

Isotope Shift

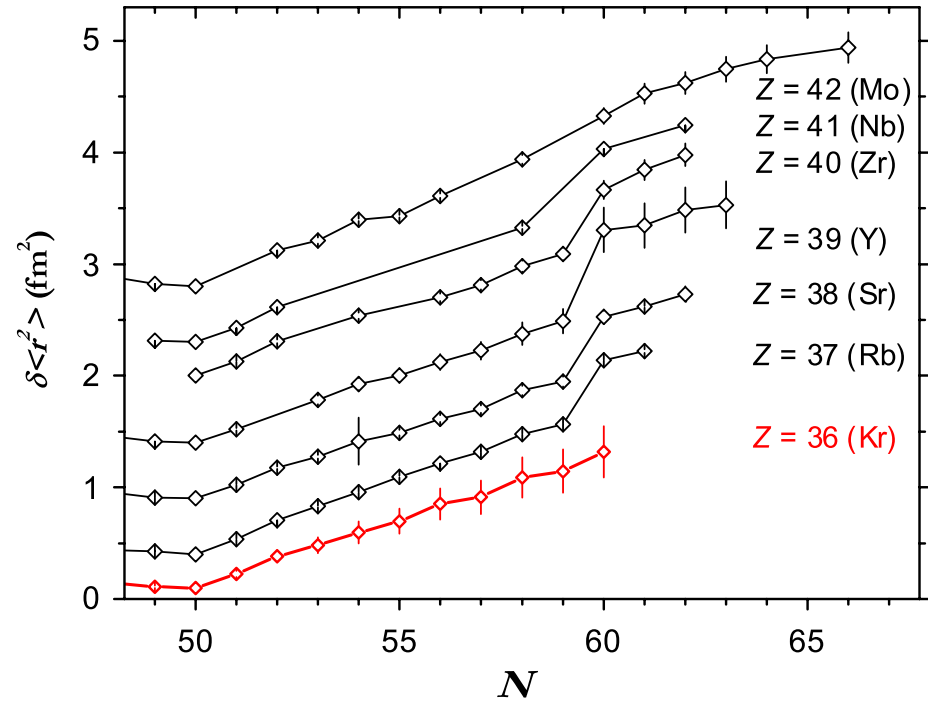
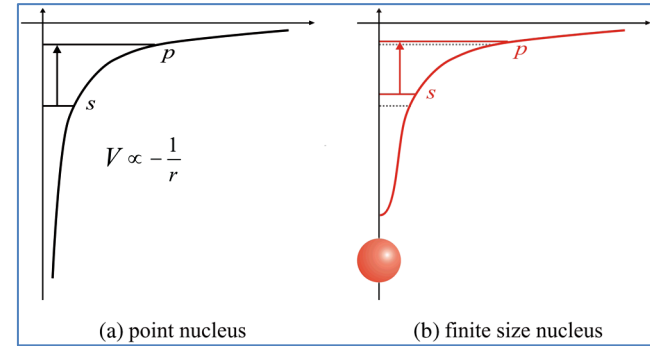
Laser trapping of exotic atoms. RMP 85, 1383 (2013)

TABLE I. Contributions to the electronic binding energy and their orders of magnitude in atomic units. a_0 is the Bohr radius, $\alpha \approx 1/137$. For helium, the atomic number $Z = 2$, and the mass ratio $\mu/M \sim 1 \times 10^{-4}$. g_I is the nuclear g factor. α_d is the nuclear dipole polarizability.

