial f}{\partial y_t} M_{tr} J_{rs} M_{su} \frac{\partial g}{\partial y_u} = \sum_{tu} \frac{\partial f}{\partial y_t} (M J M^T)_{tu} \frac{\partial g}{\partial y_u} \\ &= \sum_{tu} \frac{\partial f}{\partial y_t} J_{tu} \frac{\partial g}{\partial y_u} = \{f, g\}_y \end{aligned}$$

We now consider the other direction: Use the "old" notation

$$\begin{aligned} \{f, g\}_{(q,p)} &= \sum_i \left(\frac{\partial f}{\partial q_i} \frac{\partial g}{\partial p_i} - \frac{\partial g}{\partial q_i} \frac{\partial f}{\partial p_i} \right) \\ &= \sum_{jke} \left[\left(\frac{\partial f}{\partial q_k} \frac{\partial Q_k}{\partial q_j} + \frac{\partial f}{\partial p_k} \frac{\partial P_k}{\partial q_j} \right) \left(\frac{\partial g}{\partial Q_e} \frac{\partial Q_e}{\partial p_i} + \frac{\partial g}{\partial P_e} \frac{\partial P_e}{\partial p_i} \right) - \left(\frac{\partial g}{\partial Q_e} \frac{\partial Q_e}{\partial q_j} + \frac{\partial g}{\partial P_e} \frac{\partial P_e}{\partial q_j} \right) \left(\frac{\partial f}{\partial Q_k} \frac{\partial Q_k}{\partial p_i} + \frac{\partial f}{\partial P_k} \frac{\partial P_k}{\partial p_i} \right) \right] \\ &= \sum_{jke} \left[\frac{\partial f}{\partial Q_k} \frac{\partial g}{\partial P_e} \left(\frac{\partial Q_k}{\partial q_j} \frac{\partial P_e}{\partial p_i} - \frac{\partial P_e}{\partial q_j} \frac{\partial Q_k}{\partial p_i} \right) + \frac{\partial f}{\partial p_k} \frac{\partial g}{\partial Q_e} \left(\frac{\partial P_k}{\partial q_j} \frac{\partial Q_e}{\partial p_i} - \frac{\partial Q_e}{\partial q_j} \frac{\partial P_k}{\partial p_i} \right) \right. \\ &\quad \left. + \frac{\partial f}{\partial Q_k} \frac{\partial g}{\partial Q_e} \left(\frac{\partial Q_k}{\partial q_j} \frac{\partial Q_e}{\partial p_i} - \frac{\partial Q_e}{\partial q_j} \frac{\partial Q_k}{\partial p_i} \right) + \frac{\partial f}{\partial p_k} \frac{\partial g}{\partial P_e} \left(\frac{\partial P_k}{\partial q_j} \frac{\partial P_e}{\partial p_i} - \frac{\partial P_e}{\partial q_j} \frac{\partial P_k}{\partial p_i} \right) \right] \\ &= \sum_{ke} \left[\frac{\partial f}{\partial Q_k} \frac{\partial g}{\partial P_e} \{Q_k, P_e\}_{(q,p)} + \frac{\partial f}{\partial p_k} \frac{\partial g}{\partial Q_e} \{P_k, Q_e\}_{(q,p)} \right. \\ &\quad \left. + \frac{\partial f}{\partial Q_k} \frac{\partial g}{\partial Q_e} \{Q_k, Q_e\}_{(q,p)} + \frac{\partial f}{\partial p_k} \frac{\partial g}{\partial P_e} \{P_k, P_e\}_{(q,p)} \right] \\ &\stackrel{!}{=} \{f, g\}_{(q,p)} = \sum_k \left(\frac{\partial f}{\partial Q_k} \frac{\partial g}{\partial P_k} - \frac{\partial g}{\partial Q_k} \frac{\partial f}{\partial P_k} \right) \end{aligned}$$

$$\Rightarrow \{Q_k, P_e\} = \delta_{ke}, \quad \{Q_k, Q_e\} = 0, \quad \{P_k, P_e\} = 0$$

Note that this also implies - consider 2x2 blocks:

$$(M J M^T)_{jk} = \begin{pmatrix} \{Q_j, Q_k\} & \{Q_j, P_k\} \\ \{P_j, Q_k\} & \{P_j, P_k\} \end{pmatrix} = \begin{pmatrix} 0 & \delta_{jk} \\ -\delta_{jk} & 0 \end{pmatrix}$$

Infinitesimal canonical transformations

$$q_j \rightarrow Q_j = q_j + \alpha F_j(q_k, p_k)$$

$$p_j \rightarrow P_j = p_j + \alpha E_j(q_k, p_k)$$

What functions are allowed? The (2x2) blocks of the Jacobian are

$$M_{jl} = \begin{pmatrix} \delta_{jk} + \alpha \frac{\partial F_j}{\partial q_k} & \alpha \frac{\partial F_j}{\partial p_k} \\ \alpha \frac{\partial E_j}{\partial q_k} & \delta_{jk} + \alpha \frac{\partial E_j}{\partial p_k} \end{pmatrix}$$

