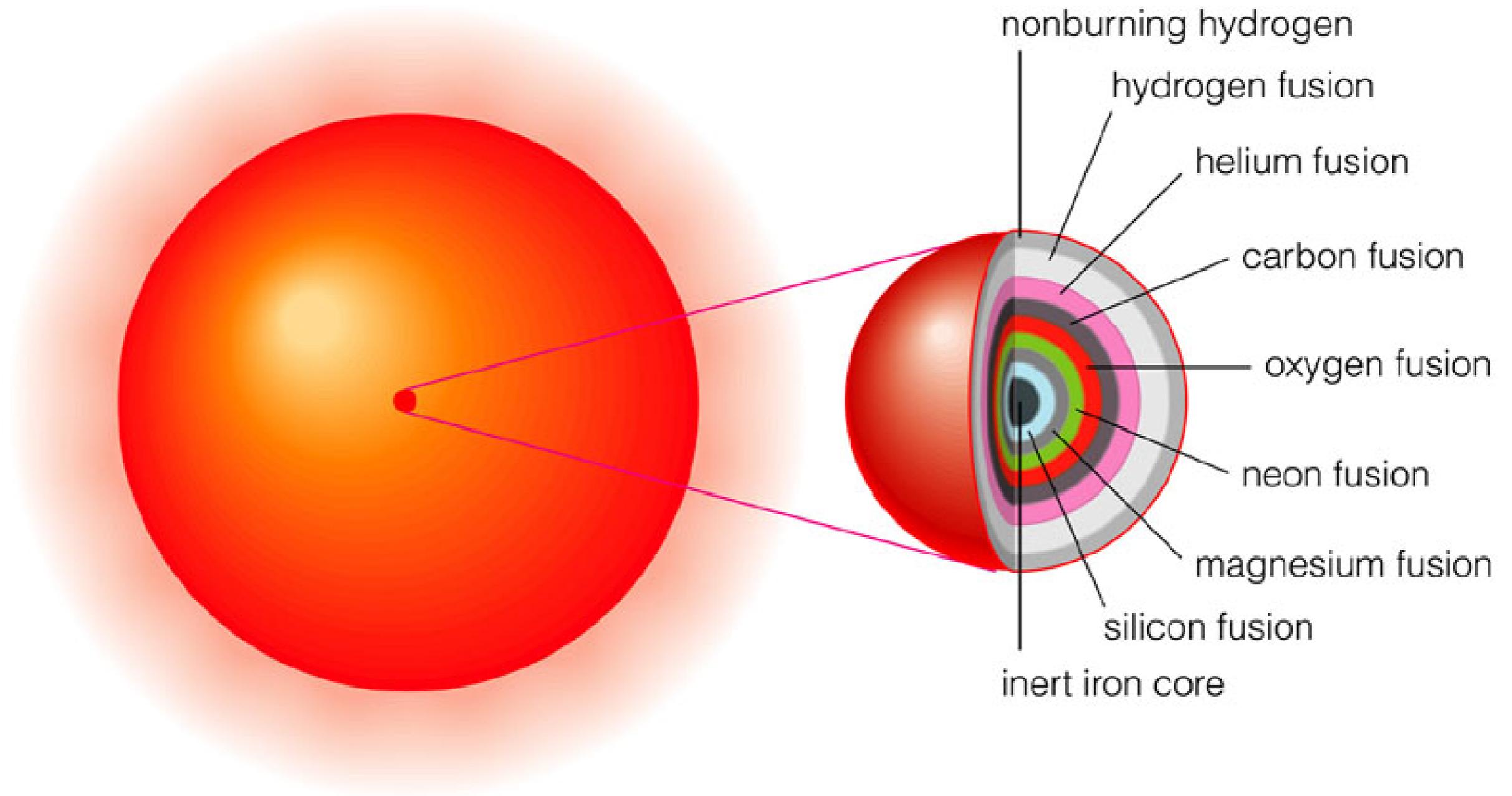


Textbook view of massive stars

Essential Cosmic Perspective, Bennett et al.



5. *The super-nova process*

We have tentatively suggested that the super-nova process represents the transition of an ordinary star into a neutron star. If neutrons are produced on the surface of an ordinary star they will "rain" down towards the center if we assume that the light pressure on neutrons is practically zero. This view explains the speed of the star's transformation into a neutron star. We are fully aware that our suggestion carries with it grave implications regarding the ordinary views about the constitution of stars and therefore will require further careful studies.

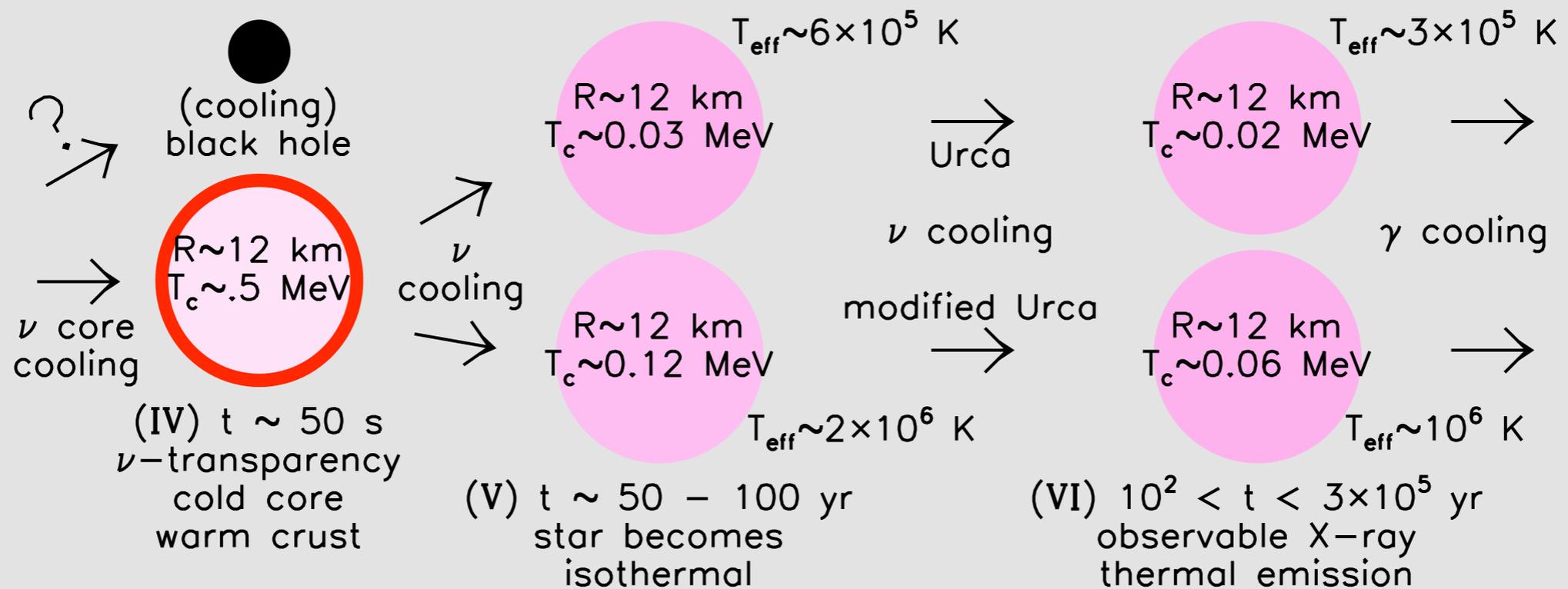
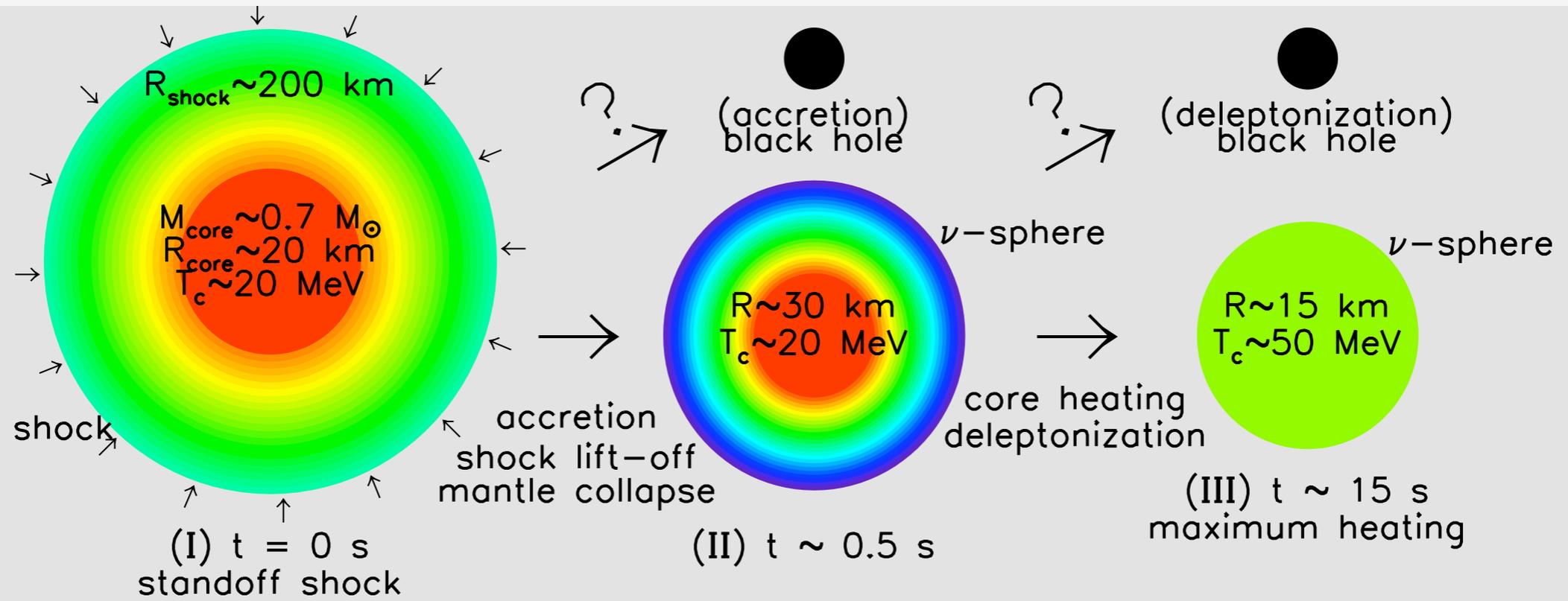
W. BAADE

