



Giant Monopole Resonance Studies with HRS

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We know K_A from E_{GMR} :

$$E_{GMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

In an approximate way, K_A may be expressed as:

$$K_A \sim K_\infty (1 + cA^{-1/3}) + K_\tau ((N - Z)/A)^2 + K_{Coul} Z^2 A^{-4/3}$$

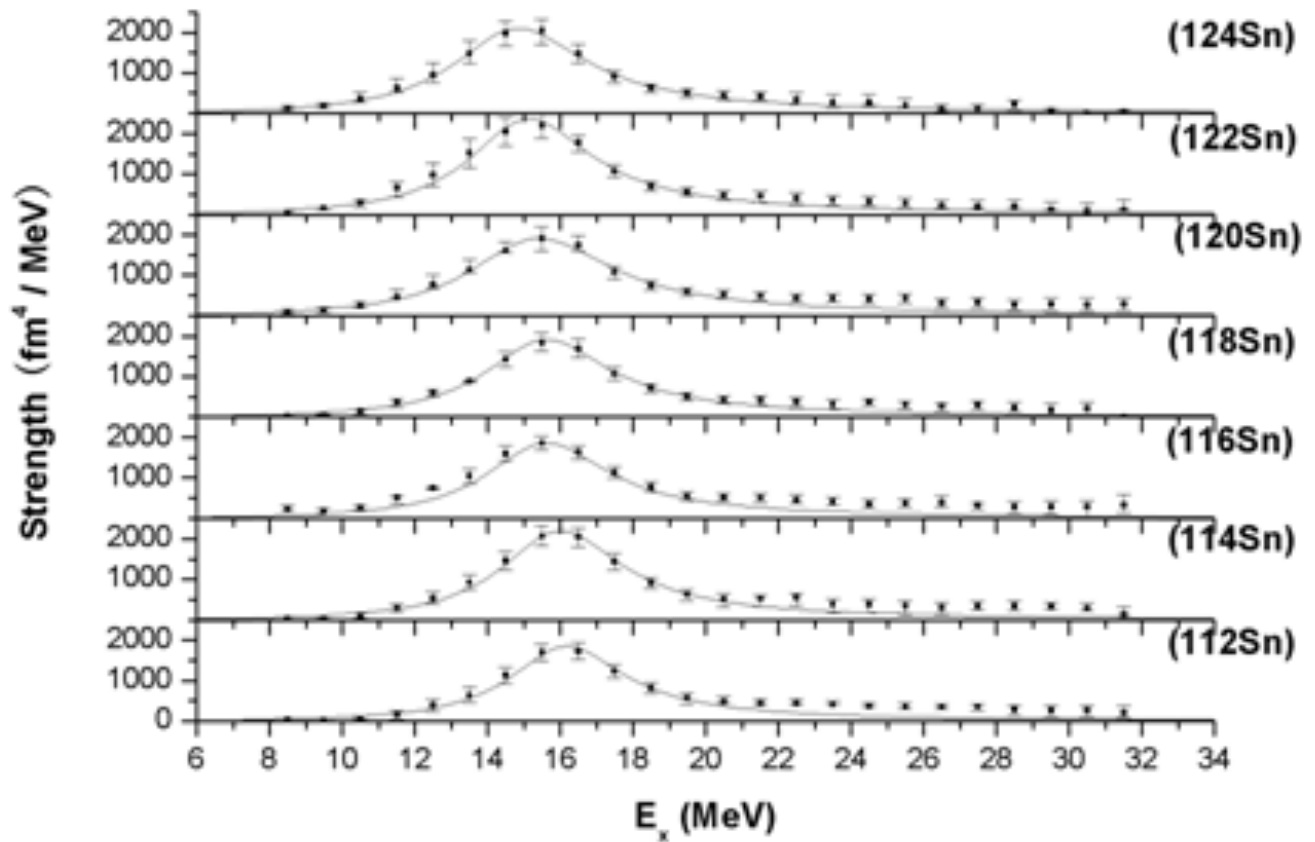
