U.S. Particle Accelerator School
Education in Beam Physics and Accelerator Technology

Self-Consistent Simulations of Beam and Plasma Systems
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A3. Special Topics

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Outline

- Particle pushers
  - Relativistic Boris pusher
  - Lorentz invariant pusher
  - Application to the modeling of electron cloud instability
- Quasistatic method
  - Concept
  - Application to the modeling of electron cloud instability
- Optimal Lorentz boosted frame
  - Concept
  - Application to the modeling of electron cloud instability
  - Generalization
  - Application to the modeling of laser-plasma accelerators

Relativistic Boris pusher

For the velocity component, the Boris pusher writes

$$u^{n+1} = u^n + \frac{q\Delta t}{m}\left(E^{n+1/2} + \frac{u^{n+1}+u^n}{2\gamma^{n+1/2}} \times B^{n+1/2}\right) \quad \text{with} \quad u = \gamma v$$

which decomposes into

one acceleration **+** ***one rotation*** **+** ***one acceleration***

$$u^- = u^n + \frac{q\Delta t}{2m}E^{n+1/2} \Rightarrow u^+ - u^- = \frac{q\Delta t}{2m\gamma^{n+1/2}}\left(u^+ + u^-\right)\times B^{n+1/2} \Rightarrow u^{n+1} = u^+ + \frac{q\Delta t}{2m}E^{n+1/2}$$

with

$$\gamma^{n+1/2} = \sqrt{1+\left(u^n + \frac{q\Delta t}{2m}E^{n+1/2}\right)^2 / c^2} = \sqrt{1+\left(u^{n+1} - \frac{q\Delta t}{2m}E^{n+1/2}\right)^2 / c^2}$$

Relativistic Boris pusher: problem with E+v×B≈0

Assuming E and B such that E+v×B=0:

$$\Rightarrow u^{n+1} = u^n \quad \Rightarrow \gamma^{n+1/2} = \gamma^n = \gamma^{n+1}$$

$$\Rightarrow \gamma^{n+1/2} = \sqrt{1+\left(u^n + \frac{q\Delta t}{2m}E^{n+1/2}\right)^2 / c^2} = \sqrt{1+\left(u^n - \frac{q\Delta t}{2m}E^{n+1/2}\right)^2 / c^2}$$

$$\Rightarrow E^{n+1/2} = -E^{n+1/2} = 0 \quad \Rightarrow B^{n+1/2} = 0$$

meaning that pusher is consistent with (E+v×B=0) only if E=B=0, and is thus inaccurate for e.g. ultra-relativistic beams.

Lorentz invariant particle pusher

Replace Boris velocity pusher

– Velocity push: $u^{n+1} = u^n + \frac{q\Delta t}{m}\left(E^{n+1/2} + \boxed{\frac{u^{n+1}+u^n}{2\gamma^{n+1/2}}} \times B^{n+1/2}\right)$ $u = \gamma v$

with

– Velocity push: $u^{n+1} = u^n + \frac{q\Delta t}{m}\left(E^{n+1/2} + \boxed{\frac{v^{n+1}+v^n}{2}} \times B^{n+1/2}\right)$

Looks implicit but solvable analytically

$$\begin{cases} \gamma^{i+1} = \sqrt{\dfrac{\sigma + \sqrt{\sigma^2 + 4(\tau^2 + u^{*2})}}{2}} \\ \mathbf{u}^{i+1} = [\mathbf{u}' + (\mathbf{u}' \cdot \mathbf{t})\mathbf{t} + \mathbf{u}' \times \mathbf{t}]/(1+t^2) \end{cases}$$

with

$$\begin{cases} \mathbf{u}' = \mathbf{u}^i + \frac{q\Delta t}{m}\left(\mathbf{E}^{i+1/2} + \frac{\mathbf{v}^i}{2} \times \mathbf{B}^{i+1/2}\right) \\ \boldsymbol{\tau} = (q\Delta t/2m)\mathbf{B}^{i+1/2} \\ u^* = \mathbf{u}' \cdot \boldsymbol{\tau}/c \\ \sigma = \gamma'^2 - \tau^2 \\ \gamma' = \sqrt{1 + u'^2/c^2} \\ \mathbf{t} = \boldsymbol{\tau}/\gamma^{i+1} \end{cases}$$