F. ZWICKY

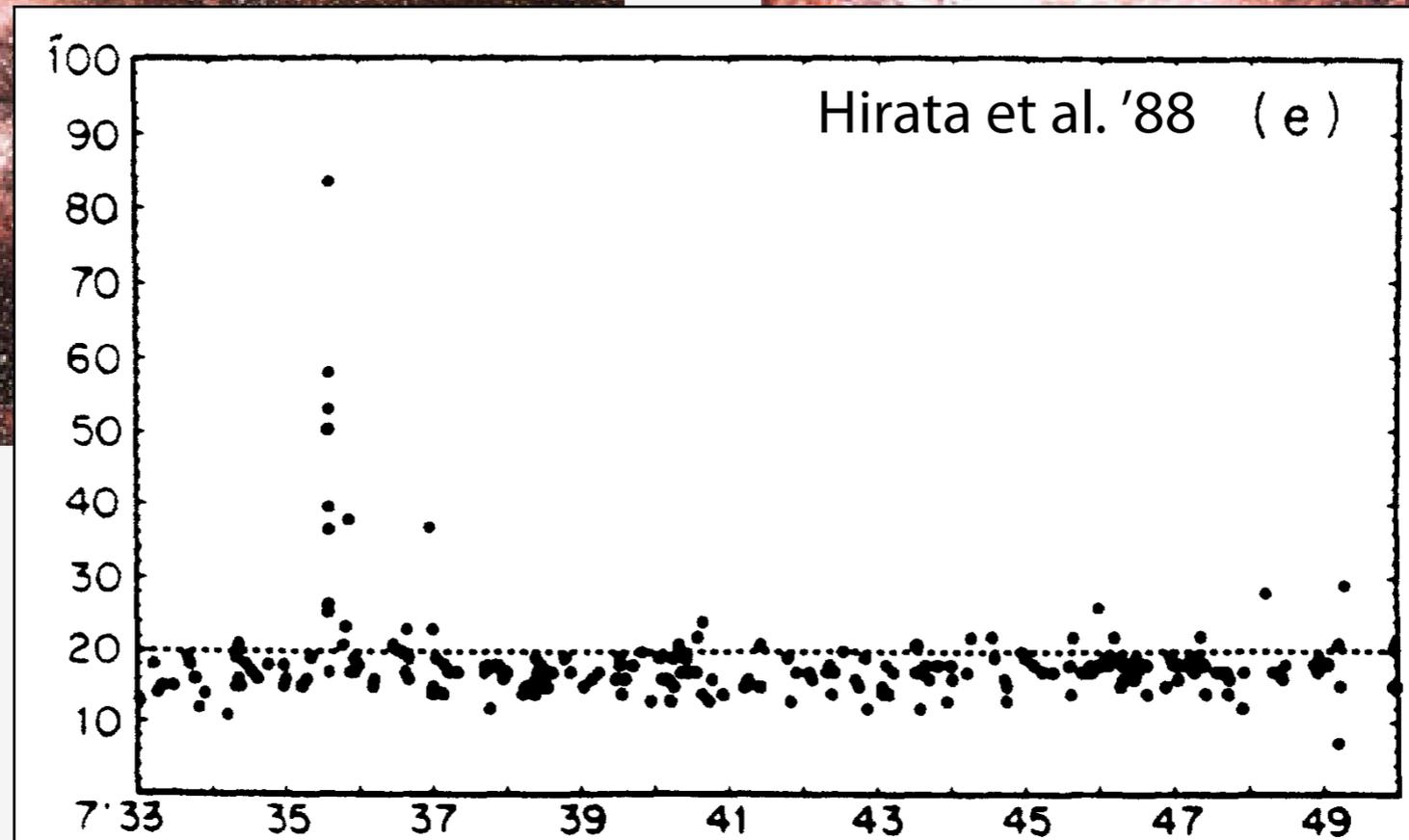
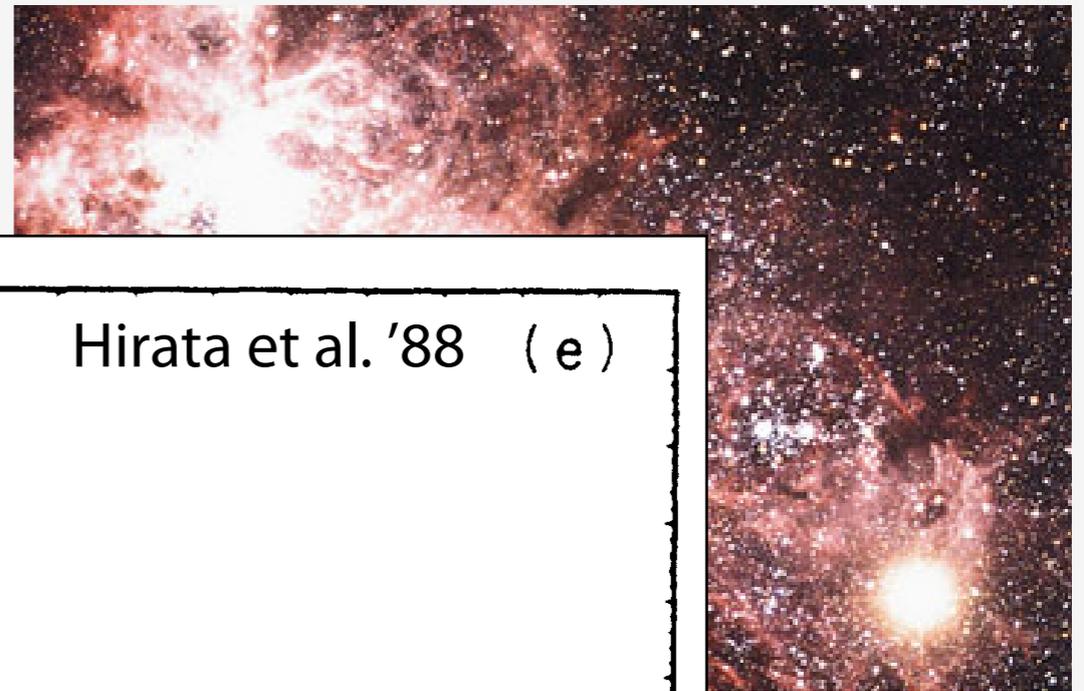
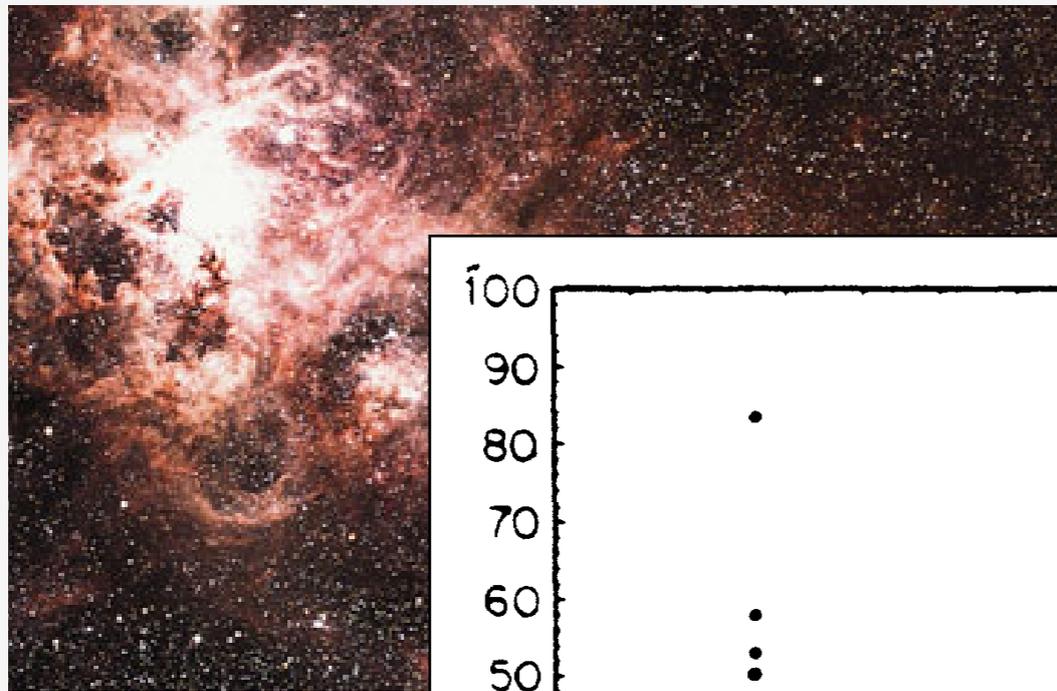
Mt. Wilson Observatory and

California Institute of Technology, Pasadena.

May 28, 1934.

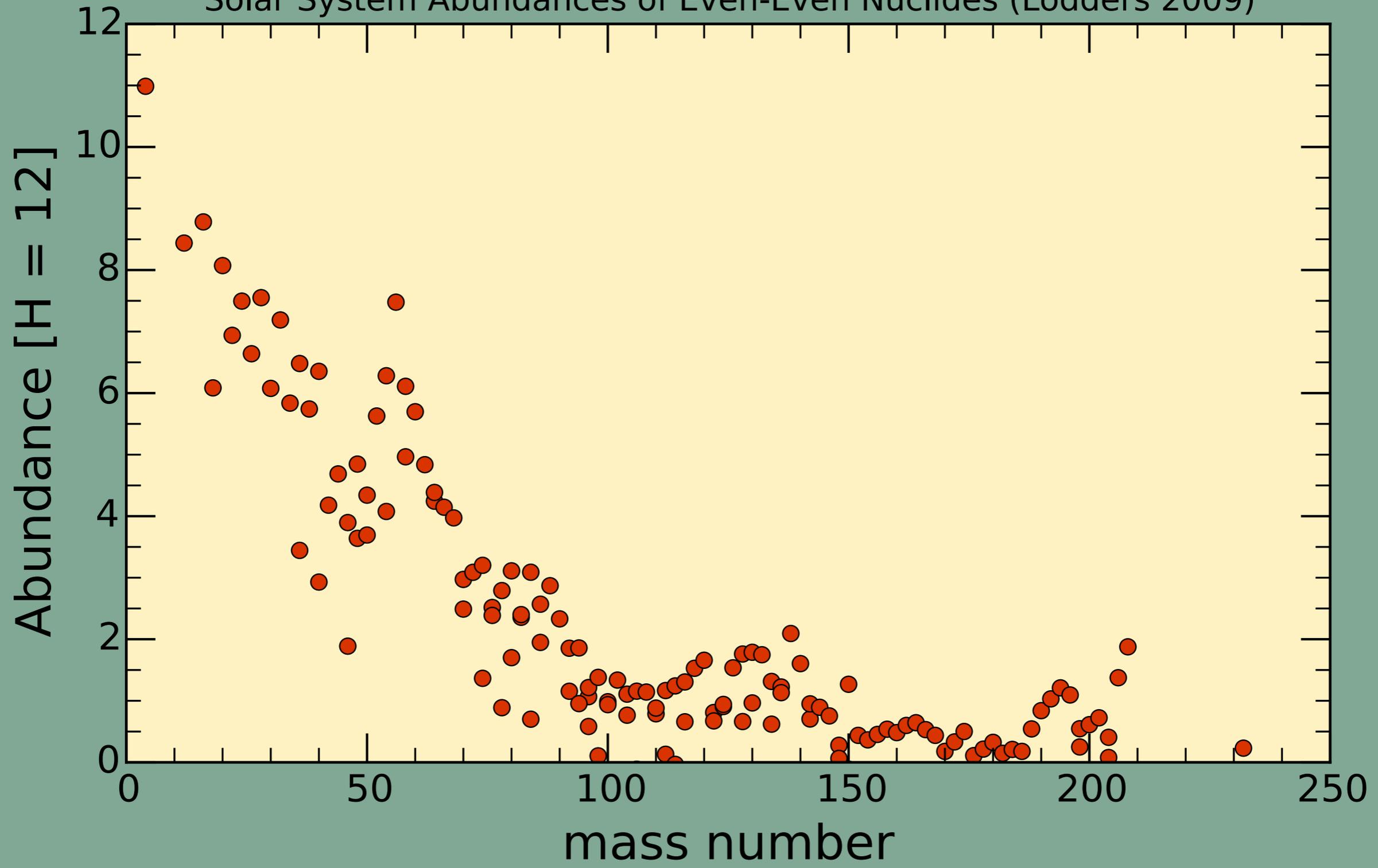


Modern view of neutron star birth
 schematic from Lattimer and Prakash

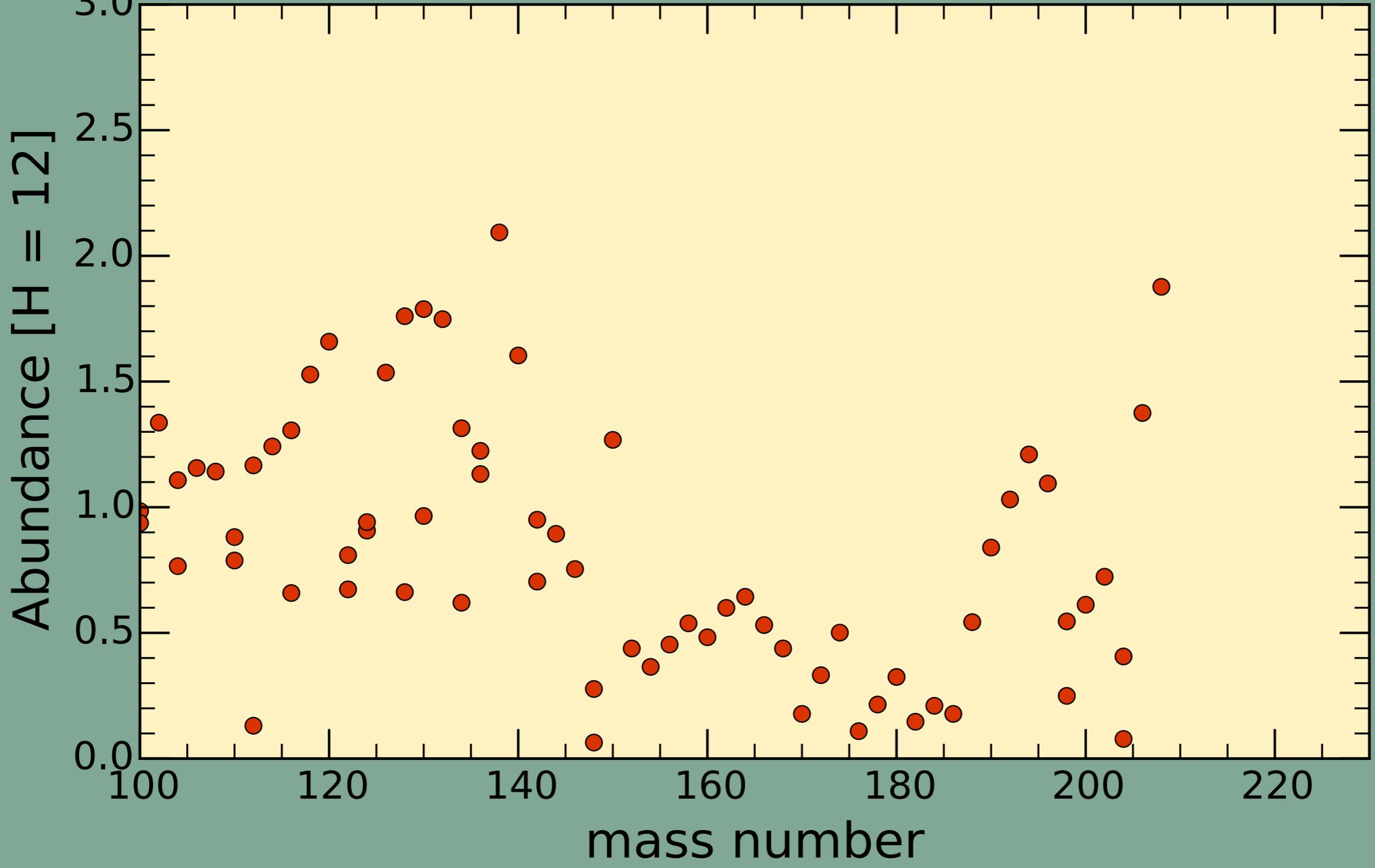


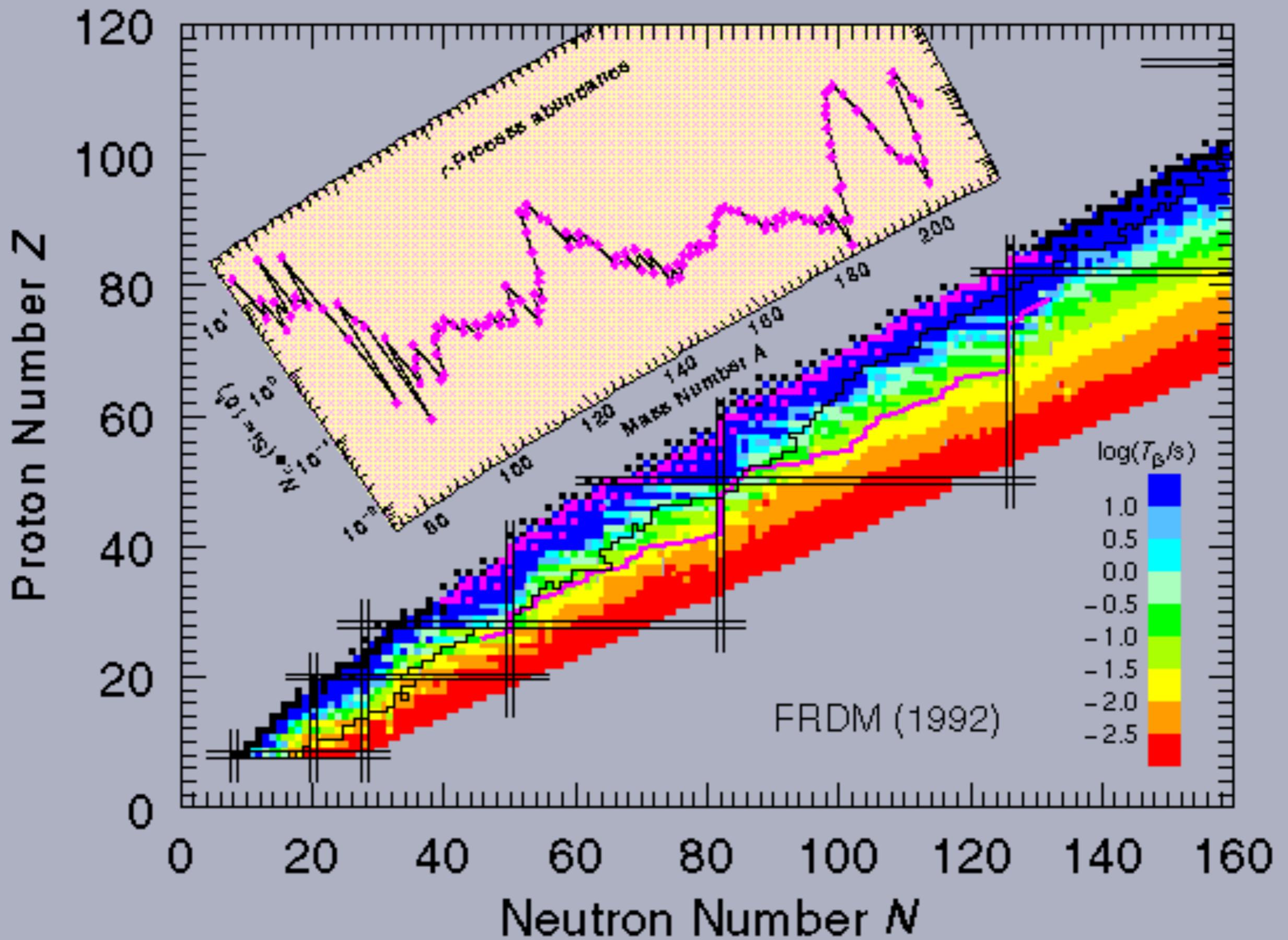
Supernova 1987a: neutrinos detected by Kamiokande!