$$c \sim -1$$

K_{Coul} is, basically, model independent

$$K_\tau ??$$

Measurements over a series of isotopes gives K_τ

K_τ is critical in our understanding the properties of neutron stars



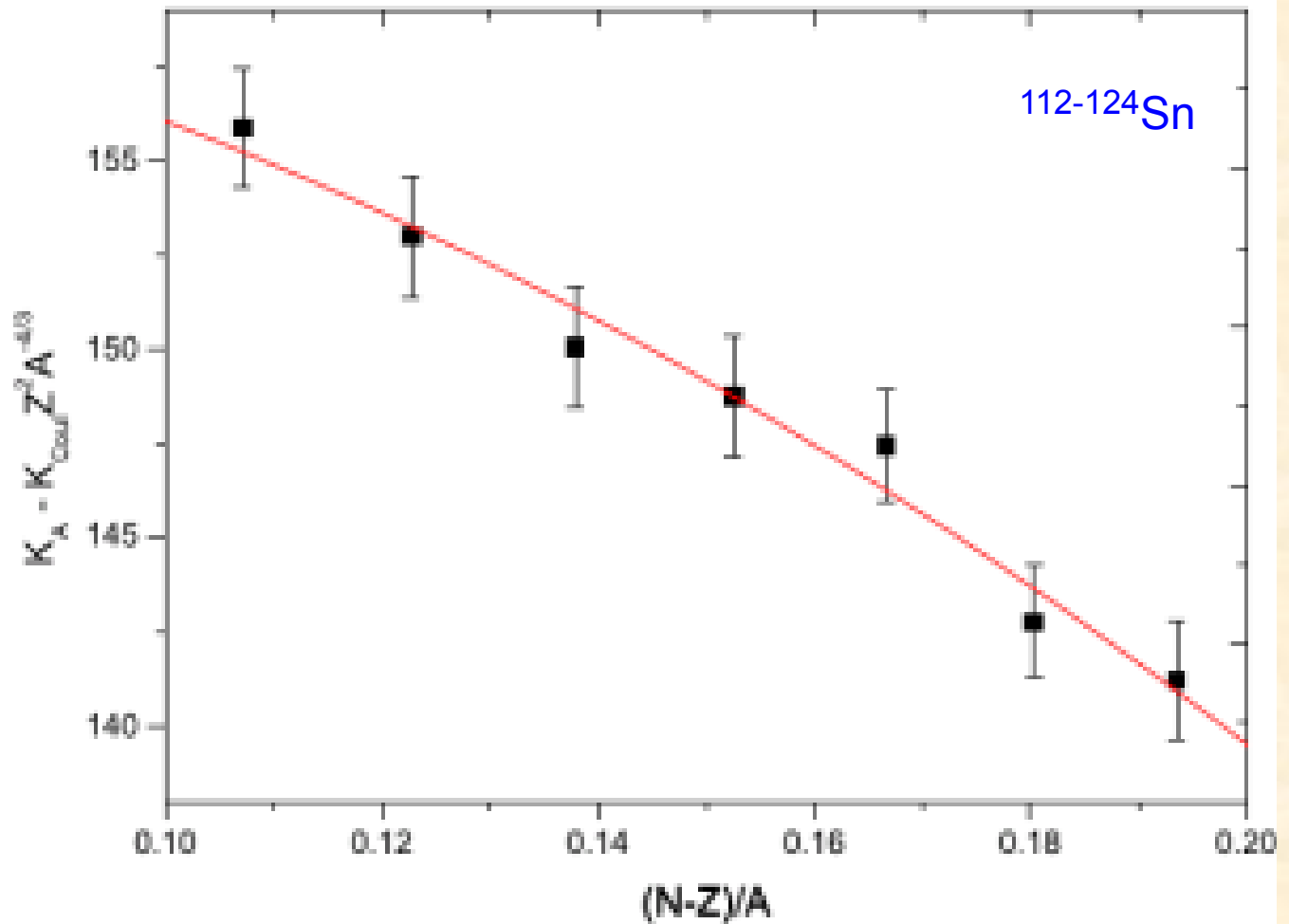


$$K_A \sim K_{\text{vol}} (1 + cA^{-1/3}) + K_{\tau} ((N - Z)/A)^2 + K_{\text{Coul}} Z^2 A^{-4/3}$$