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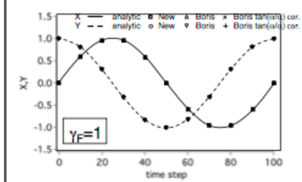
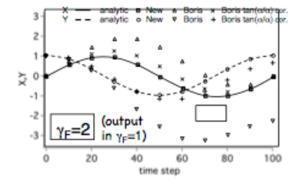
Lorentz invariant particle pusher: test w/ 1 particle

Lab frame — particle cycling in constant B field ⇨ **Boosted frame γ=2** — ExB drift adds to gyration





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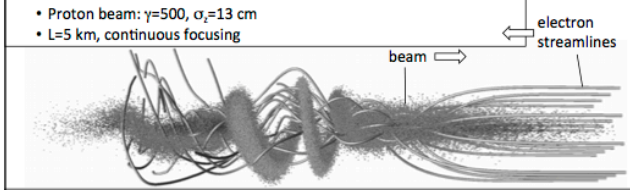
Application to modeling of two-stream instability

Calculation of e-cloud induced instability of a proton bunch

- Proton beam: γ=500, σ_z=13 cm
- L=5 km, continuous focusing



Beam was lost after a few betatron oscillations with Boris pusher.

Accurate result was obtained with new pusher.

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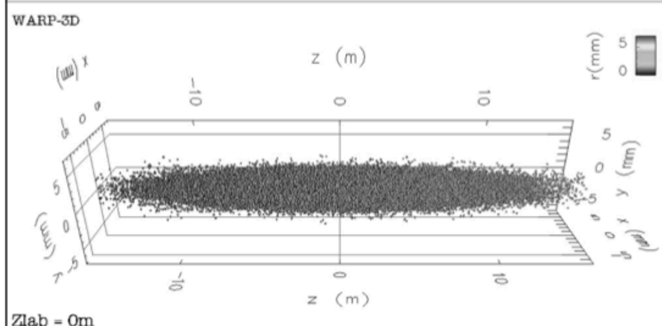
Application to modeling of two-stream instability



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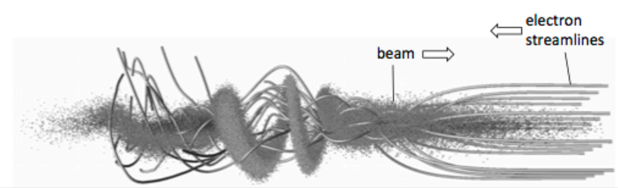
Modeling of two-stream instability is expensive

Need to follow short (σ_z=13 cm) and stiff (γ=500) proton beam for 5 km:

- mobile background electrons react in fraction of beam ➔ small time steps



Two solutions:

- separate treatment of slow (beam) and fast (electrons) components ➔ quasistatic approx.
- solve in a Lorentz boosted frame which matches beam & electrons time scales

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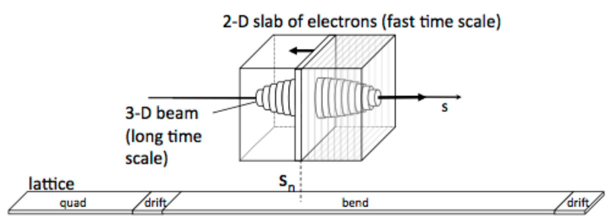
Quasistatic approximation



repeat:

1. 2-D slab of electrons is stepped backward (with small time steps) through the beam field and its self-field (solving 2-D Poisson at each step),
2. 2-D electron fields are stacked in a 3-D array and added to beam self-field,
3. 3-D field is used to kick the 3-D beam,
4. 3-D beam is pushed to next station with large time steps,
5. Solve Poisson for 3-D beam self-field.

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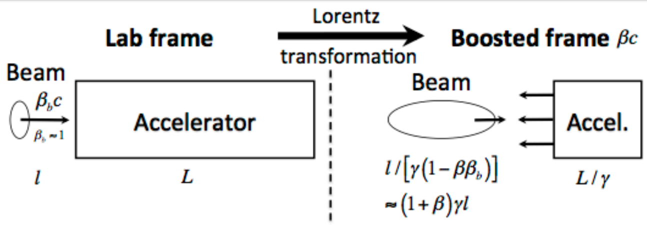
Optimal Lorentz boosted frame



Many time steps needed to follow short stiff high-energy beam into long accelerator filled with fast reacting electron clouds.

Much less time steps needed to follow long low-energy beam into shorter accelerator filled with stiffer electron clouds.