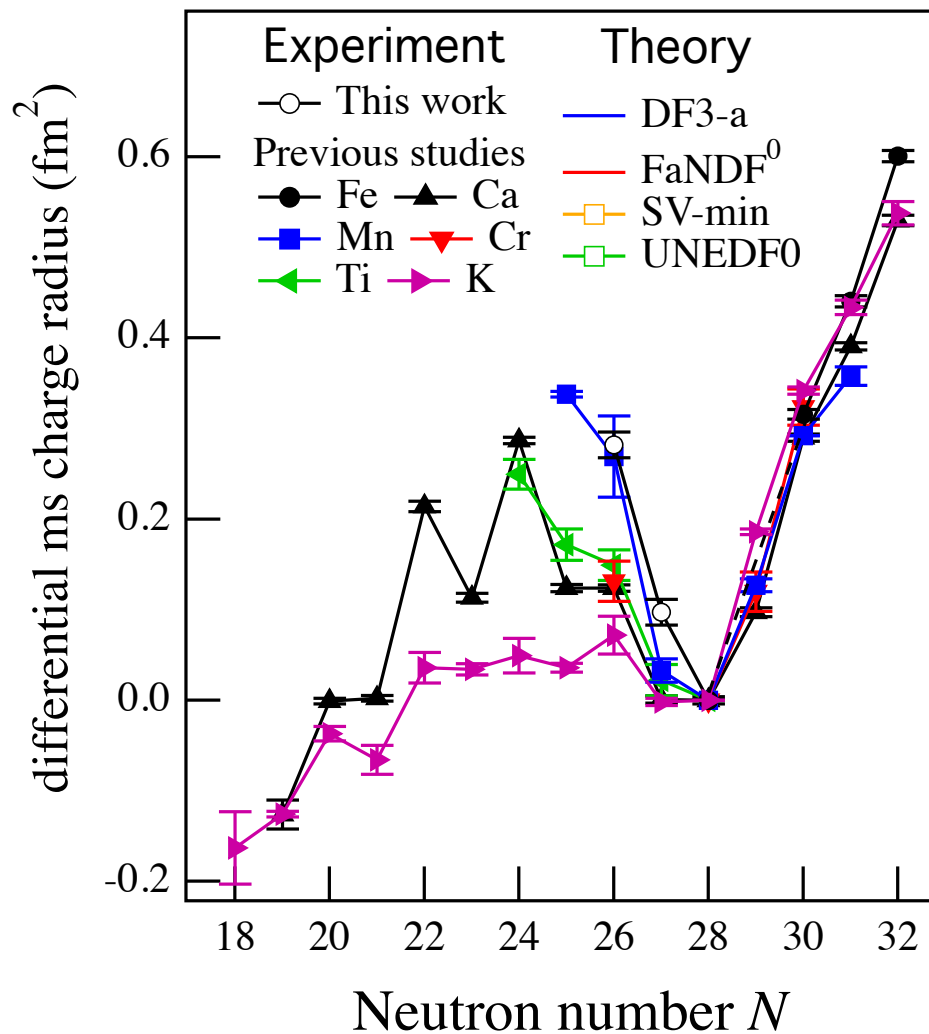
Contribution	Magnitude
Nonrelativistic energy	Z^2
Mass polarization	$Z^2 \mu/M$
Second-order mass polarization	$Z^2 (\mu/M)^2$
Relativistic corrections	$Z^4 \alpha^2$
Relativistic recoil	$Z^4 \alpha^2 \mu/M$
Anomalous magnetic moment	$Z^4 \alpha^3$
Hyperfine structure	$Z^3 g_I \mu_0^2$
Lamb shift	$Z^4 \alpha^3 \ln \alpha + \dots$
Radiative recoil	$Z^4 \alpha^3 (\ln \alpha) \mu/M$
Finite nuclear size	$Z^4 \langle r_c / a_0 \rangle^2$
Nuclear polarization	$Z^3 e^2 \alpha_d / (\alpha a_0^4)$