Solar System Abundances of Even-Even Nuclides (Lodders 2009)

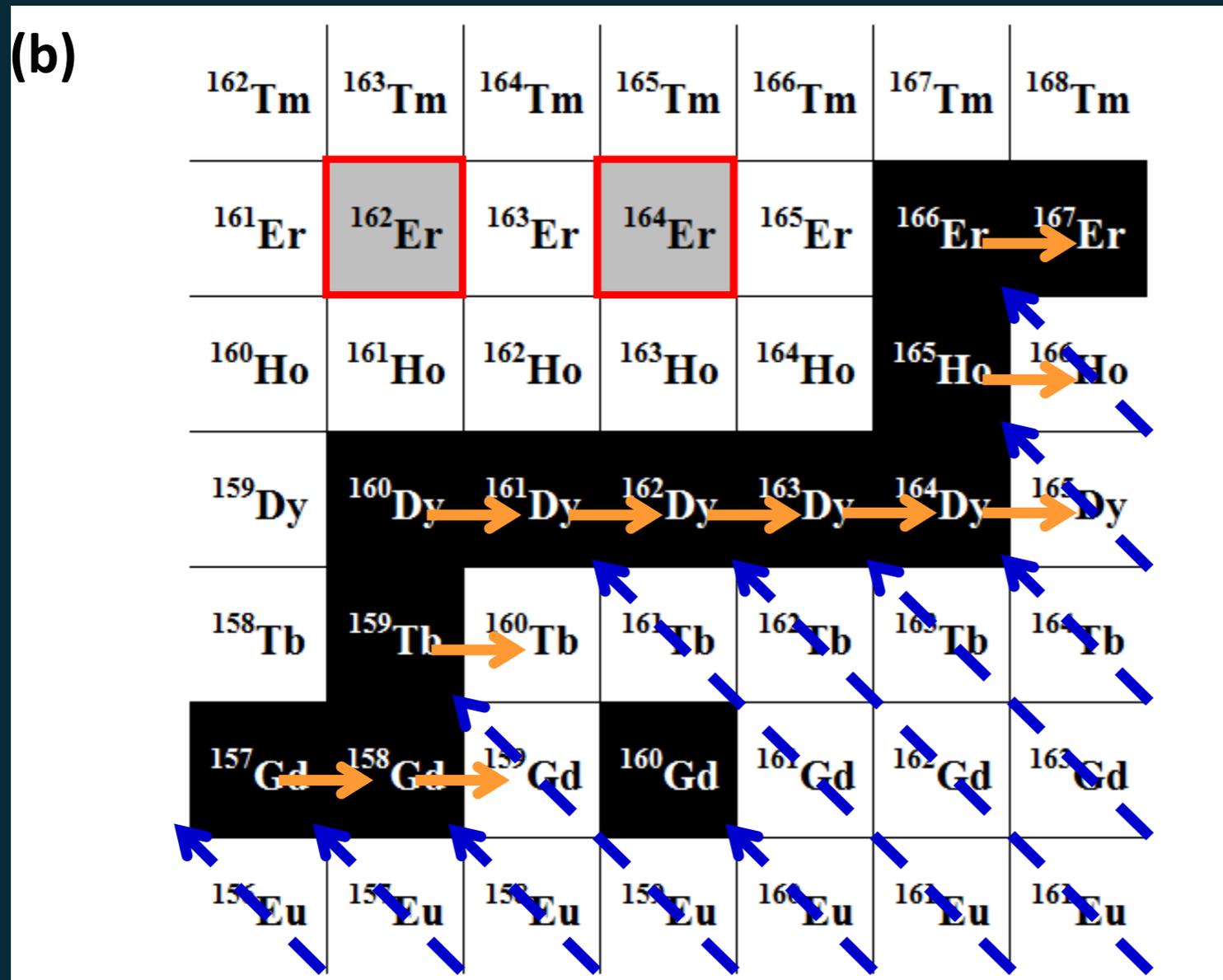


Solar System Abundances of Even-Even Nuclides (Lodders 2009)



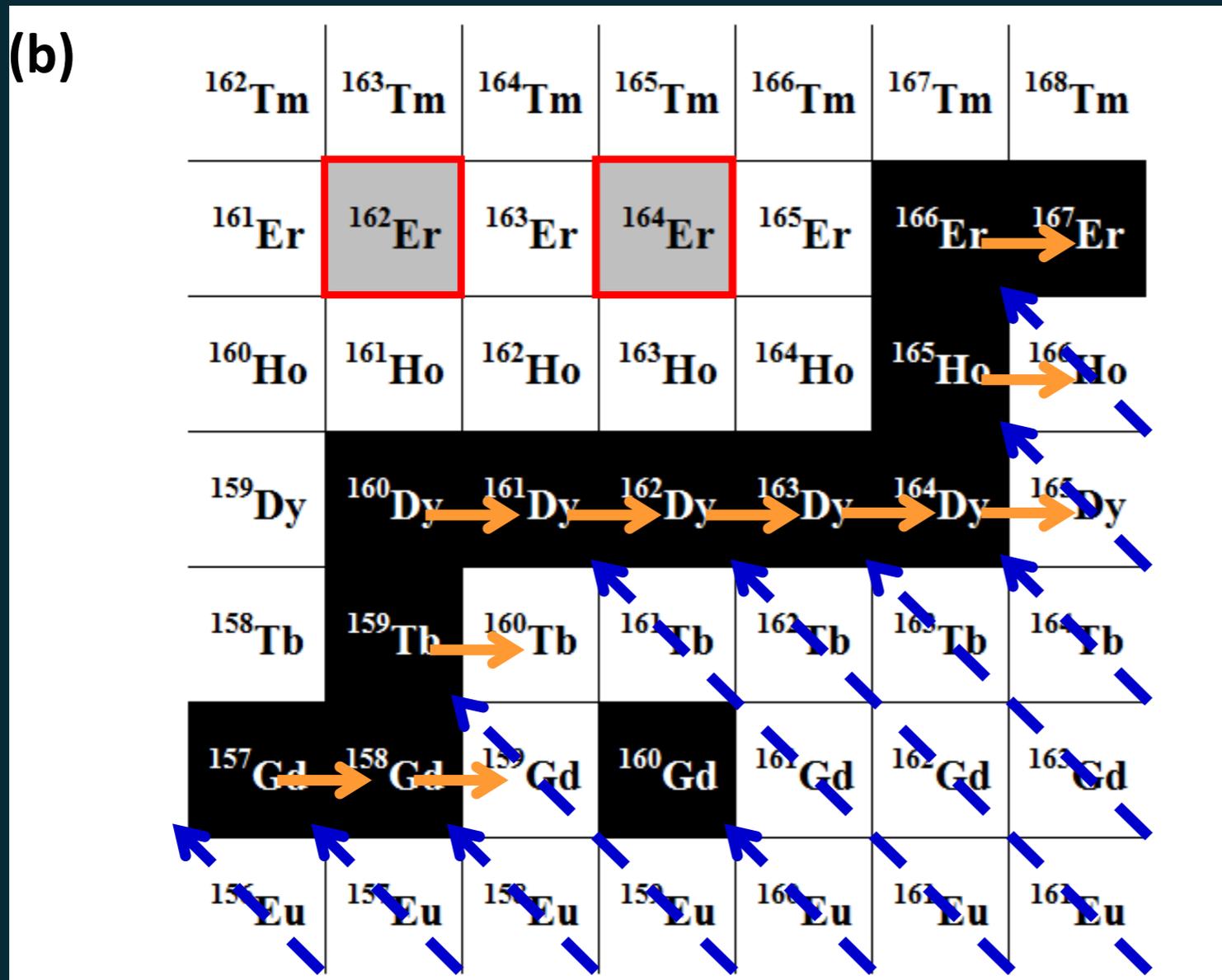


which nuclei are s-only?



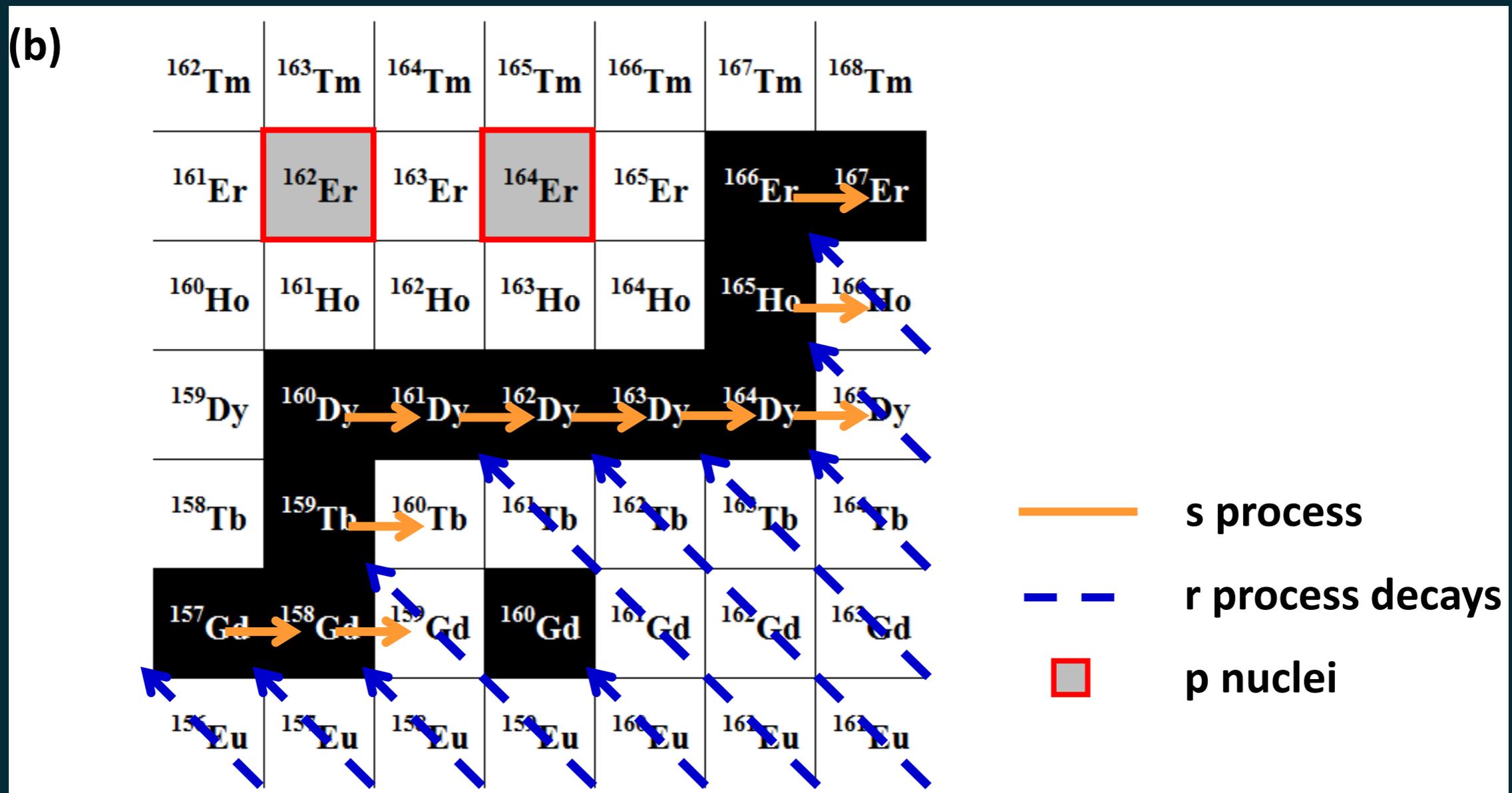
S. Quinn 2015

which nuclei are r-only?



S. Quinn 2015

which nuclei are r-only?



S. Quinn 2015

neutron star basics

A solar mass consists of $\sim 10^{57}$ nucleons. If they are separated by typical inter-nucleon distances, what would the radius of the volume containing them be?

neutron star basics



$$R \sim 1 \text{ fm} \cdot (10^{57})^{1/3} \sim 10 \text{ km}$$

$$\frac{GM}{Rc^2} \approx 0.2 \approx \frac{\Delta\lambda}{\lambda}$$

$$\frac{GMm_H}{R} \approx 200 \text{ MeV}$$

from nuclei to bulk matter

$$\frac{E}{A} \approx -a_V + a_S A^{-1/3} + a_A \left(\frac{N-Z}{N+Z} \right)^2 + a_C \frac{Z^2}{A^{4/3}}$$

Thermodynamics near saturation density

see review by Lattimer & Prakash

Let's examine properties of npe matter near saturation density $n = 0.16 \text{ fm}^{-3}$. The proton fraction is

$$x = n_p / (n_n + n_p),$$

and charge neutrality requires that

$$n_e = n_p = xn.$$

We write the energy per nucleon as

$$\varepsilon(n, x) = \varepsilon_S(n) + \varepsilon_A(n)(1 - 2x)^2.$$

Why so neutron-rich?