Number of time steps divided by $(1+\beta)\gamma^2$

With high γ, orders of magnitude speedups are possible.

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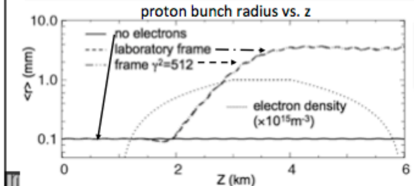

CPU time (on 8 cores in 2006):
- **lab frame: >2 weeks**
- **frame with γ^2=512: <30 min**

Speedup x1000

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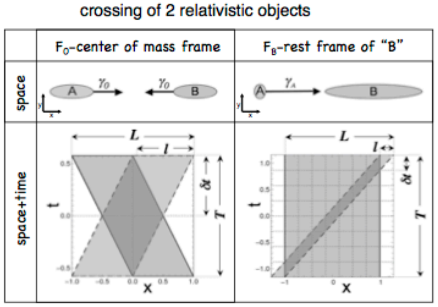
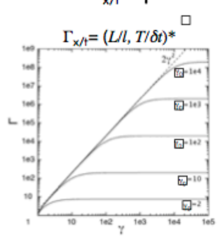
Generalization of optimal boosted frame approach

General formulation:
crossing of 2 relativistic objects



Range of space/time scales $\Gamma_{x/t} \propto \gamma^2$

$\Gamma_{x/t}$= $(L/l,\ T/\delta t)$*



The range of space and time scales is not a Lorentz invariant and scales as γ^2 for the crossing of two relativistic objects (matter of photons).

Applicable to study of electron cloud effects, plasma accelerators, free electron lasers, etc.

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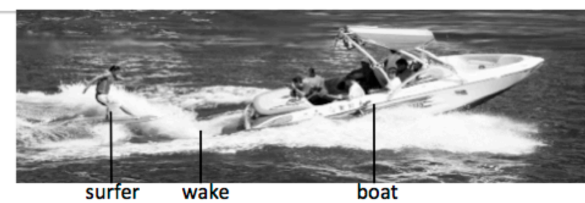
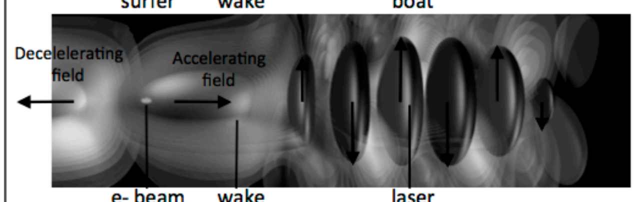
Laser plasma accelerators "surf" electrons on plasma waves
for acceleration on ultra short distances





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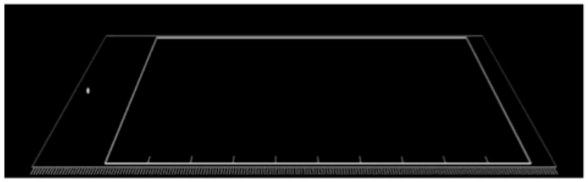
Modeling from first principle is very challenging



For a 10 GeV scale stage:

~**1µm** wavelength laser propagates into ~**1m** plasma

➔ millions of time steps needed

(similar to modeling **5m** boat crossing ~**5000** km Atlantic Ocean)

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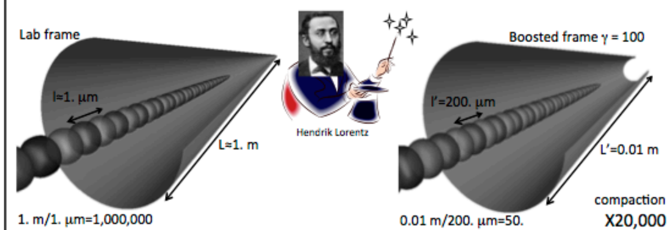
Optimal boosted frame enables large speedup



Alternate or complementary solutions: quasistatic, laser envelope, azimuthal Fourier decomposition ("Circ"), …

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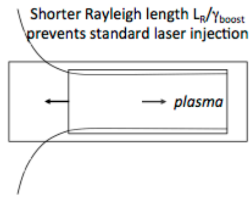
Laser injection through moving plane solves initialization issue in LBF

Lab frame

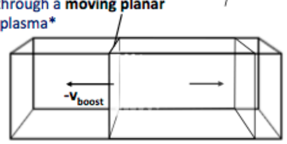
Standard laser injection from left boundary or all at once

Boosted frame

Shorter Rayleigh length L_R/γ_{boost} prevents standard laser injection



Solution: injection through a **moving planar antenna** in front of plasma*



- Laser injected using macroparticles using Esirkepov current deposition ==> verifies Gauss' Law.
- For high γ_{boost}, backward radiation is blue shifted and unresolved.