$$\alpha = \frac{1}{137}$$

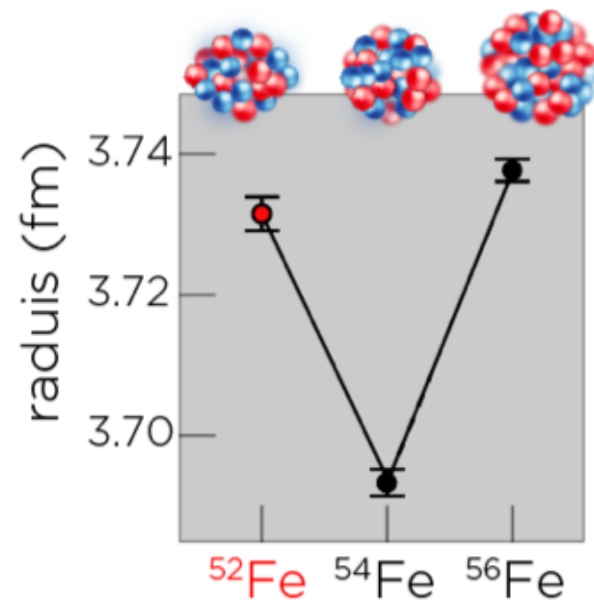
μ =reduced electron mass



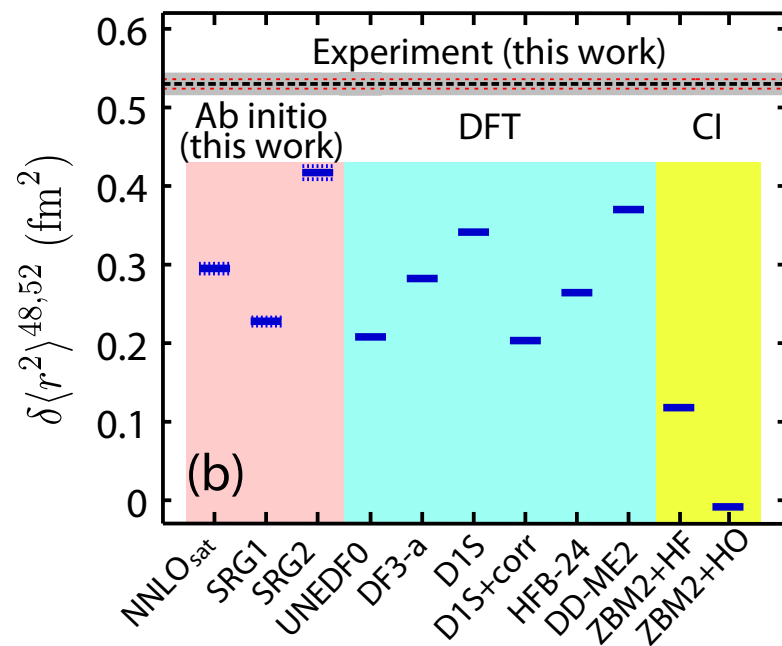
Difference in mean-square charge radii for the $N \sim 60$ region, PRL 105, 032502 (2010)



Phys. Rev. Lett. 117, 252501 (2016)
 BECOLA @ NSCL



Nature Physics 12, 594 (2016)



Neutron radii

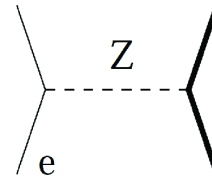
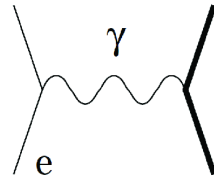
- Proton-Nucleus elastic
 - Pion, alpha, d scattering
 - Pion photoproduction
- } Involve strong probes

Phys. Rev. Lett. 112, 242502 (2014)

<http://physics.aps.org/synopsis-for/10.1103/PhysRevLett.112.242502>

Parity-violating electron scattering

Z₀ of Weak Interaction



$M_Z = 90.19 \text{ GeV!}$

Parity Violating Asymmetry

$$A = \frac{\left(\frac{d\sigma}{d\Omega}\right)_R - \left(\frac{d\sigma}{d\Omega}\right)_L}{\left(\frac{d\sigma}{d\Omega}\right)_R + \left(\frac{d\sigma}{d\Omega}\right)_L} = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \left[\underbrace{1 - 4\sin^2\theta_W}_{\approx 0} - \frac{F_n(Q^2)}{F_p(Q^2)} \right] \sim 7 \cdot 10^{-7}$$