In β -equilibrium,

$$\begin{aligned}\mu_e &= \mu_n - \mu_p = \left(\frac{\partial \varepsilon}{\partial n_n} \right)_{n_p} - \left(\frac{\partial \varepsilon}{\partial n_p} \right)_{n_n} \\ &= -\frac{\partial \varepsilon}{\partial x} = 4\varepsilon_A(1 - 2x).\end{aligned}$$

For relativistic electrons, $\mu_e = (3\pi^2 n_e)^{1/3} \hbar c$, so we can solve for the proton fraction

$$x = \left[6 + \frac{3\pi^2}{64} \left(\frac{\hbar c}{\varepsilon_A} \right)^3 n \right]^{-1} \approx 0.04$$

for $\varepsilon_A = 30$ MeV.

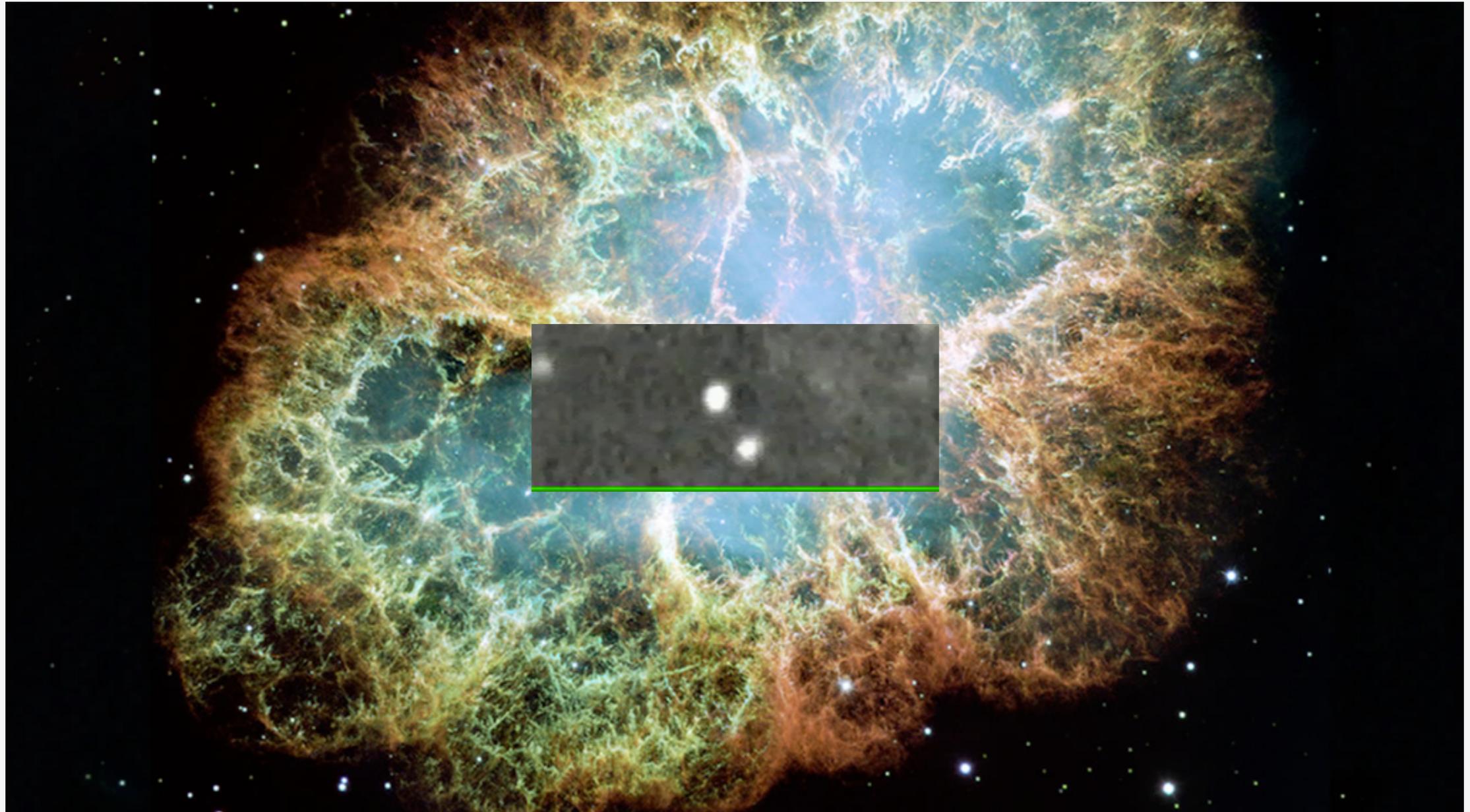
What is the pressure at saturation?

$$\begin{aligned} P &= n^2 \frac{\partial \varepsilon}{\partial n} + \frac{n_e \mu_e}{4} \\ &= \cancel{n^2 \frac{\partial \varepsilon_s}{\partial n}} + n^2 \frac{\partial \varepsilon_A}{\partial n} (1 - 2x)^2 + \varepsilon_A (1 - 2x) x n \\ &= n(1 - 2x) \left[n \frac{\partial \varepsilon_A}{\partial n} (1 - 2x) + \cancel{x \varepsilon_A} \right]. \end{aligned}$$

Discovery!

radio pulsations discovered (Hewish, Bell, et al. 1968)

Gold (1968): explained as due to rotation of a neutron star (WHY?)



Crab Nebula: Remnant of supernova in 1054

Accreting neutron stars

PHYSICAL REVIEW LETTERS

9

DECEMBER 1, 1962

NUMBER

EVIDENCE FOR X RAYS FROM SOURCES OUTSIDE THE SOLAR SYSTEM*

Riccar
Americ

Massac

THE ASTROPHYSICAL JOURNAL LETTERS TO THE EDITOR

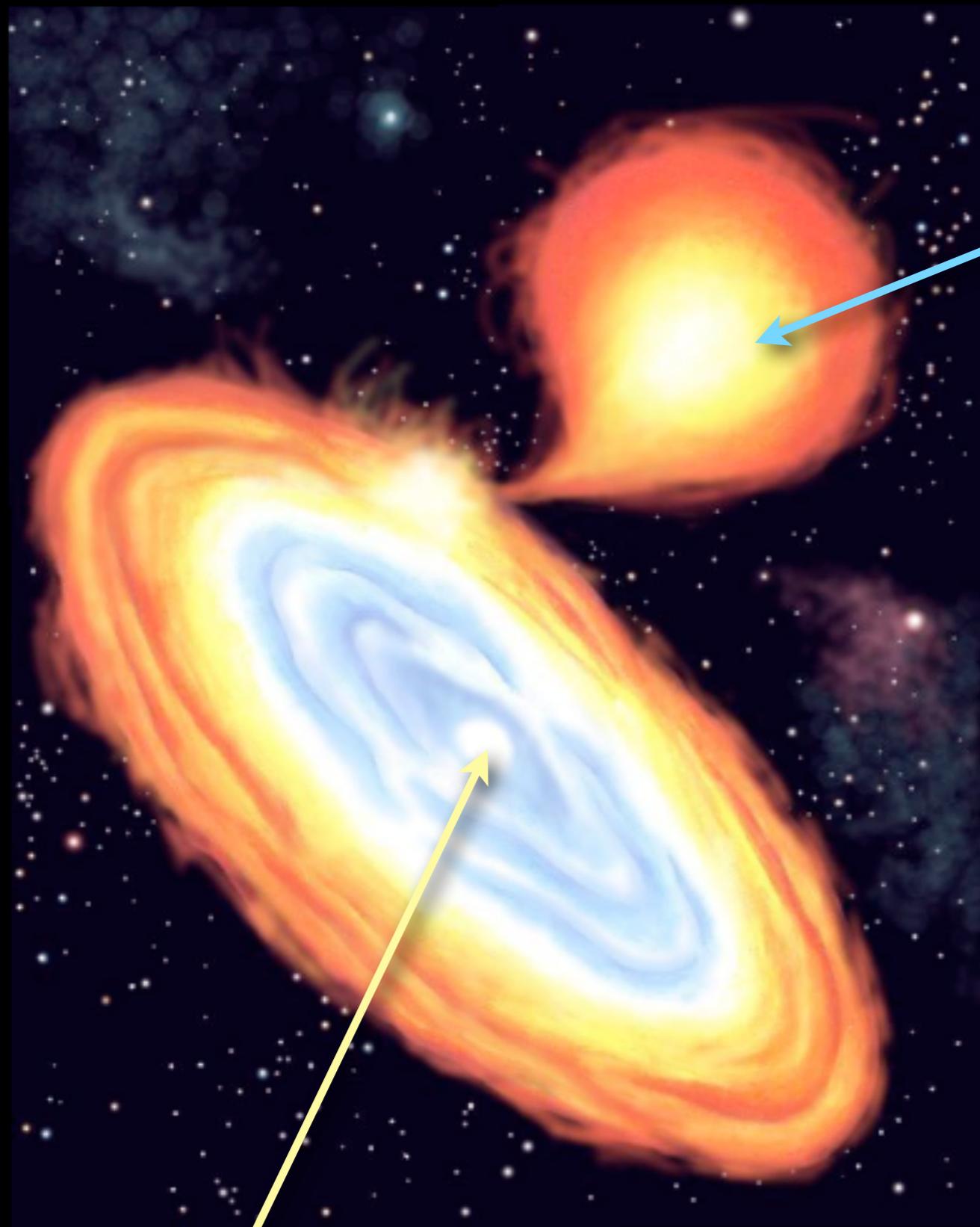
VOLUME 148

APRIL 1967

NUMBER 1, PART 2

ON THE NATURE OF THE SOURCE OF X-RAY EMISSION OF SCO XR-1

Artwork courtesy T. Piro



Neutron star

\approx solar mass star

$P_{\text{orb}} = \text{minutes-hours}$

Each accreted H releases

$$\approx \frac{GMm_{\text{H}}}{R} \approx 200 \text{ MeV.}$$

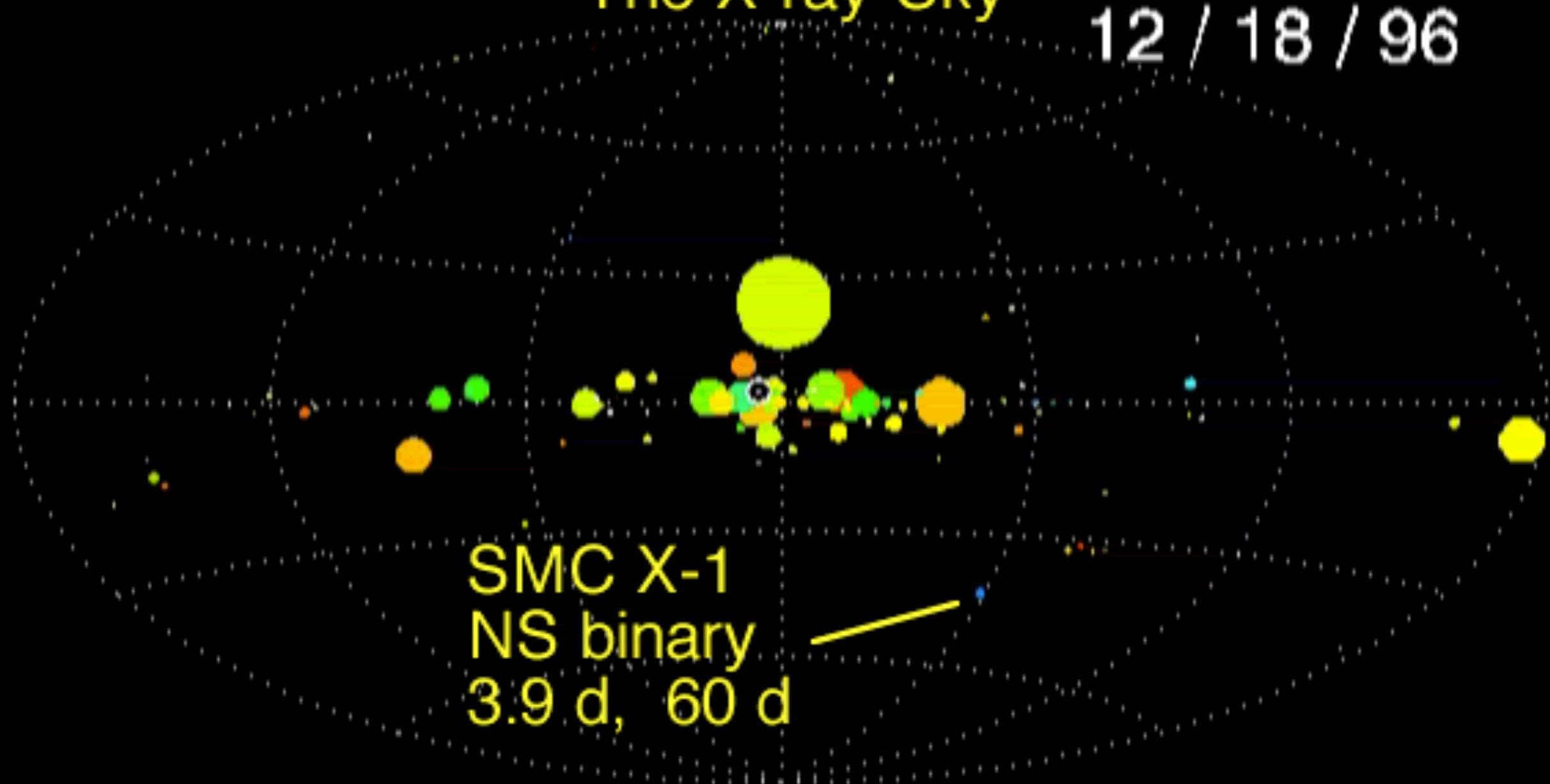
Fusing H to He releases

$$\approx 7 \text{ MeV}$$

per nucleon.