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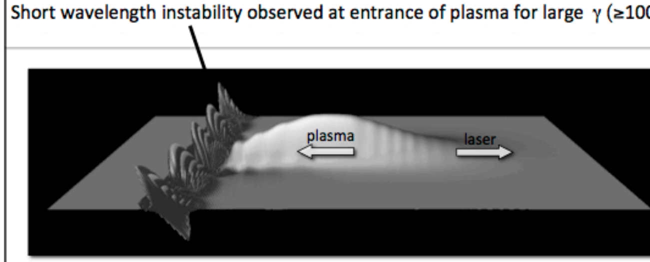
Short wavelength instability observed at entrance of plasma for large γ (≥100)



Is it numerical Cherenkov instability?

BTW, what is "numerical Cherenkov instability"?

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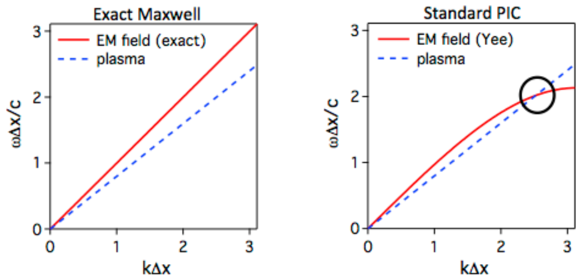
Relativistic plasmas PIC subject to "numerical Cherenkov"

B. B. Godfrey, "Numerical Cherenkov instabilities in electromagnetic particle codes", *J. Comput. Phys.* **15** (1974)

Numerical dispersion leads to crossing of EM field and plasma modes -> instability.



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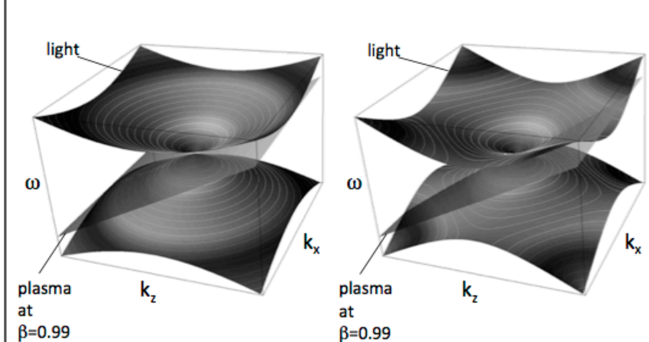
Space/time discretization aliases ➔ more crossings in 2/3-D

Exact Maxwell **Standard PIC**



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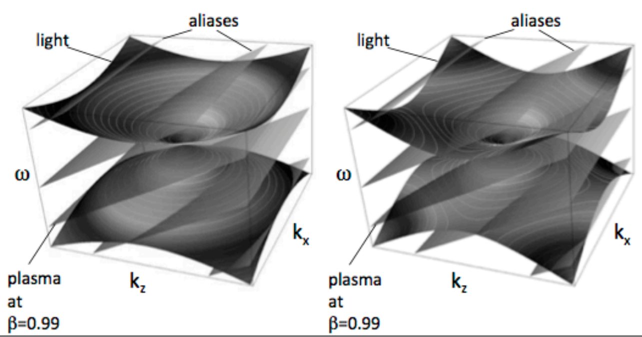
Space/time discretization aliases ➔ more crossings in 2/3-D

Exact Maxwell **Standard PIC**



Analysis calls for full PIC numerical dispersion relation

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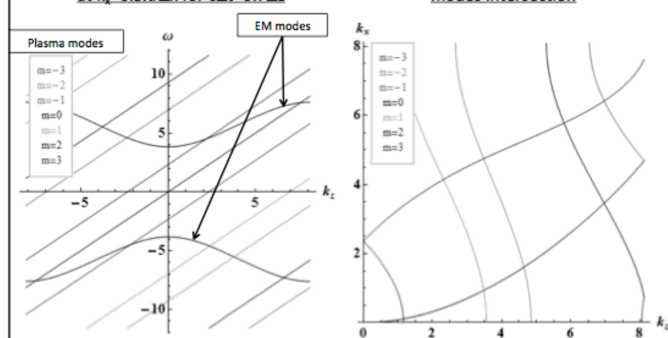
Maps of unstable modes