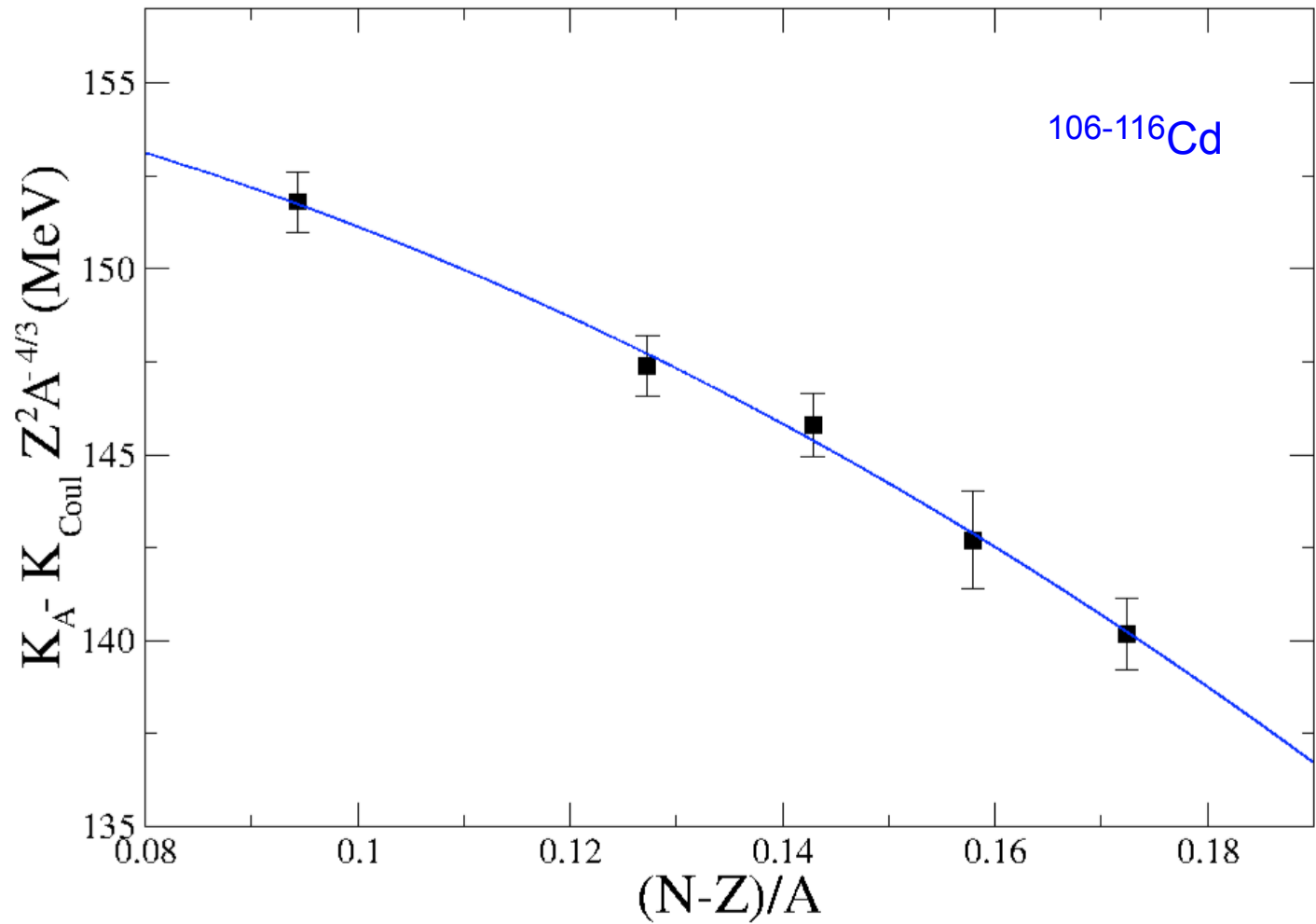
$$K_A - K_{\text{Coul}} Z^2 A^{-4/3} \sim K_{\text{vol}} (1 + cA^{-1/3}) + K_{\tau} ((N - Z)/A)^2$$