The X-ray Sky

12 / 18 / 96



Detection: Isolated neutron stars

about 1700 pulsars detected; about 50 are in binary systems with some mass information

very precise mass information;

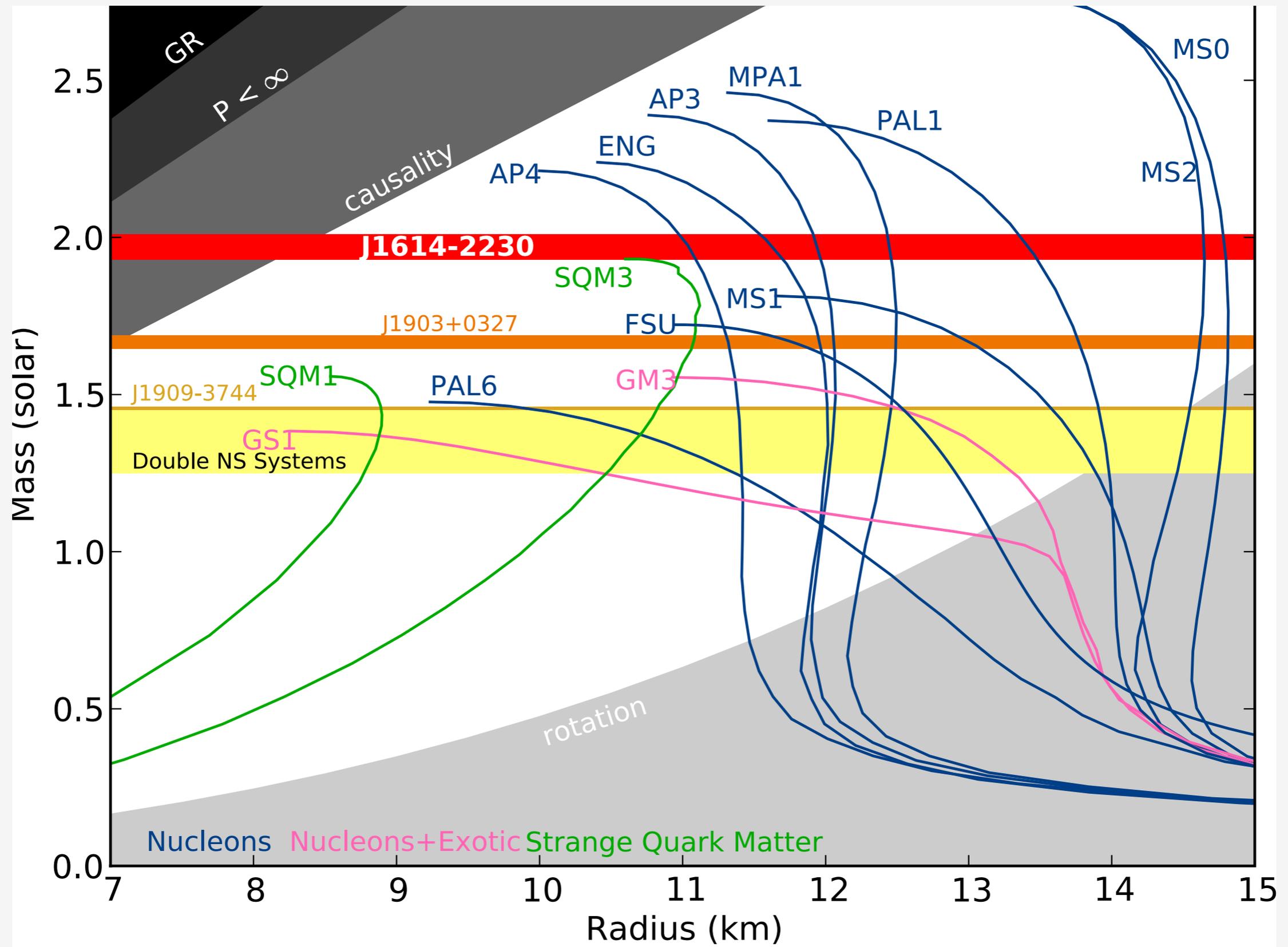
but, no radius information

fastest spin is 716 Hz (Hessels et al. 2006; faster than household blender!)



John Rowe Animation/Australia Telescope

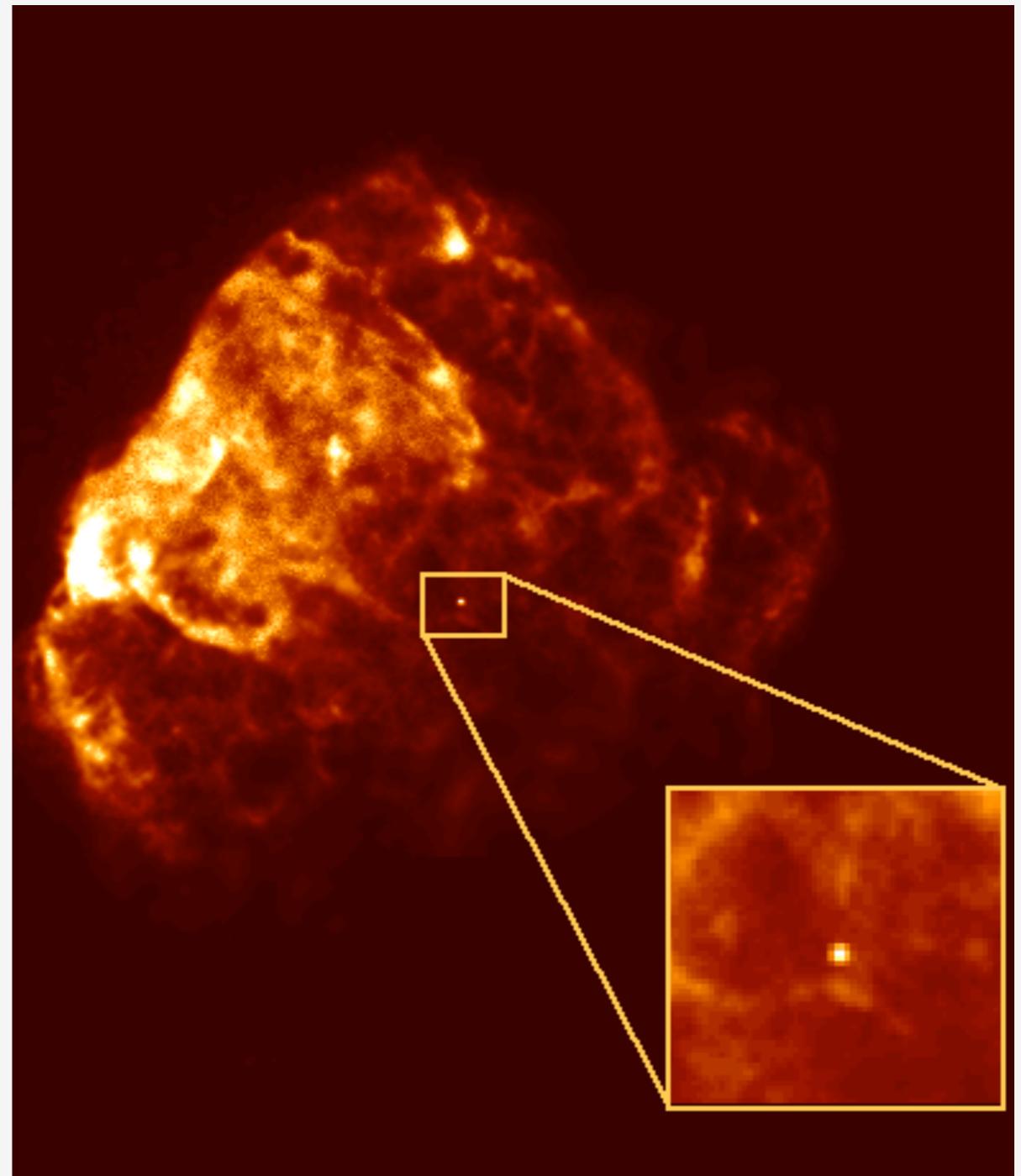
National Facility, CSIRO



Neutron stars can reach $2 M_{\text{sun}}$! Demorest et al. 2010

A neutron star cooling from its fiery birth

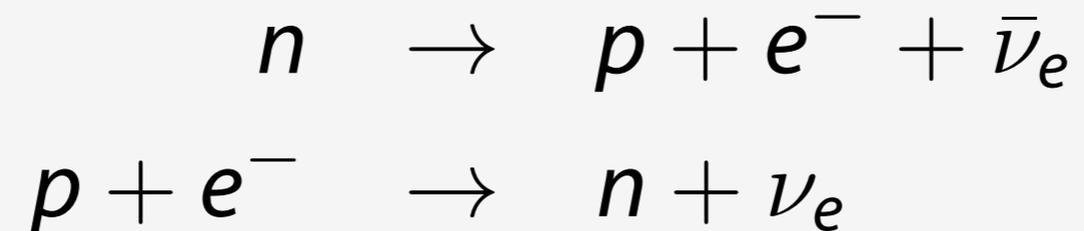
ROSAT Image of thermal emission
from neutron star in Puppis A
supernova remnant



cooling: the Urca process

Gamow & Schoenberg 1941

In npe-matter, maintain β -equilibrium via



Integration of the rate over phase space gives a T^6 dependence of the cooling rate.



but this is blocked...

Chiu & Salpeter, Bahcall & Wolff

mom. cons.

$$\rho_{F,n} < \rho_{F,e} + \rho_{F,p}$$

β -equil.

$$\mu_e = \mu_n - \mu_p$$

charge neut.

$$n_e = n_p$$

Need a “bystander” particle, e.g., $n + n \rightarrow n + p + e^- + \bar{\nu}_e$.

This rate is $> 10^6$ times slower at typical $T < 10^8$ K.

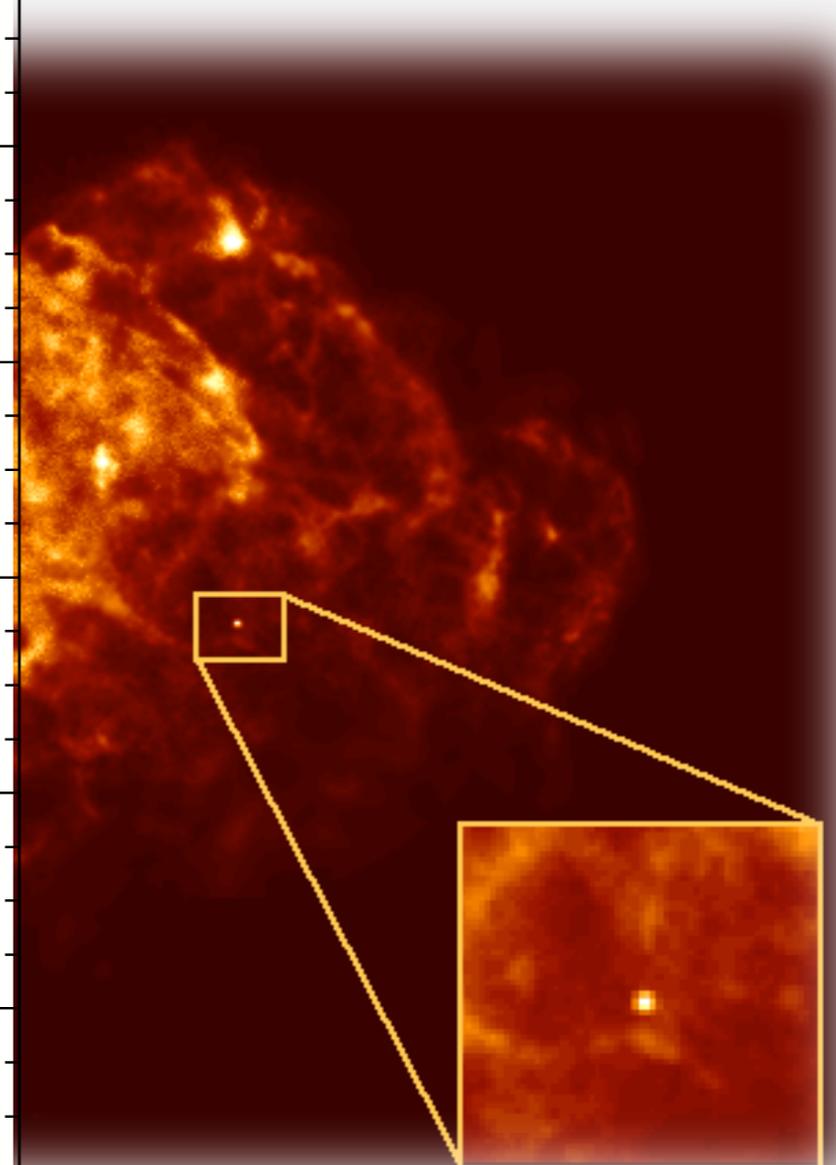
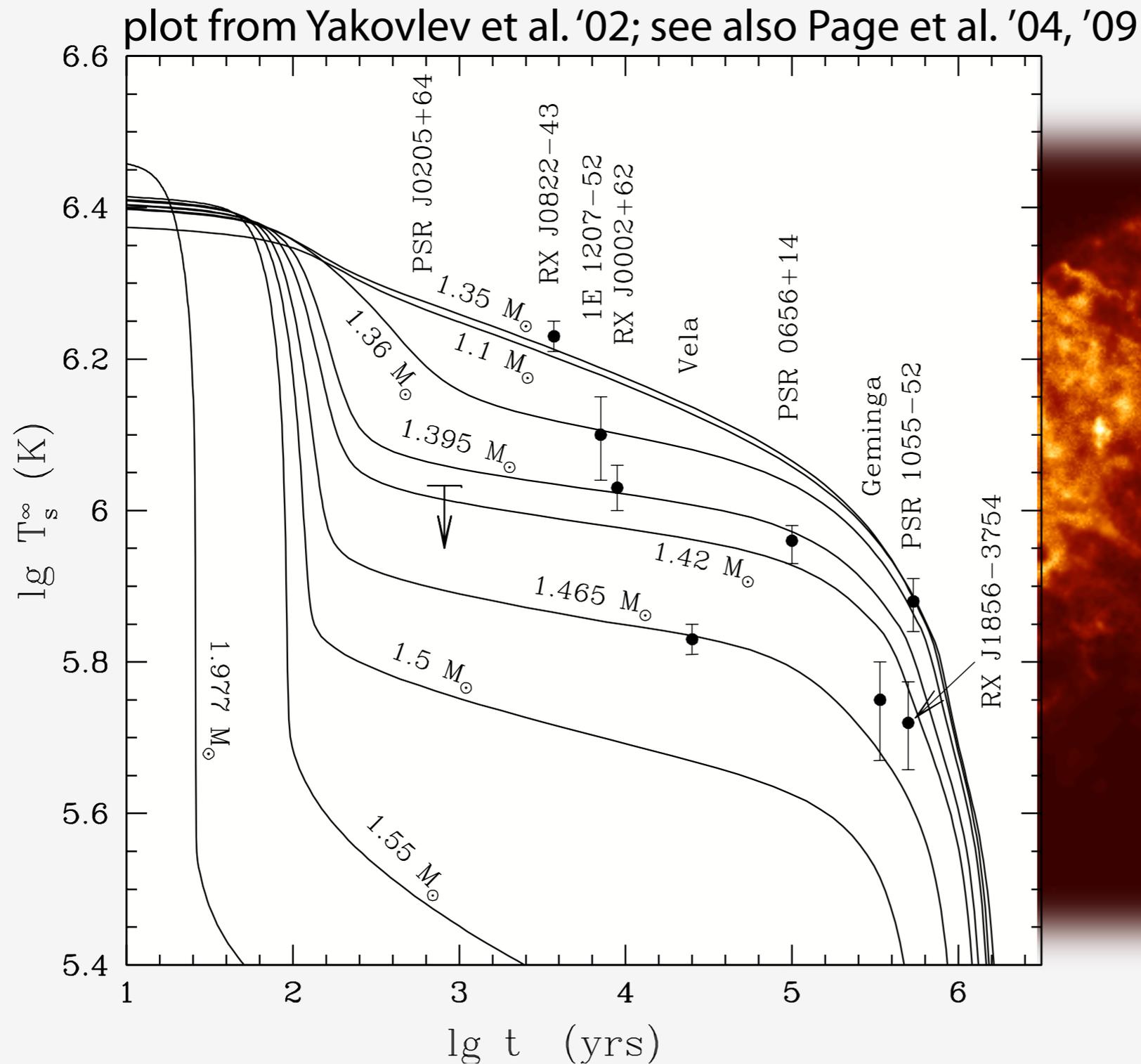
If $n_p/n > 0.11$, direct process can go.
Also if other channels, e.g.
hyperons, are available.