Weinberg angle: $\sin^2\theta_W = 0.23120 \pm 0.00015$

	proton	neutron
Electric charge	1	0
Weak charge	0.08	1

A comment: Yukawa potential

$$V_Y(r) = -g^2 \frac{e^{-\mu r}}{r}$$

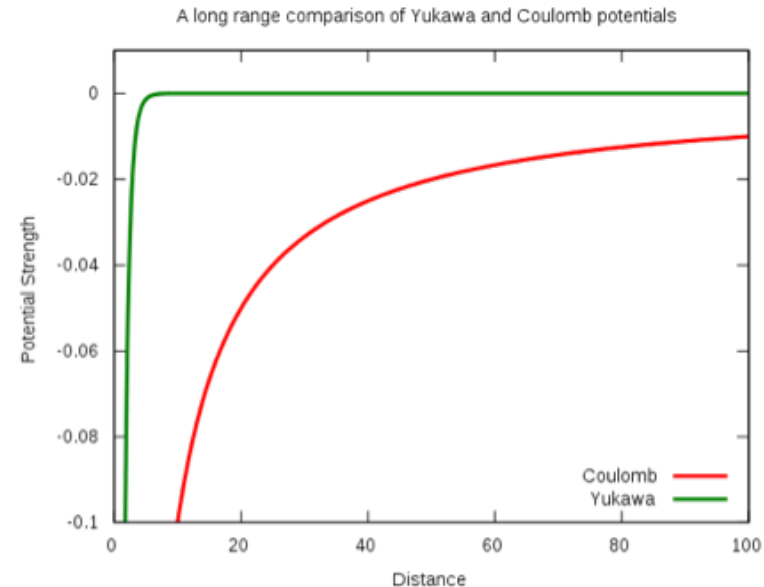
$$\lambda_B = \frac{1}{\mu} = \frac{\hbar}{m_B c}$$

Compton wavelength of
the boson (force carrier)

Mass of the boson

$$-\frac{1}{c^2} \frac{\partial^2 \Psi}{\partial t^2} + \nabla^2 \psi = \mu^2 \Psi$$

Klein-Gordon equation



Lead (^{208}Pb) Radius Experiment : PREX

Analysis is clean, like electromagnetic scattering:

1. Probes the entire nuclear volume
2. Perturbation theory applies

$E = 850 \text{ MeV}, \theta = 6^\circ$
electrons on lead

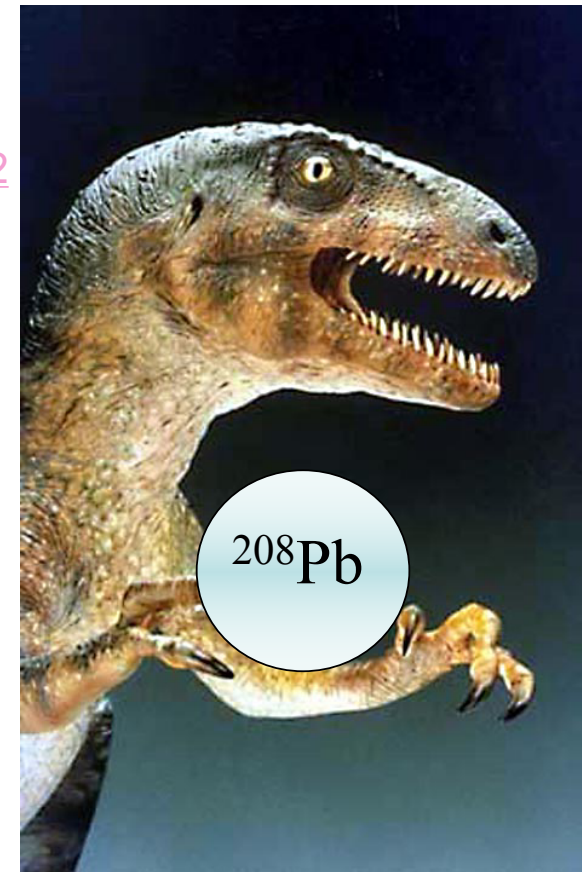
Phys. Rev. Lett. 108, 112502 (2012)