Normal modes at k_x=0.5π/Δx for cΔt=0.7Δz

Projection of normal modes intersection



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Numerical dispersion relation of full-PIC algorithm

2-D relation (Fourier space):

$$\begin{pmatrix} \xi_{z,z}+[\omega] & \xi_{z,x} & \xi_{z,y}+[k_x] \\ \xi_{x,z} & \xi_{x,x}+[\omega] & \xi_{x,y}-[k_z] \\ [k_x] & -[k_z] & [\omega] \end{pmatrix} \begin{pmatrix} E_z \\ E_x \\ B_y \end{pmatrix} = 0.$$

$$[\omega] = \sin\left(\omega\frac{\Delta t}{2}\right) / \left(\frac{\Delta t}{2}\right) \quad [k_z] = k_z \sin\left(k\frac{\Delta t}{2}\right) / \left(k\frac{\Delta t}{2}\right) \quad [k_x] = k_x \sin\left(k\frac{\Delta t}{2}\right) / \left(k\frac{\Delta t}{2}\right)$$

$$S^J = \left[\sin\left(k_z'\frac{\Delta z}{2}\right) / \left(k_z'\frac{\Delta z}{2}\right)\right]^{\ell_z+1} \left[\sin\left(k_x'\frac{\Delta x}{2}\right) / \left(k_x'\frac{\Delta x}{2}\right)\right]^{\ell_x+1},$$

$$S^{E_z} = \left[\sin\left(k_z'\frac{\Delta z}{2}\right) / \left(k_z'\frac{\Delta z}{2}\right)\right]^{\ell_z} \left[\sin\left(k_x'\frac{\Delta x}{2}\right) / \left(k_x'\frac{\Delta x}{2}\right)\right]^{\ell_x+1} (-1)^{m_z},$$

$$S^{E_x} = \left[\sin\left(k_z'\frac{\Delta z}{2}\right) / \left(k_z'\frac{\Delta z}{2}\right)\right]^{\ell_z+1} \left[\sin\left(k_x'\frac{\Delta x}{2}\right) / \left(k_x'\frac{\Delta x}{2}\right)\right]^{\ell_x} (-1)^{m_x},$$

$$S^{B_y} = \cos\left(\omega\frac{\Delta t}{2}\right) \left[\sin\left(k_z'\frac{\Delta z}{2}\right) / \left(k_z'\frac{\Delta z}{2}\right)\right]^{\ell_z} \left[\sin\left(k_x'\frac{\Delta x}{2}\right) / \left(k_x'\frac{\Delta x}{2}\right)\right]^{\ell_x} (-1)^{m_z+m_x}.$$

*B. B. Godfrey, J. L. Vay, I. Haber, J. Comp. Phys. 248 (2013)

Numerical dispersion relation of full-PIC algorithm (II)

$$\xi_{z,z} = -n\gamma^{-2}\sum_m S^J S^{E_z} \csc^2\left[(\omega - k_z'v)\frac{\Delta t}{2}\right] \left(kk_z^2\Delta t + \zeta_z k_x^2 \sin(k\Delta t)\right)\Delta t\,[\omega]\,k_z'/4k^3k_z,$$

$$\xi_{z,x} = -n\sum_m S^J S^{E_x} \csc\left[(\omega - k_z'v)\frac{\Delta t}{2}\right] \eta_z k_x'/2k^3k_z,$$

$$\xi_{z,y} = nv\sum_m S^J S^{B_y} \csc\left[(\omega - k_z'v)\frac{\Delta t}{2}\right] \eta_z k_x'/2k^3k_z,$$

$$\xi_{x,z} = -n\gamma^{-2}\sum_m S^J S^{E_z} \csc^2\left[(\omega - k_z'v)\frac{\Delta t}{2}\right] \left(k\Delta t - \zeta_z \sin(k\Delta t)\right)\Delta t\,[\omega]\,k_xk_z'/4k^3,$$

$$\xi_{x,x} = -n\sum_m S^J S^{E_x} \csc\left[(\omega - k_z'v)\frac{\Delta t}{2}\right] \eta_x k_x'/2k^3k_x,$$

$$\xi_{x,y} = nv\sum_m S^J S^{B_y} \csc\left[(\omega - k_z'v)\frac{\Delta t}{2}\right] \eta_x k_x'/2k^3k_x,$$