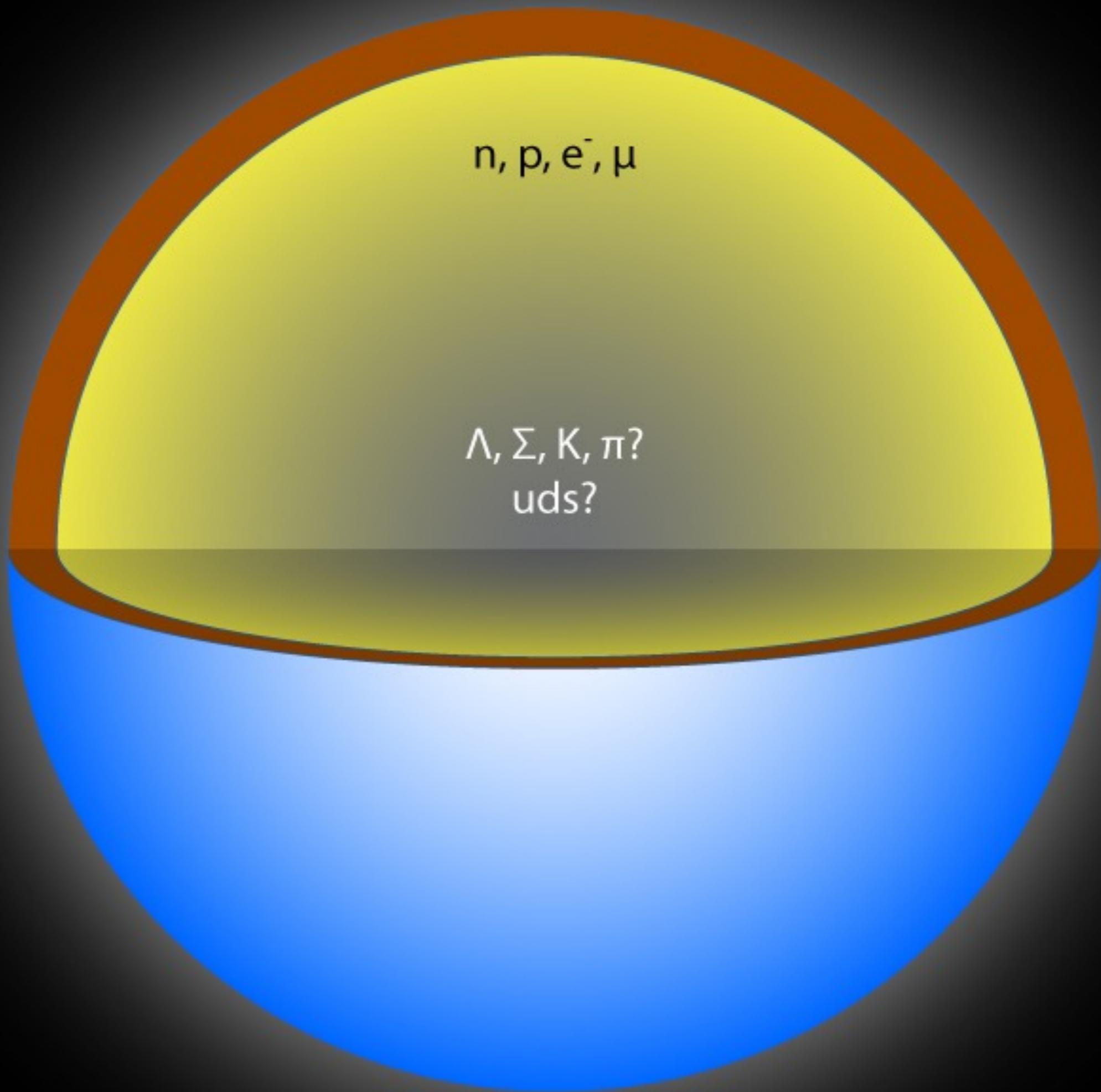
**A high symmetry energy implies
that neutron stars should cool
rapidly!**

—Lattimer & Prakash 2007



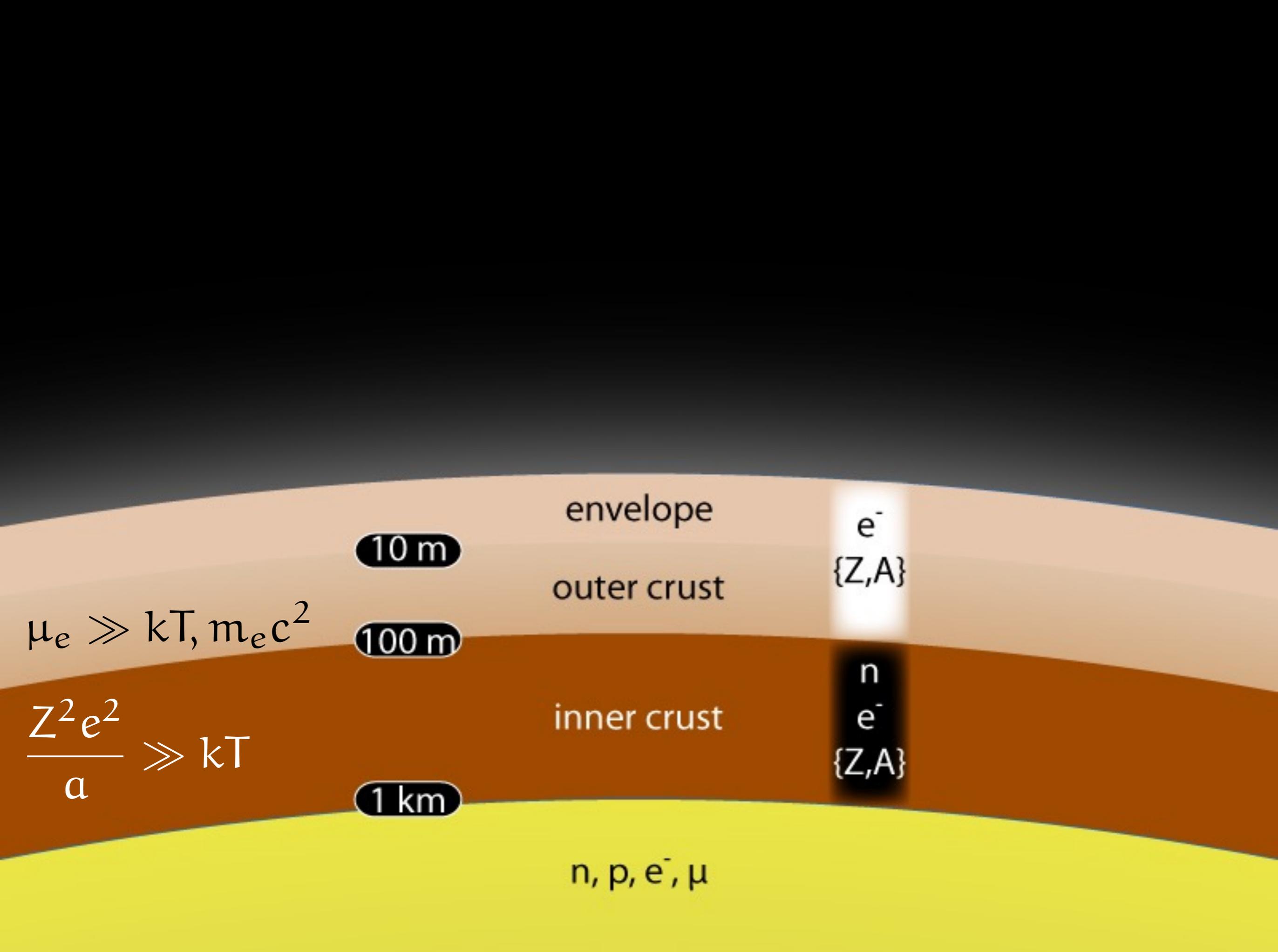
Neutron stars cooling from their fiery birth: “fast” and “slow” neutrino emissivities

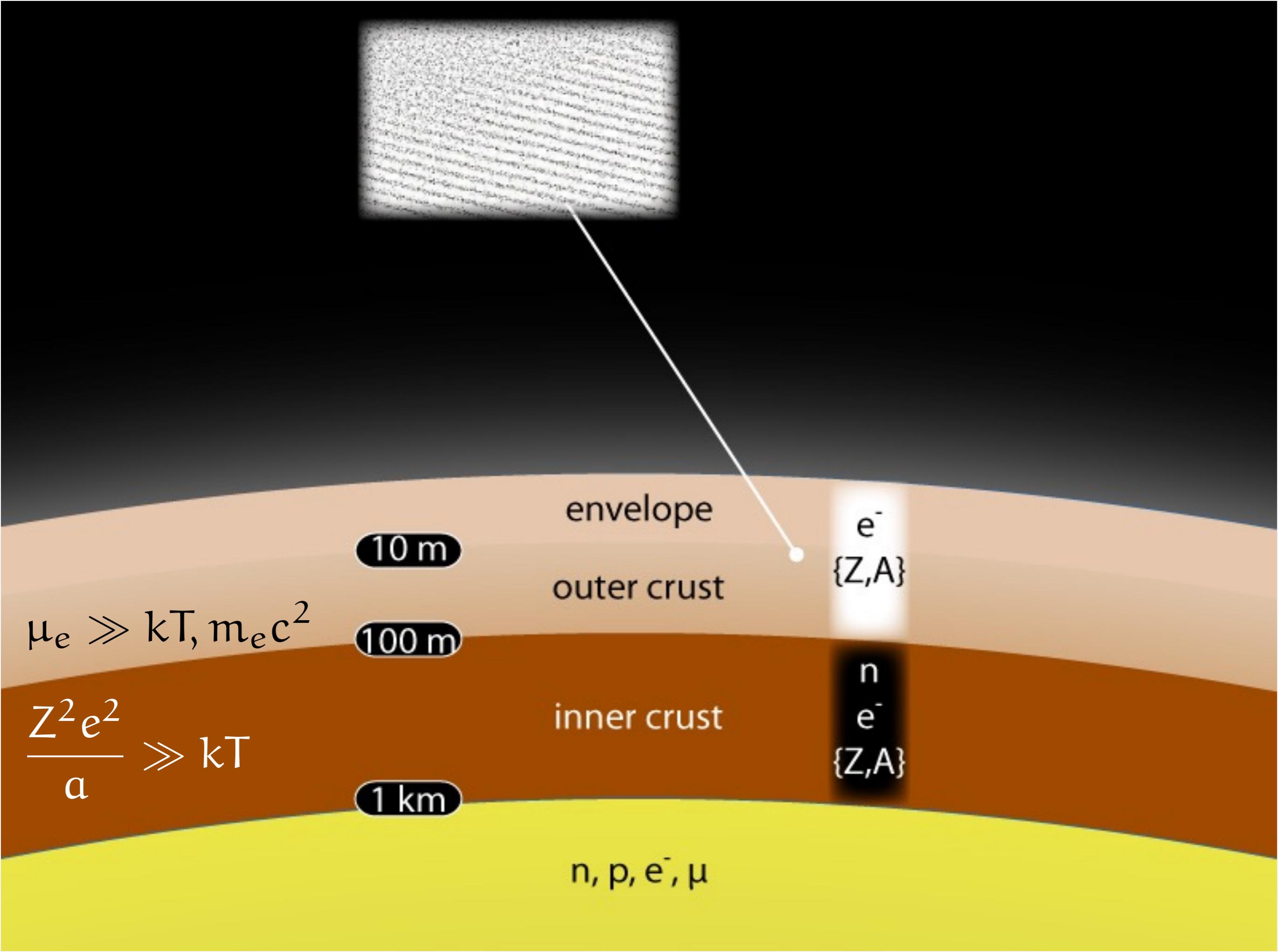
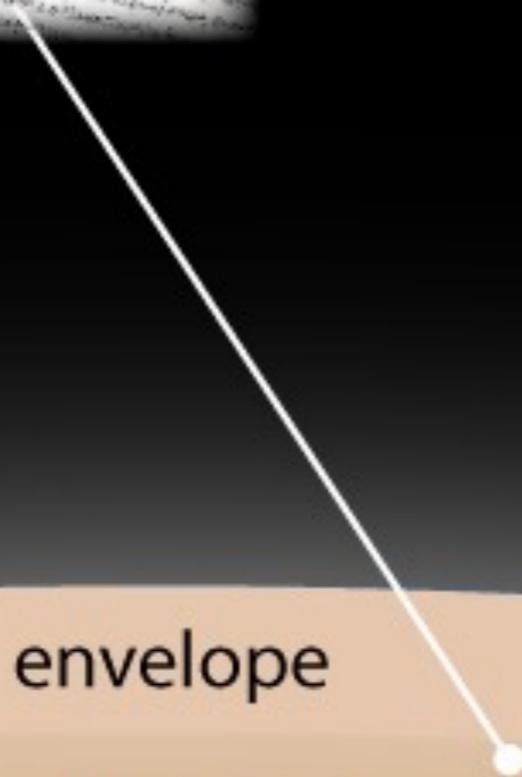
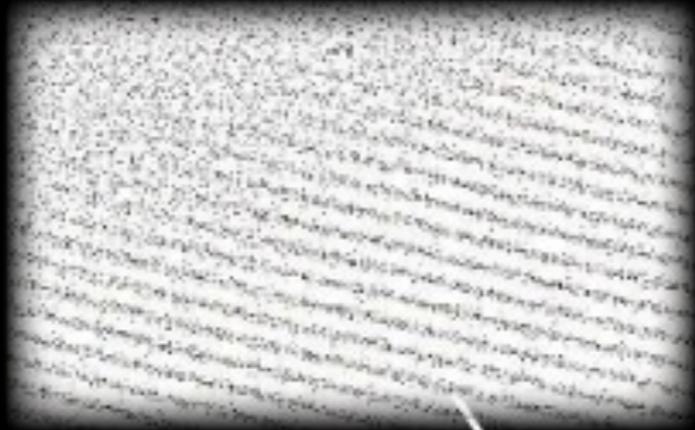