$$\sim \text{Constant} + K_{\tau} ((N - Z)/A)^2$$

We use $K_{\text{Coul}} = - 5.2 \text{ MeV}$ (from Sagawa)



$$K_{\tau} = -550 \pm 100 \text{ MeV}$$



$$K_{\tau} = -555 \pm 75 \text{ MeV}$$



Towards very neutron-rich nuclei

❖ K_{τ}

❖ K_{core} and K_{skin}

“soft GMR” akin to pigmy GDR’s.

❖ Need inverse reactions

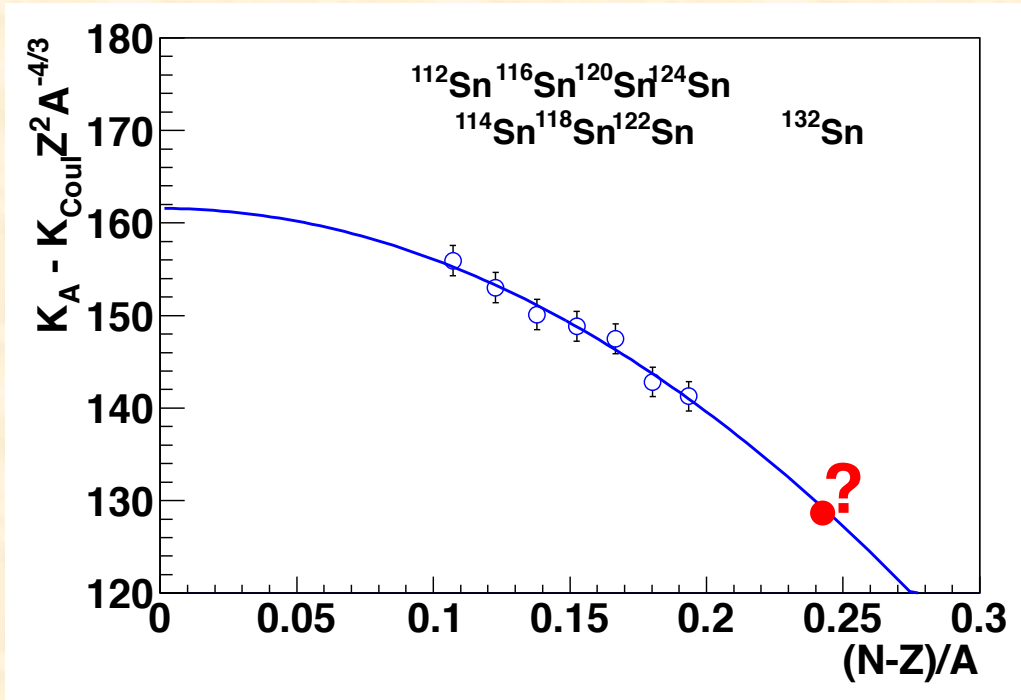
${}^2\text{H}$, ${}^4\text{He}$, or ${}^6\text{Li}$ targets