*B. B. Godfrey, J. L. Vay, I. Haber, J. Comp. Phys. 248 (2013)

Numerical dispersion relation of full-PIC algorithm (III)

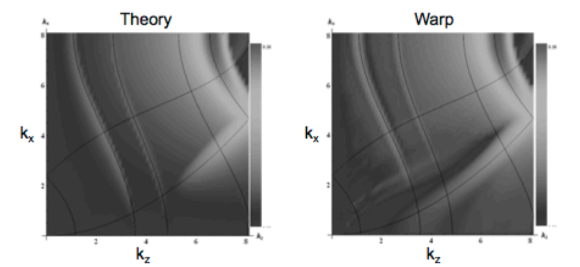
$$\eta_z = \cot\left[(\omega - k_z'v)\frac{\Delta t}{2}\right]\left(kk_z^2\Delta t + \zeta_z k_x^2\sin(k\Delta t)\right)\sin\left(k_z'v\frac{\Delta t}{2}\right) + \left(k\Delta t - \zeta_x\sin(k\Delta t)\right)k_z^2\cos\left(k_z'v\frac{\Delta t}{2}\right),$$

$$\eta_x = \cot\left[(\omega - k_z'v)\frac{\Delta t}{2}\right]\left(k\Delta t - \zeta_z\sin(k\Delta t)\right)k_x^2\sin\left(k_z'v\frac{\Delta t}{2}\right) + \left(kk_x^2\Delta t + \zeta_x k_z^2\sin(k\Delta t)\right)\cos\left(k_z'v\frac{\Delta t}{2}\right).$$

Then simplify and solve with Mathematica…

*B. B. Godfrey, J. L. Vay, I. Haber, J. Comp. Phys. 248 (2013)

Growth rates from theory match Warp simulations



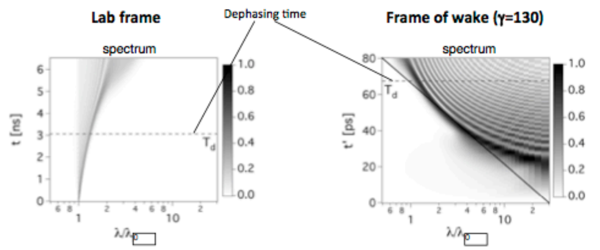
Warp run uses uniform drifting plasma with periodic BC.
Yee finite difference, energy conserving gather (cΔt/Δx=0.7)

Latest theory has led to ne insight and the development of very effective methods to mitigate the instability.

Physics in boosted frame also allows the use of wideband filtering

Time history of laser spectrum (relative to laser λ_0 in vacuum)

Spectrum very different in lab and boosted frames



Content concentrated around λ_0

Content concentrated at much larger λ.

More filtering possible without altering physics*.

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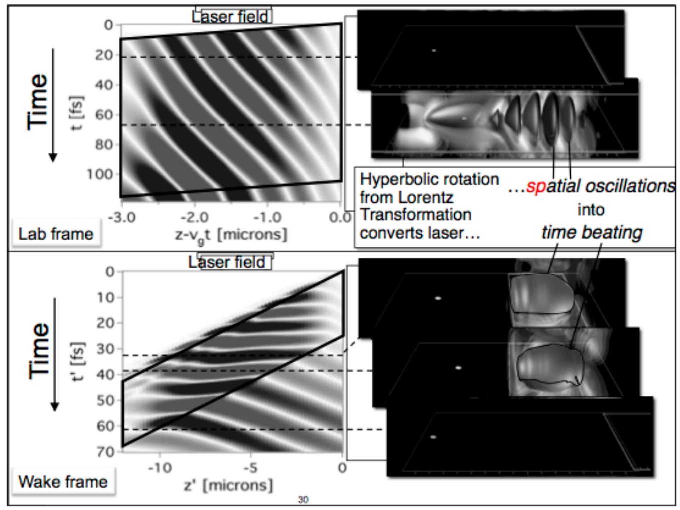

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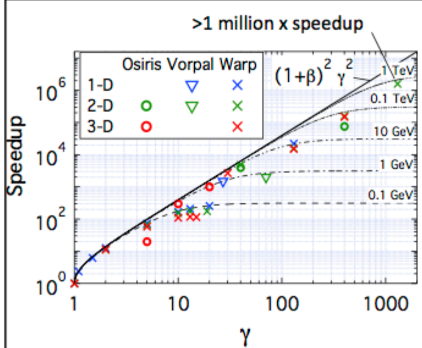
Speedup verified by us and others to over a million



