

Fission

Fission

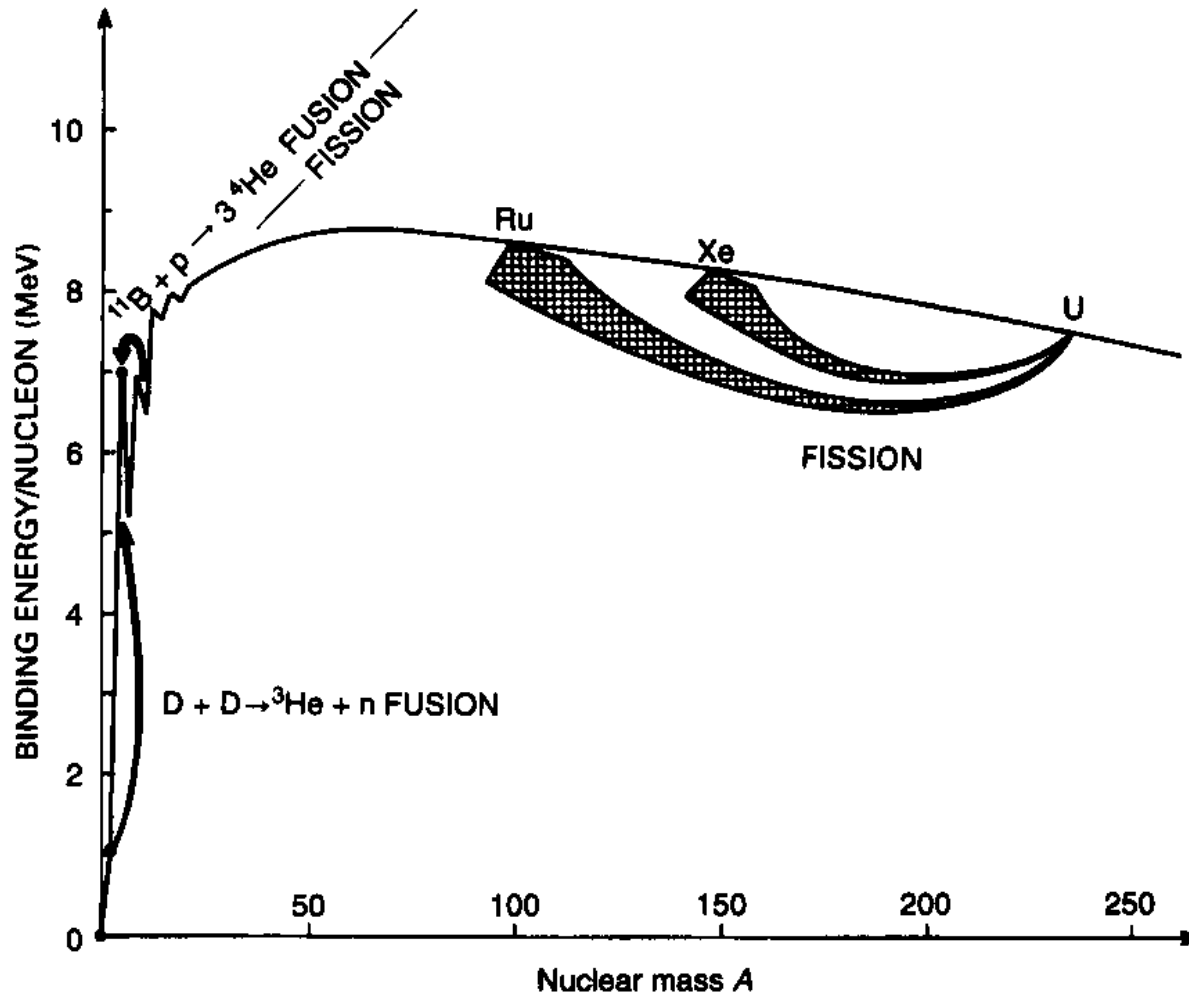