beams of 35-100 MeV/A

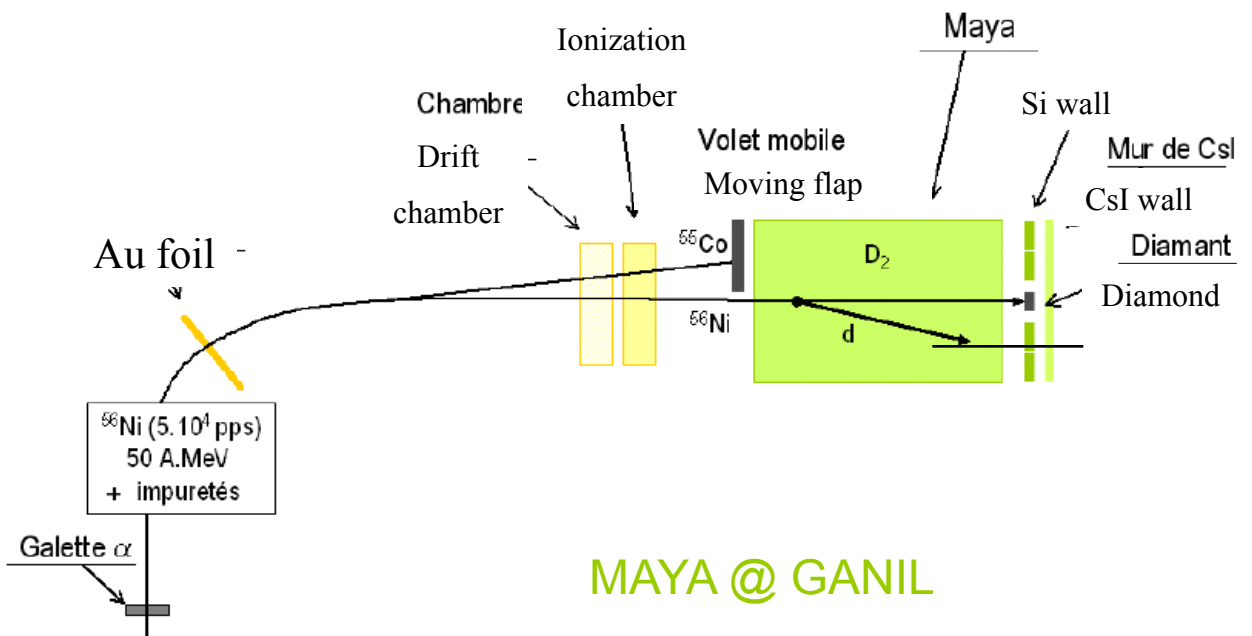
❖ First experiments performed at
GANIL

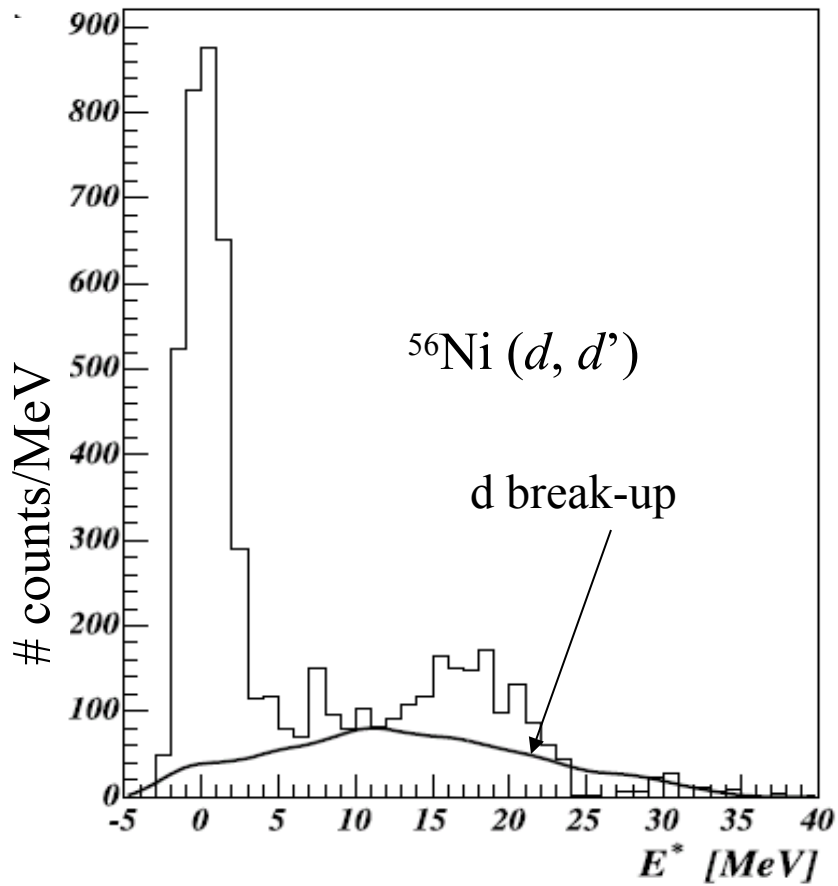
${}^{56}\text{Ni} + {}^2\text{H}$, ${}^{68}\text{Ni} + {}^4\text{He}$