So $M J M^T = J$ implies

$$\begin{aligned} & \sum_k \begin{pmatrix} \delta_{jk} + \alpha \frac{\partial F_j}{\partial q_k} & \alpha \frac{\partial F_j}{\partial p_k} \\ \alpha \frac{\partial E_j}{\partial q_k} & \delta_{jk} + \alpha \frac{\partial E_j}{\partial p_k} \end{pmatrix} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \begin{pmatrix} \delta_{ke} + \alpha \frac{\partial F_k}{\partial q_e} & \alpha \frac{\partial E_k}{\partial q_e} \\ \alpha \frac{\partial F_k}{\partial p_e} & \delta_{ke} + \alpha \frac{\partial E_k}{\partial p_e} \end{pmatrix} \\ &= \sum_k \begin{pmatrix} \alpha \frac{\partial F_k}{\partial p_e} & \delta_{ke} + \alpha \frac{\partial E_k}{\partial p_e} \\ -(\delta_{ke} + \alpha \frac{\partial F_k}{\partial q_e}) & -\alpha \frac{\partial E_k}{\partial q_e} \end{pmatrix} \\ &= \sum_k \begin{pmatrix} (\delta_{jk} + \alpha \frac{\partial F_j}{\partial q_k}) \alpha \frac{\partial F_k}{\partial p_e} - \alpha \frac{\partial F_j}{\partial p_k} (\delta_{ke} + \alpha \frac{\partial F_k}{\partial q_e}) & (\delta_{jk} + \alpha \frac{\partial F_j}{\partial q_k}) (\delta_{ke} + \alpha \frac{\partial E_k}{\partial p_e}) - \alpha^2 \frac{\partial F_j}{\partial p_k} \frac{\partial E_k}{\partial q_e} \\ \alpha^2 \frac{\partial E_j}{\partial q_k} \frac{\partial F_k}{\partial p_e} - (\delta_{jk} + \alpha \frac{\partial E_j}{\partial p_k}) (\delta_{ke} + \alpha \frac{\partial F_k}{\partial q_e}) & \alpha \frac{\partial E_j}{\partial p_k} (\delta_{ke} + \alpha \frac{\partial E_k}{\partial p_e}) - (\delta_{jk} + \alpha \frac{\partial E_j}{\partial p_k}) \alpha \frac{\partial E_k}{\partial q_e} \end{pmatrix} \\ &= \sum_k \begin{pmatrix} \alpha (\delta_{jk} \frac{\partial F_k}{\partial p_e} - \frac{\partial F_j}{\partial p_k} \delta_{ke}) & \delta_{jk} \delta_{ke} + \alpha (\delta_{jk} \frac{\partial E_k}{\partial p_e} + \frac{\partial F_j}{\partial q_k} \delta_{ke}) \\ -\delta_{jk} \delta_{ke} - \alpha (\frac{\partial E_j}{\partial p_k} \delta_{ke} + \frac{\partial F_k}{\partial q_e} \delta_{jk}) & \alpha (\frac{\partial E_j}{\partial p_k} \delta_{ke} - \delta_{jk} \frac{\partial E_k}{\partial q_e}) \end{pmatrix} \\ &= \begin{pmatrix} \alpha (\frac{\partial F_j}{\partial p_e} - \frac{\partial F_j}{\partial p_e}) & \delta_{je} + \alpha (\frac{\partial E_j}{\partial p_e} + \frac{\partial F_j}{\partial q_e}) \\ -\delta_{je} - \alpha (\frac{\partial E_j}{\partial p_e} + \frac{\partial F_j}{\partial q_e}) & \alpha (\frac{\partial E_j}{\partial p_e} - \frac{\partial E_j}{\partial q_e}) \end{pmatrix} \\ &= \begin{pmatrix} 0 & \delta_{je} \\ -\delta_{je} & 0 \end{pmatrix} + \alpha \begin{pmatrix} 0 & (\frac{\partial E_j}{\partial p_e} + \frac{\partial F_j}{\partial q_e}) \\ -(\frac{\partial E_j}{\partial p_e} + \frac{\partial F_j}{\partial q_e}) & 0 \end{pmatrix} \doteq J \end{aligned}$$

This means

$$\frac{\partial F_1}{\partial p_i} = - \frac{\partial E_1}{\partial q_i}$$

which is true if

$$F_1 = \frac{\partial G}{\partial p_i}, \quad E_1 = - \frac{\partial G}{\partial q_i}$$

for some generator $G(q_k, p_k)$

Suppose we have a one-parameter family of transformations

$$q_j \rightarrow Q_j(q_k, p_k; \alpha) \quad p_j \rightarrow P_j(q_k, p_k; \alpha)$$

that is an active transformation in phase space, from

$$(q_j, p_j) \rightarrow (Q_j(q_k, p_k; \alpha), P_j(q_k, p_k; \alpha))$$

As we vary α , we map out a line in phase space - a flow generated by G . This is evident from the tangents:

$$\frac{dq_j}{d\alpha} = \frac{\partial G}{\partial p_j}, \quad \frac{dp_j}{d\alpha} = - \frac{\partial G}{\partial q_j}$$

\Rightarrow Looks just like ~~the~~ Hamilton's eqns. with $t \rightarrow \alpha$, $H \rightarrow G$.

Thus, time evolution can be thought of as a canonical transformation generated by the Hamiltonian (analogous to unitary time evolution in QM.)