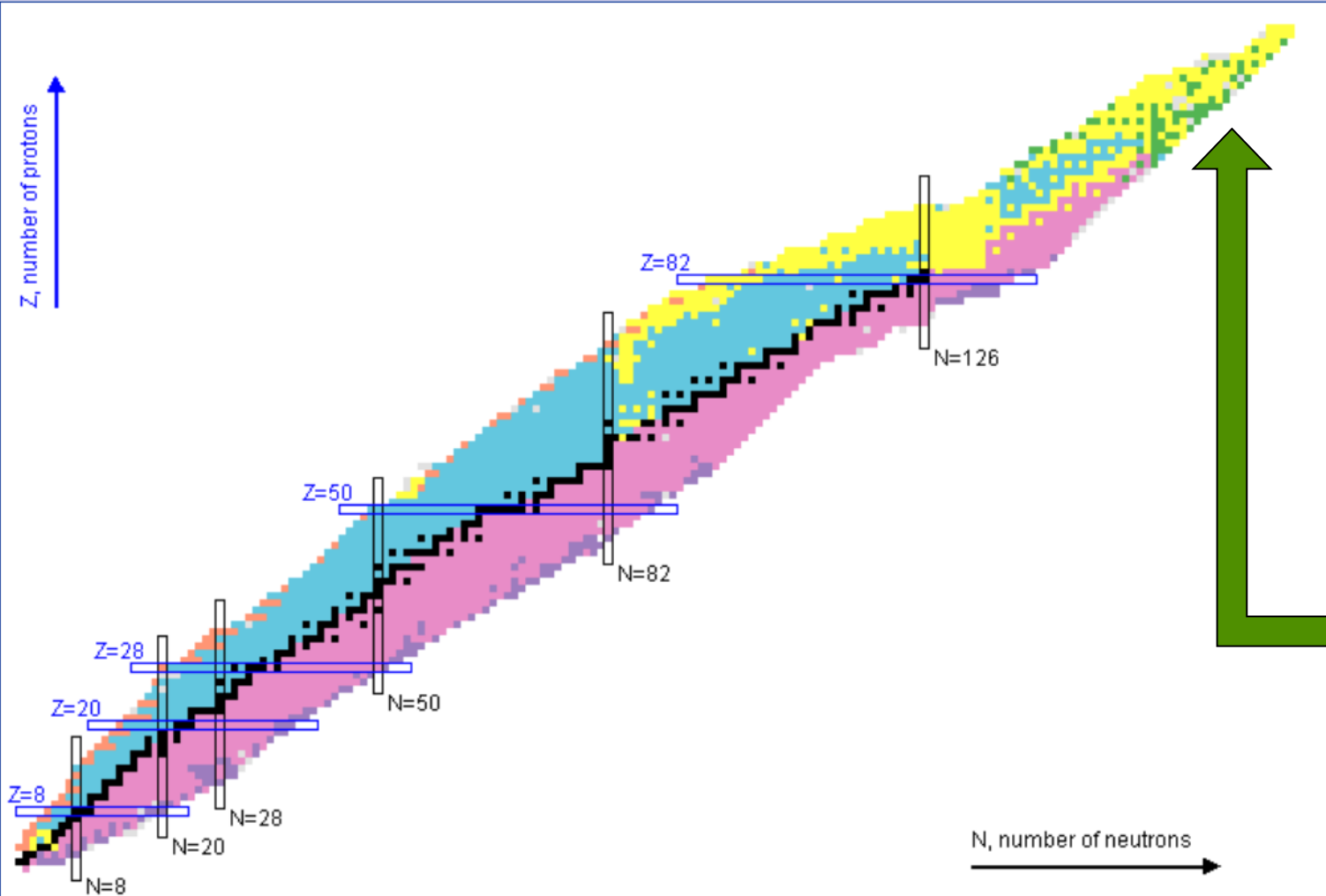