with active target MAYA



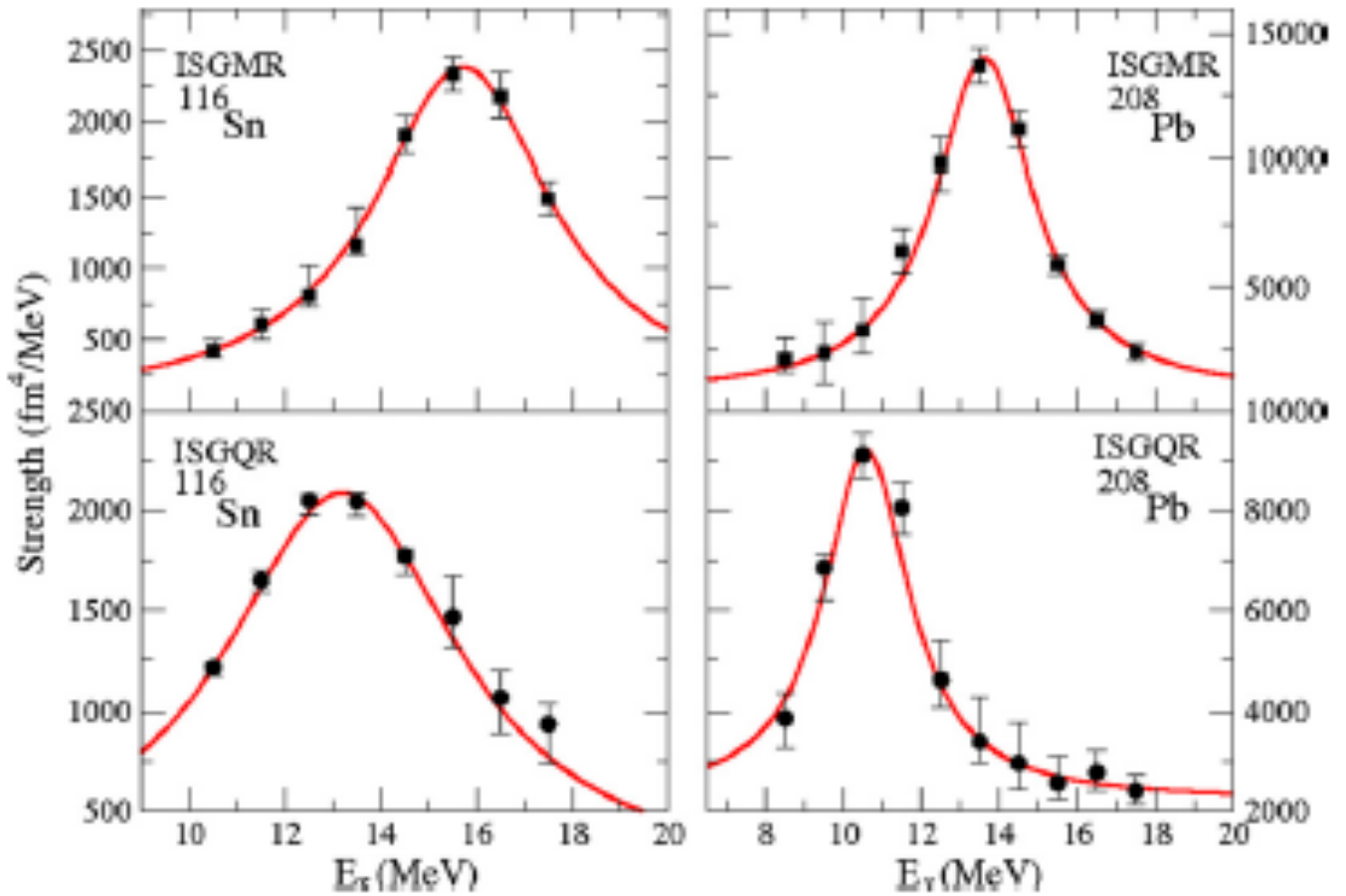
⇒ error of K_{τ} can be reduced to 50 MeV from 75+ MeV.