n, p, e^-, μ

$\Lambda, \Sigma, K, \pi?$
 $uds?$





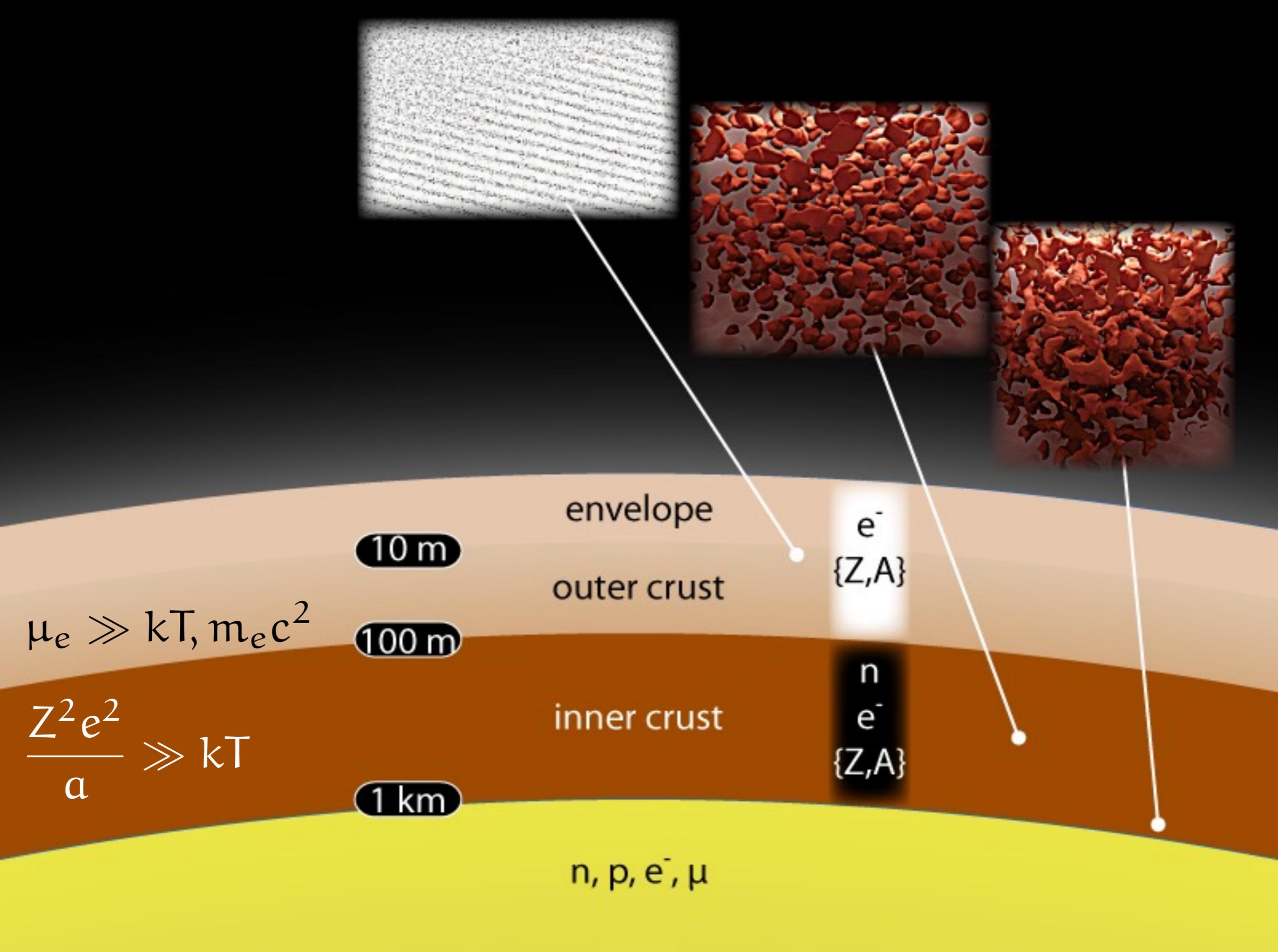
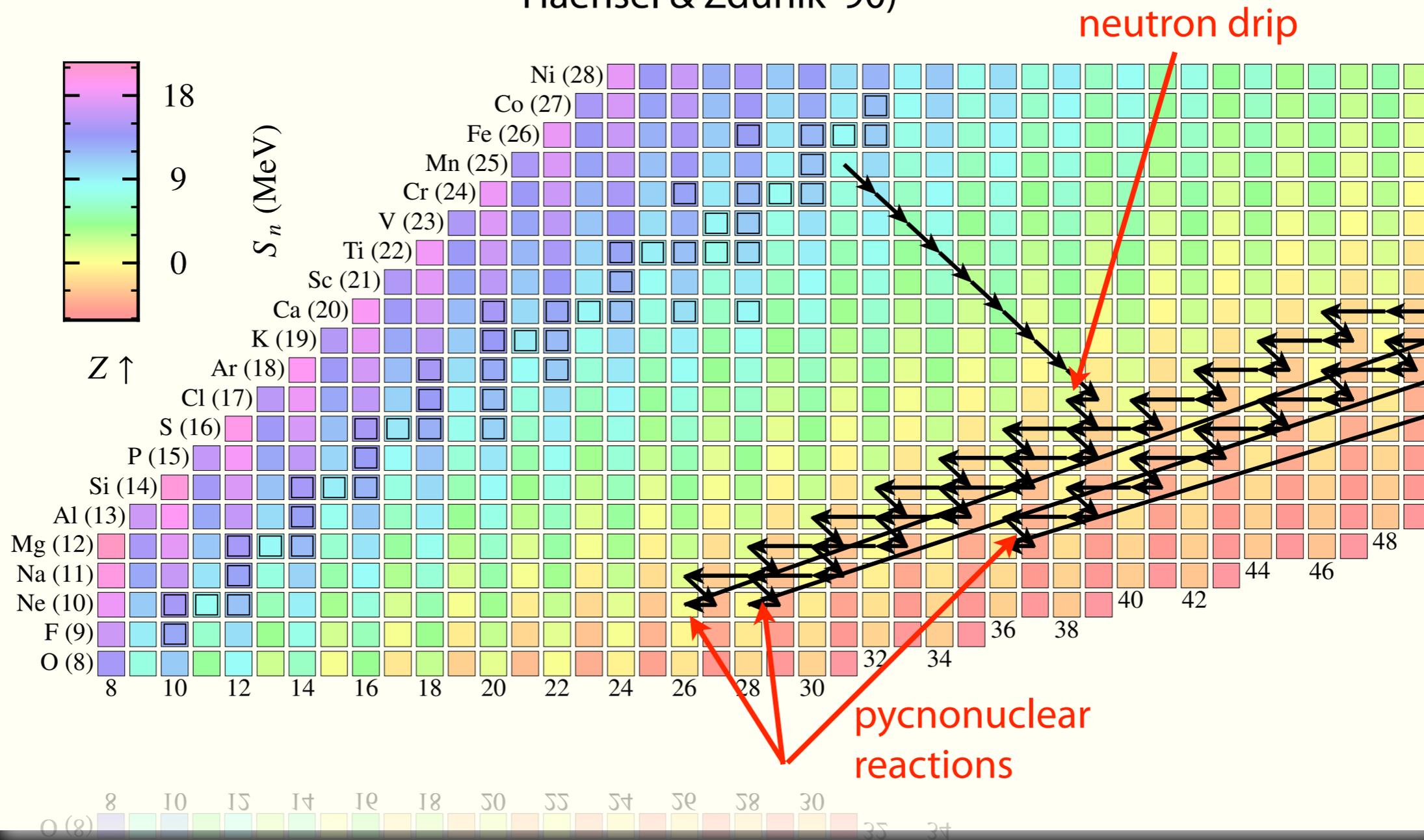


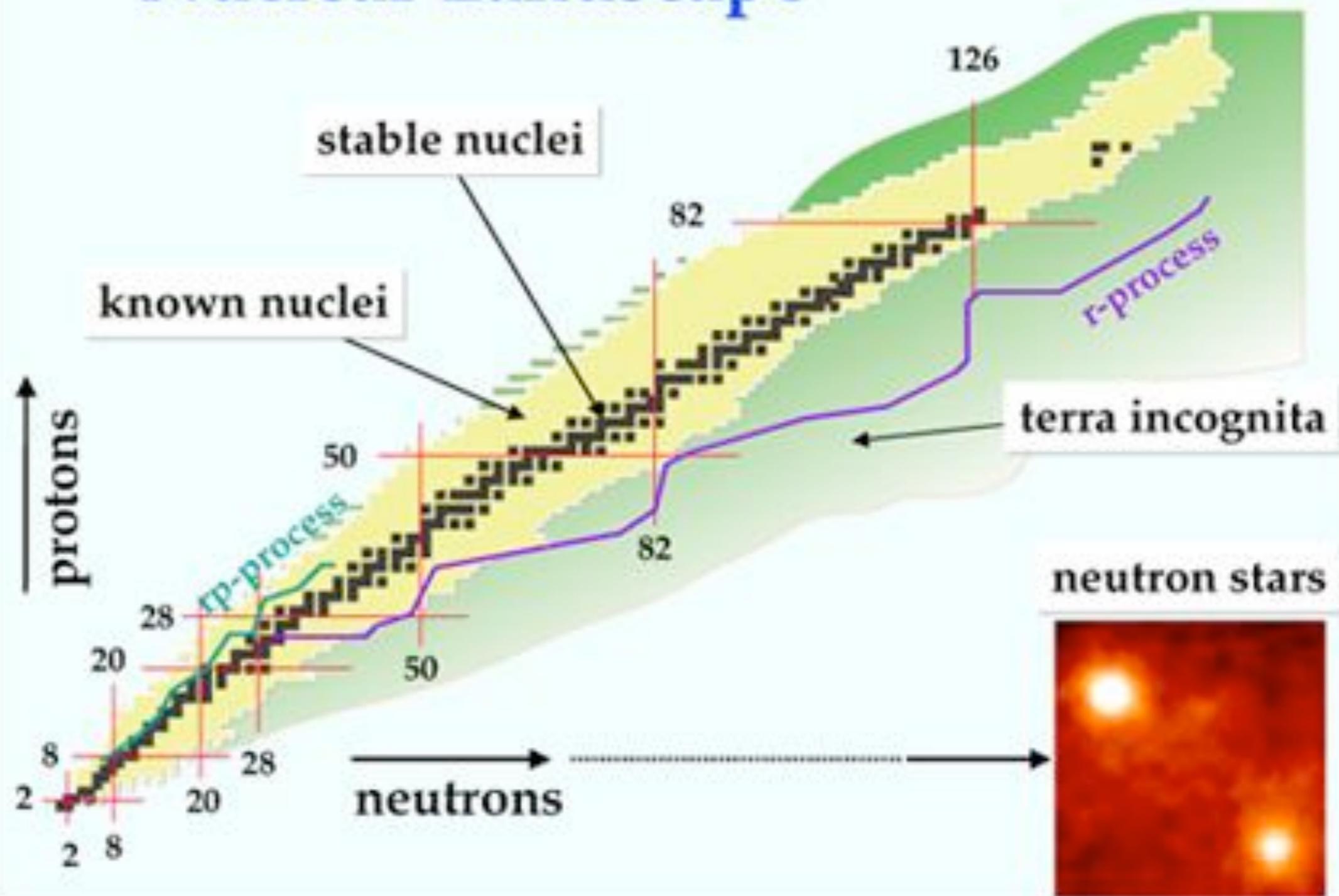
illustration with a simple liquid-drop model (Mackie & Baym '77, following Haensel & Zdunik '90)



crust reactions

Sato '79; Haensel & Zdunik '90; Gupta et al. '07; Steiner '12; Schatz et al. '13; Lau et al. (in prep)

Nuclear Landscape



Many of these reactions are within reach of FRIB

Disordered nuclear pasta, magnetic field decay, and crust cooling in neutron stars

