

Liouville's Theorem

We have rewritten mechanics in terms of a flow in phase space, a time evolution generated by H .

Consider an infinitesimal volume

$$dV = dq_1 \dots dq_n dp_1 \dots dp_n$$

$$\Rightarrow \text{for } q_i \rightarrow q_i + \dot{q}_i dt = q_i + \frac{\partial H}{\partial p_i} dt = q_i'$$

$$p_i \rightarrow p_i + \dot{p}_i dt = p_i - \frac{\partial H}{\partial q_i} dt = p_i'$$

$$\Rightarrow dV' = dq_1' \dots dq_n' dp_1' \dots dp_n' = \det M \cdot dV$$

~~$$\begin{aligned} &= \det \begin{pmatrix} \frac{\partial q_1'}{\partial q_1} & \dots & \frac{\partial q_1'}{\partial p_1} \\ \vdots & \ddots & \vdots \\ \frac{\partial q_n'}{\partial q_1} & \dots & \frac{\partial q_n'}{\partial p_1} \\ \vdots & \ddots & \vdots \\ \frac{\partial p_1'}{\partial q_1} & \dots & \frac{\partial p_1'}{\partial p_1} \\ \vdots & \ddots & \vdots \\ \frac{\partial p_n'}{\partial q_1} & \dots & \frac{\partial p_n'}{\partial p_1} \end{pmatrix} dV \\ &= \det \begin{pmatrix} 1 + \frac{\partial^2 H}{\partial p_i \partial q_j} dt & & \\ & \ddots & \\ \frac{\partial^2 H}{\partial q_i \partial q_j} dt & & 1 - \frac{\partial^2 H}{\partial p_i \partial p_j} dt \end{pmatrix} dV \\ &= \det M \cdot dV \end{aligned}$$~~

Recall $M J M^T = J \Rightarrow \det(M J M^T) = \det J$
 $\det M \cdot \det J \cdot \det M^T = \det J$
 $(\det M)^2 \cdot \det J = \det J$
 $\Rightarrow (\det M)^2 = 1$

This leaves the possibilities $\det M = \pm 1$, but we must have $M \rightarrow \mathbb{I}$ for infinitesimal canonical maps, so $\det M = 1$. This means canonical transformations preserve the volume element (or general volumes) in phase space, and ~~the~~ time evolution via H does so in particular! This is Liouville's Theorem.
Can use it to evolve distributions $\rho(q, p, t)$:

• e.g. single system, but we don't know precise values of q, p :

$$\int \rho(q_i, p_i, t) dp_1 \dots dp_n dq_1 \dots dq_n = 1$$

• large number of particles (Stat-Mech.!) \Rightarrow

$$\int \rho(q_i, p_i, t) dp_1 \dots dp_n dq_1 \dots dq_n = N$$

\Rightarrow particles / probabilities must be conserved, so

~~$$\frac{d\rho}{dt} = 0 = \frac{\partial \rho}{\partial t} + \{ \rho, H \}$$~~

$$\frac{d\rho}{dt} = \{ \rho, H \} + \frac{\partial \rho}{\partial t} = 0$$

Liouville's equation