<http://physics.aps.org/synopsis-for/10.1103/PhysRevLett.108.112502>

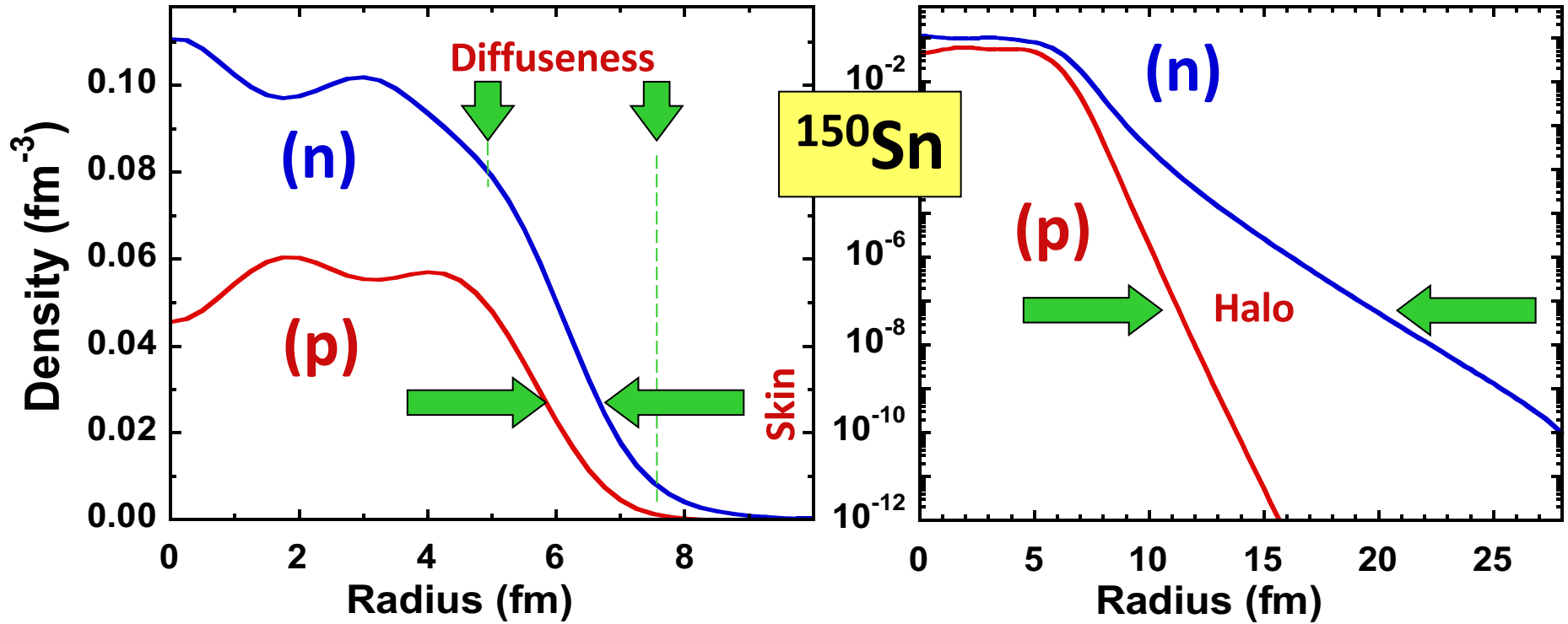
$$R_{skin} = R_n - R_p$$

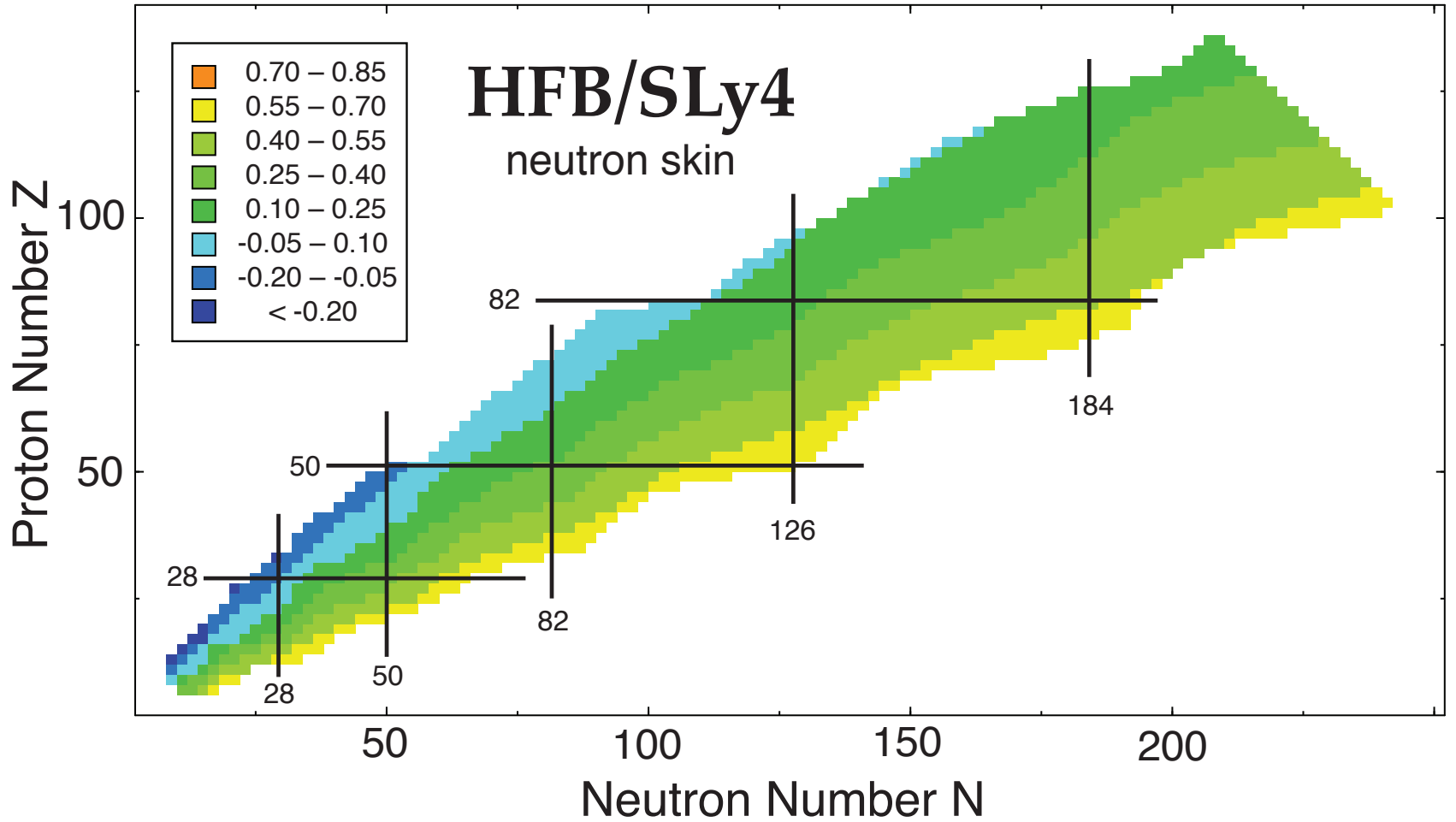
PREX: $0.34^{+0.15}_{-0.17} \text{ fm}$

Theory: $0.168 \pm 0.022 \text{ fm}$



Neutron & proton density distributions







HW: Assuming the nucleus of mass number A to be a spherical object with a sharp surface and constant nucleonic density $\rho_0 = 0.16$ nucleons/ fm^3 , find the relation between nuclear radius and A .

Test the performance of the resulting expression by comparing with experimental data for charge radii:

<http://www.sciencedirect.com/science/article/pii/S0092640X12000265>

Assume that the radius of the mass distribution is the same as the radius of the charge distribution. Note that this reference discusses root-mean-square (rms) nuclear charge radii not geometric radii.