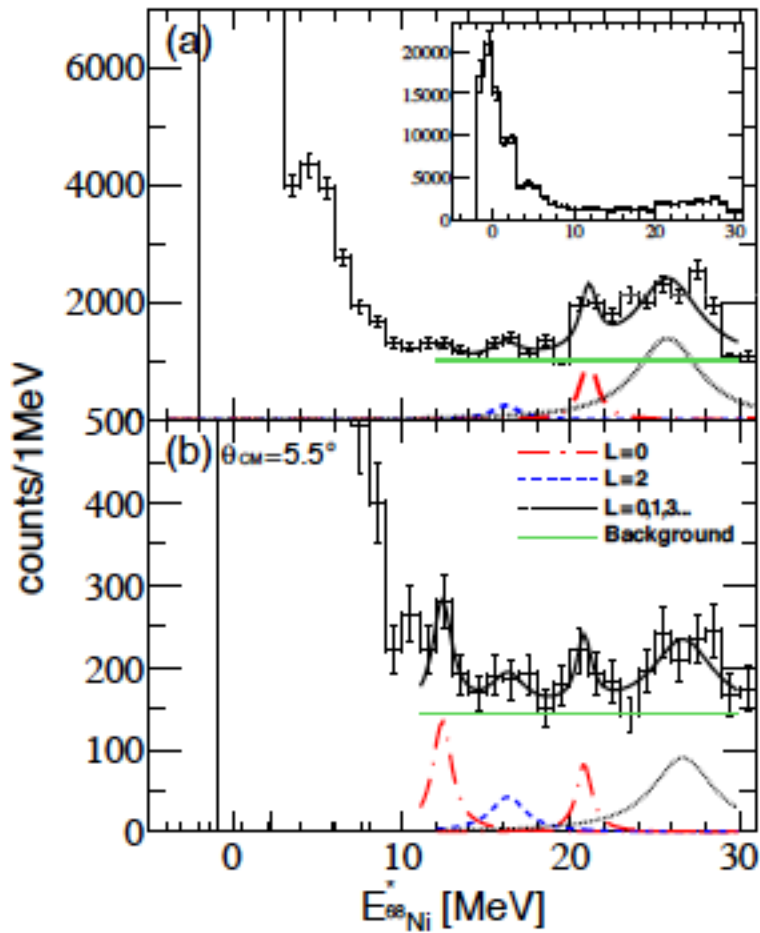


MAYA @ GANIL

C. Monrozeau *et al.*, Phys. Rev. Lett. 100, 042501 (2008)



(d, d') @ 100 MeV/A



$^{68}\text{Ni} (\alpha, \alpha') @ 50 \text{ MeV/A}$

M. Vandebrouck *et al.*, Phys. Rev. Lett. (to appear)



So, how can HRS help?

- ✓ For light nuclei (up to $Z \sim 28$), one can use HRS in coincidence with charged-particle and neutron decays in forward detectors.
- ✓ Provide cleaner spectra by gating on beam-like particles.



धन्यवाद

Thanks!