Chart of Nuclides

Click on a nucleus for information

Color code	Half-life	Decay Mode	Q_{β^-}	Q_{EC}	Q_{β^+}	S_n	S_p	Q_{α}	S_{2n}	S_{2p}	$Q_{2\beta^-}$	Q_{2EC}	Q_{ECp}
$Q_{\beta-n}$	BE/A	(BE-LDM Fit)/A	$E_{1st\ ex. st.}$	E_{2+}	E_{3-}	E_{4+}	E_{4+}/E_{2+}	β_2	$B(E2)_{42}/B(E2)_{20}$	$\sigma(n,\gamma)$	$\sigma(n,F)$	235U FY	239Pu FY



Tooltips
 On
 Off

Zoom
 1
 2
 3
 4
 5
 6
 7

Uncertainty
 NDS
 Standard

Screen Size
 Narrow
 Wide

Nucleus

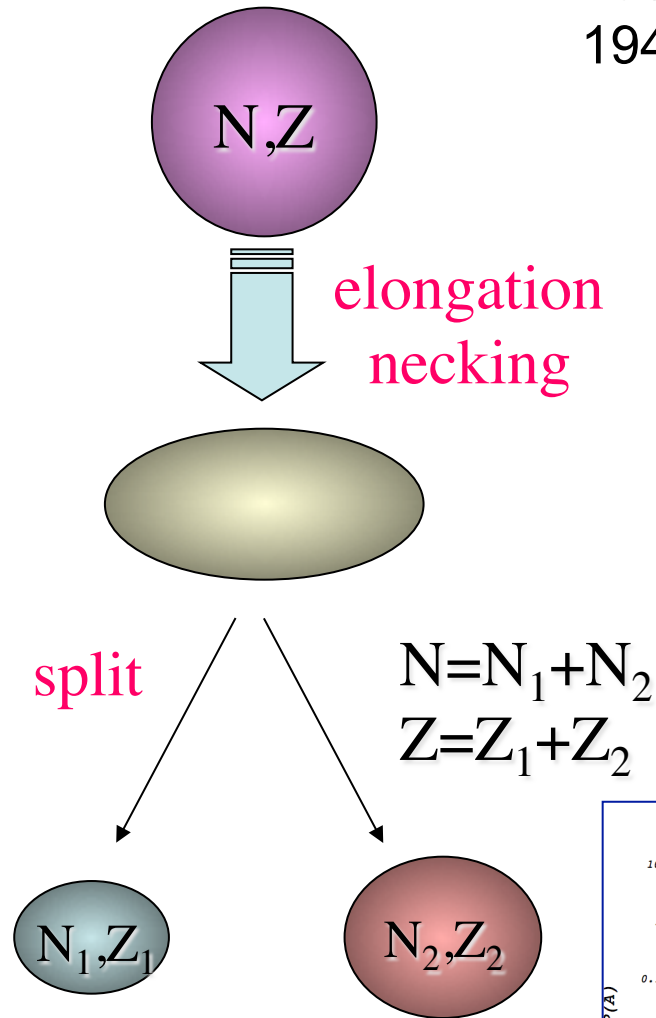
- Stable
- EC+β+
- β-
- α
- P
- N
- SF
- Unknown

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1938 Hahn & Strassmann
 1939 Meitner & Frisch
 1939 Bohr & Wheeler
 1940 Petrzhak & Flerov



Fission yields (fragments)

The neutron strikes the nucleus and is absorbed.

The absorbed neutron causes the nucleus to undergo deformation.

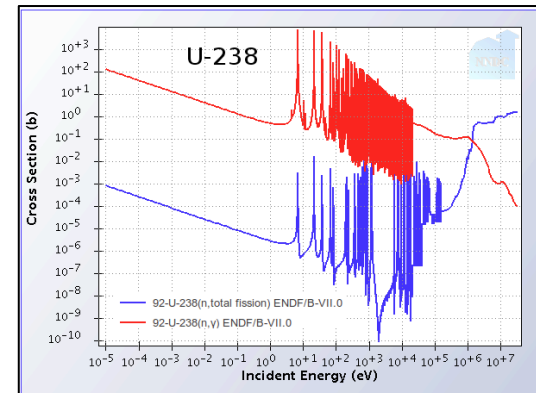
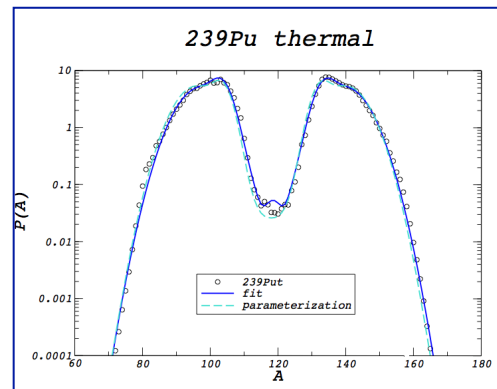
In about 10^{-14} second, one of the deformations is so drastic that the nucleus cannot recover.

The nucleus fissions, releasing two or more neutrons.

In about 10^{-12} second, the fission fragments lose their kinetic energy and come to rest, emitting a number of gamma rays. Now the fragments are called fission products.

The fission products lose their excess energy by radioactive decay, emitting beta particles and gamma rays over a lengthy time period (seconds to years).

● neutrons ● protons ● beta particles 🌊 gamma rays



Understanding the fission process is crucial for many areas of science and technology:

- Fission governs the production and existence of many transuranium elements, including the predicted long-lived super-heavy species.
- Fission influences the formation of heavy elements in a neutron rich environment.
- Fission produces reactor antineutrinos
- Improved understanding of the fission process will enable scientists to enhance the safety and reliability of nuclear reactors.
- Fission is important for stockpile stewardship

The new phase in fission theory is expected to rely heavily on advanced modeling and simulation capabilities utilizing massively parallel leadership-class computers



Letters to the Editor

The Editor does not hold himself responsible for opinions expressed by his correspondents. He cannot undertake to return, or to correspond with the writers of, rejected manuscripts intended for this or any other part of NATURE. No notice is taken of anonymous communications.

NOTES ON POINTS IN SOME OF THIS WEEK'S LETTERS APPEAR ON P. 337.

CORRESPONDENTS ARE INVITED TO ATTACH SIMILAR SUMMARIES TO THEIR COMMUNICATIONS.

Disintegration of Heavy Nuclei

THROUGH the kindness of the authors I have been informed of the content of the letters¹ recently sent to the Editor of NATURE by Prof. Meitner and Dr. Frisch. In the first letter, these authors propose an interpretation of the remarkable findings of Hahn and Strassmann as indication for a new type of disintegration of heavy nuclei, consisting in a fission of the nucleus into two parts of approximately equal masses and charges with release of enormous energy.