Warp:

1. J.-L. Vay, et al., *Phys. Plasmas* **18** 123103 (2011)
2. J.-L. Vay, et al., *Phys. Plasmas (letter)* **18** 030701 (2011)
3. J.-L. Vay, et al., *J. Comput. Phys.* **230** 5908 (2011)
4. J.-L. Vay et al, PAC Proc. (2009)

Osiris:

1. S. Martins, et al., *Nat. Phys.* **6** 311 (2010)
2. S. Martins, et al., *Comput. Phys. Comm.* **181** 869 (2010)
3. S. Martins, et al., *Phys. Plasmas* **17** 056705 (2010)
4. S. Martins et al, PAC Proc. (2009)

Vorpal:

1. D. Bruhwiler, et al., *AIP Conf. Proc* **1086** 29 (2009)

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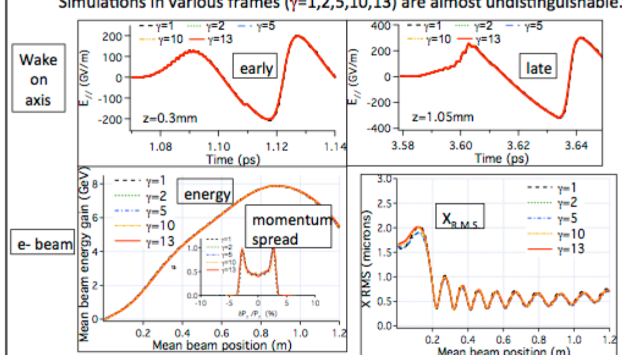
Very high precision validation of BF method with Warp

Simulations in various frames (γ=1,2,5,10,13) are almost undistinguishable.



Warp-3D – a_0=1, $n_0=10^{19}cm^{-3}$ (~100 MeV) scaled to $10^{17}cm^{-3}$ (~10 GeV).
Detailed validation for a0>1 (non-linear regime) is underway.

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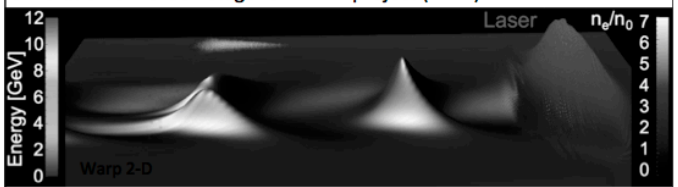
Enabling simulations that were previously untractable

Simulation of 10 GeV stage for BELLA project (LBNL)



State-of-the-art PIC simulations of 10 GeV stages:

2006 (lab) in 1D: ~ 5k CPU-hours ➔ 2011 (boost) in 3D: ~ 1k CPU-hours

Current state-of-the-art in lab: 2-D RZ simulations in ~2 weeks on thousands of cores.

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Special topics summary

- Modeling of relativistic beams/plasmas with full PIC may benefit from "non-standard" algorithms
  - Lorentz invariant particle pusher
  - Quasistatic approximation
  - Optimal Lorentz boosted frame
- Quasistatic is well established method, but requires writing dedicated code or module
- Boosted frame approach is newer and uses standard PIC at core, needing only extensions

References

1. Brendan B. Godfrey, Jean-Luc Vay, "Improved numerical Cherenkov instability suppression in the generalized PSTD PIC algorithm", Computer Physics Communications, 196, 221 (2015) http://dx.doi.org/10.1016/j.cpc.2015.06.008.
2. B. B. Godfrey, J.-L. Vay, "Suppressing the numerical Cherenkov instability in FDTD PIC codes", Journal of Computational Physics, 267, 1-6 (2014) http://dx.doi.org/10.1016/j.jcp.2014.02.022
3. B. B. Godfrey, J.-L. Vay, I. Haber, "Numerical Stability Improvements for the Pseudospectral EM PIC Algorithm," IEEE Transactions on Plasma Science 42, 1339-1344 (2014) http://dx.doi.org/10.1109/TPS.2014.2310654
4. B. B. Godfrey, J.-L. Vay, I. Haber, "Numerical stability analysis of the pseudo-spectral analytical time-domain PIC algorithm" , J. Comput. Phys. 258, 689-704 (2014) http://dx.doi.org/10.1016/j.jcp.2013.10.053
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