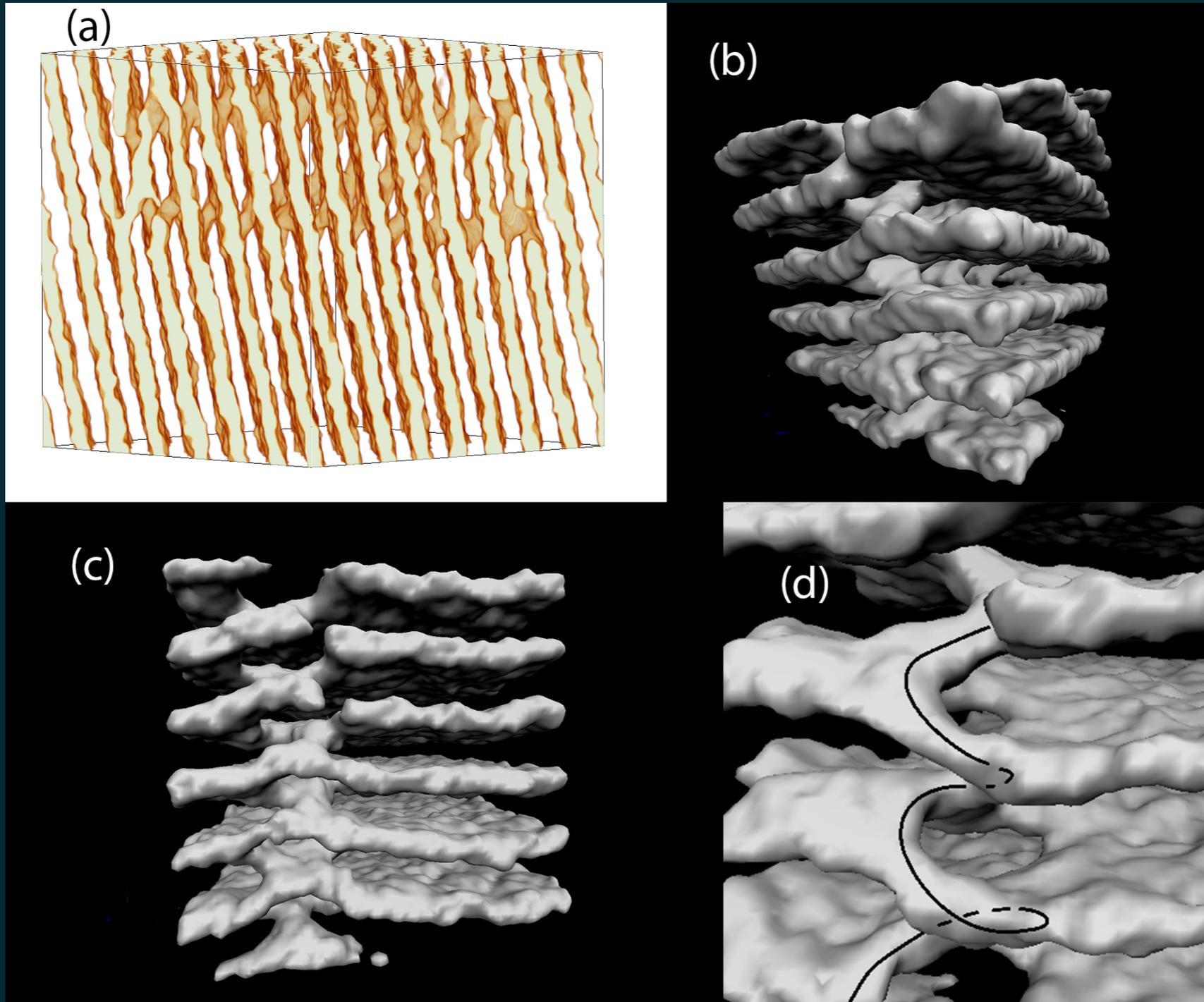
C. J. Horowitz,^{1,*} D. K. Berry,² C. M. Briggs,¹ M. E. Caplan,¹ A. Cumming,³ and A. S. Schneider¹

¹*Department of Physics and Center for the Exploration of Energy and Matter,
Indiana University, Bloomington, IN 47405, USA*

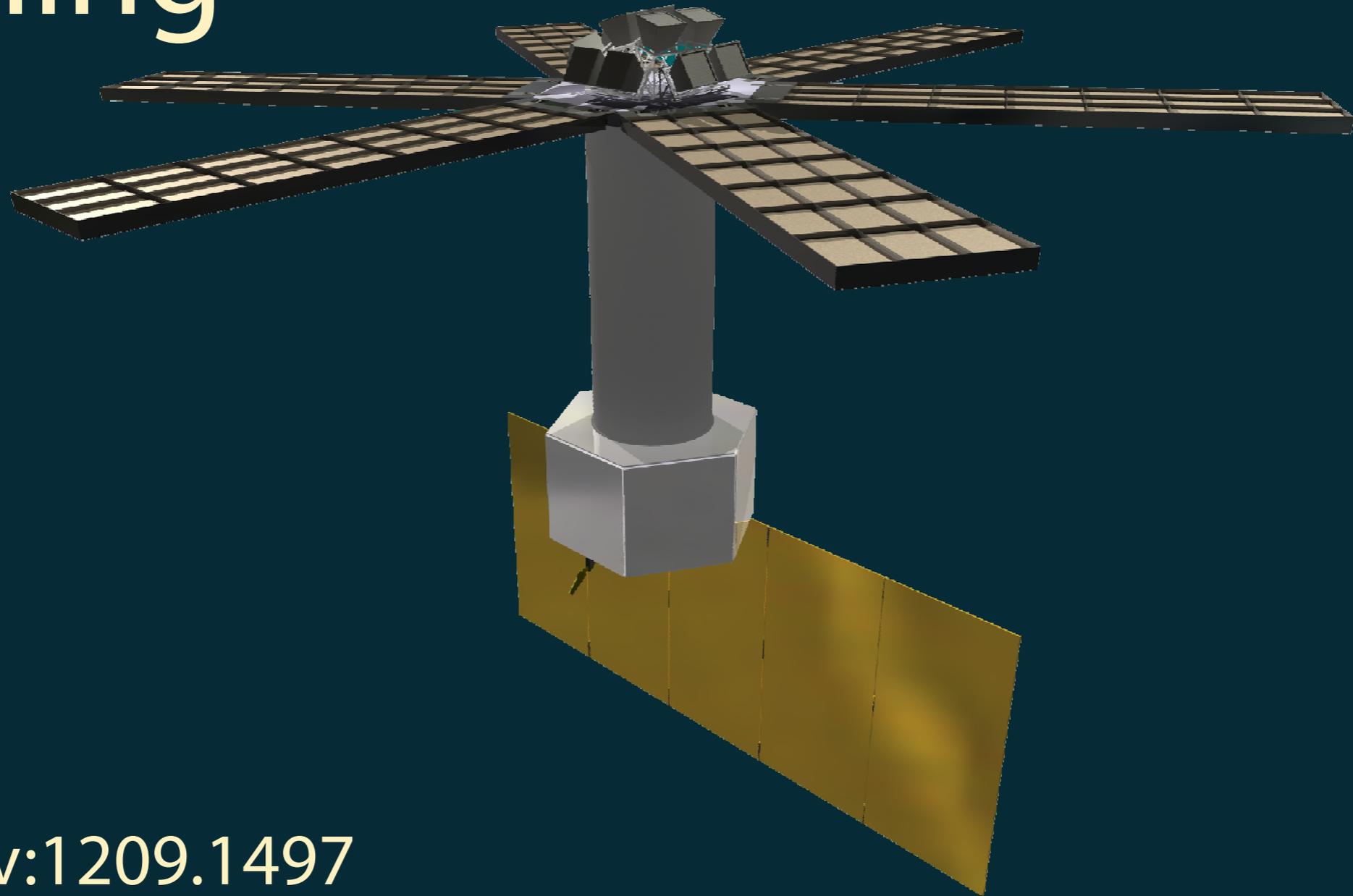
²*University Information Technology Services, Indiana University, Bloomington, IN 47408, USA*

³*Department of Physics, McGill University, 3600 rue University, Montreal QC, H3A 2T8 Canada*

(Dated: October 9, 2014)



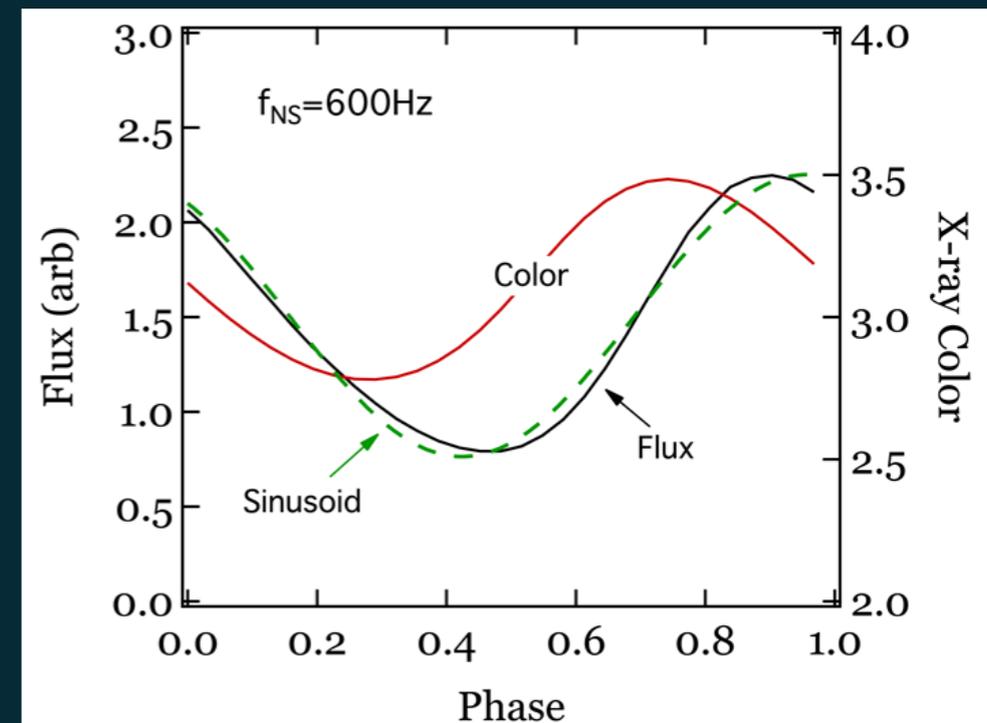
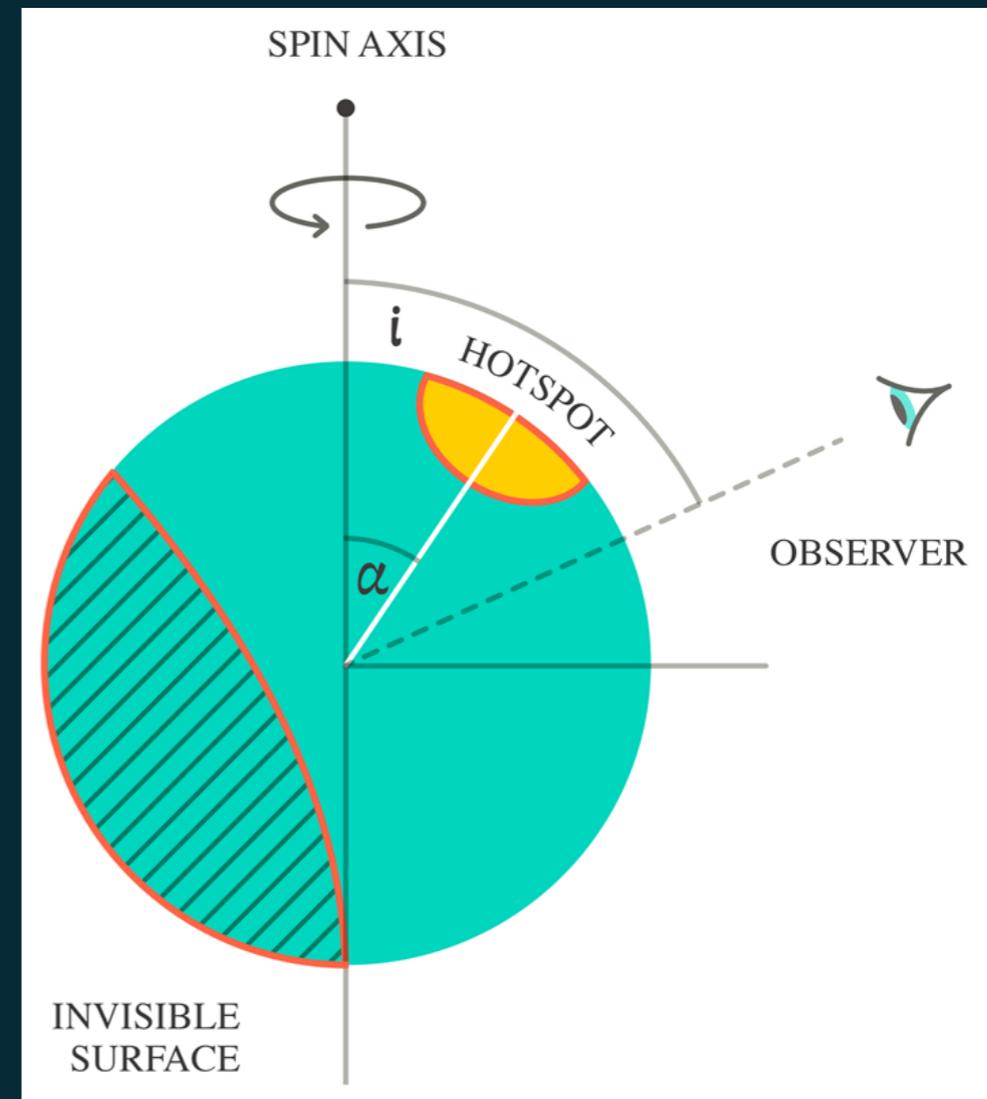
Large Observatory For x-ray Timing



[arxiv:1209.1497](https://arxiv.org/abs/1209.1497)

Figure 2. A pictorial view of the LOFT mission currently being studied by the LOFT Consortium, showing the deployed LAD panels, including the detector modules, as well as the 10 cameras comprising the 5 units of the WFM. The structure of 21 Modules in each Panel is shown, although the individual detector tiles are not visible. A structural tower supports the optical bench, with the service module supporting the solar panel array.

EOS from pulse phase spectroscopy





LIGO. Hanford, WA site: will it observe neutron stars?

Summary

- A number of observational probes of the nuclear EOS are available
 - pulsar masses
 - masses and radii from X-ray bursts, seismology
 - cooling of isolated neutron stars, thermal relaxation of accreting transients
 - gravitational wave emission
- No single observation is ideal and there are substantial systematic uncertainties—it's astrophysics; but
- These observations, taken together, offer interesting constraints and complement theoretical and experimental efforts in nuclear physics
- *Stay tuned! There are lots of opportunities to make advances in the next few years.*