In the second letter, Dr. Frisch describes experimentally these parts are directly detected by their ionization they produce. Due to the importance of this discovery, I should be glad to give a few comments on the mechanism of the process from the point of view of the general ideas of recent years, to account for the results of the nuclear reactions hitherto known.

In accordance with these ideas, any nuclear reaction induced by collisions or radiation involves as an intermediate stage the formation of a compound nucleus in which the excitation energy is distributed among the various degrees of freedom of the nucleus, including the thermal agitation of a solid body. The relative probabilities of the various possible courses of the reaction will therefore depend on the facility with which this energy can be released as radiation or converted into a form which produces the disintegration of the nucleus. In the case of ordinary reactions, the disintegration consists in the escape of a particle, i.e., this conversion means the concentration of a part of the energy on some particle of the nucleus, and resembles therefore the evaporation of a molecule from a liquid drop. In the case of disintegrations comparable to those which split a drop into two droplets, it is evident that a large amount of energy must be converted into a form of vibration of the compound nucleus involving a considerable deformation of the nucleus.

It is possible, however, that the quasi-thermal distribution of energy may be largely converted into a mode of vibration of the compound nucleus involving a considerable deformation of the nucleus. In such cases, the course of the disintegration may be determined to a large extent by the fluctuation in the distribution of the energy between the various degrees of freedom of the system, the probability of occurrence of which is essentially determined by the amount of energy to be concentrated on a particular type of motion considered and by the probability corresponding to the nuclear structure of the nucleus. Since the effective cross-sections for the capture of neutrons of different velocities are of about the same order of magnitude as those for ordinary nuclear reactions, we may conclude that for the heaviest nuclei an energy sufficient for the fission is of

the same order of magnitude as the energy necessary for the escape of a single nuclear particle. For somewhat lighter nuclei, however, where only evaporation-like disintegrations have so far been observed, the former energy should be considerably larger than the binding energy of a particle.

These circumstances find their straightforward explanation in the fact, stressed by Meitner and Frisch, that the mutual repulsion between the electric charges in a nucleus will for highly charged nuclei counteract to a large extent the effect of the short-range forces between the nuclear particles in opposing a deformation of the nucleus. The nuclear problem concerned reminds us indeed in several ways of the question of the stability of a charged liquid drop, and in particular, any deformation of a nucleus, sufficiently large for its fission, may be treated approximately as a classical mechanical problem, since the corresponding amplitude must evidently be large compared with the quantum mechanical zero-point oscillations. Just this condition would in fact seem to provide an understanding of the remarkable stability of heavy nuclei in their normal state or in the states of low excitation, in spite of the large amount of energy which would be liberated by an imaginable division of such nuclei.

The continuation of the experiments on the new type of nuclear disintegrations, and above all the closer examination of the conditions for their occurrence, should certainly yield most valuable information as regards the mechanism of nuclear excitation.

N. BOHR.

At the Institute for Advanced Study,
Princeton, N.J. Jan. 20.

¹ (NATURE, 143, 239 and 275 (1939)).

Photoactivation of Solids and its Effect on Adsorption

CONSIDERABLE attention has been given recently to chemical processes involving an activating influence of a crystal excited by irradiation¹. The mechanism of such photosensitized reactions, although unknown in detail, is generally believed to be a more or less complete transfer of the energy absorbed by the crystal to the reacting components, physically or chemically. Accordingly, the essential difference between photosensitized processes and real photochemical ones is the distance between the place of absorption and the place of reaction. But there must also be another, more general, effect of irradiation on the activity of crystals. Due to the change in the electronic state of the particles in the lattice by the absorption of light (change in charge and degree of polarization, formation of space charges, etc.), the forces between the particles are changed and, consequently, there is also a change in potential of the

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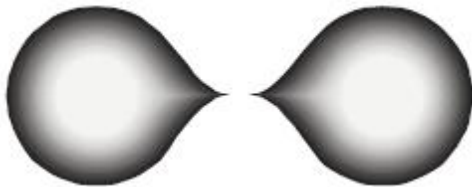
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Deformed liquid drop (Bohr & Wheeler, 1939)

fission of nuclear
droplet



$$E_{LDM}(def) = E_S(0) \left[B_S(def) - 1 + 2x(B_C(def) - 1) \right]$$

$$B_S(def) = \frac{E_S(def)}{E_S(0)}, \quad B_C(def) = \frac{E_C(def)}{E_C(0)}$$

$$x = \frac{E_C(0)}{2E_S(0)} = \frac{Z^2 / A}{(Z^2 / A)_{crit}} \approx \frac{Z^2}{50 A}$$

x: fissility parameter